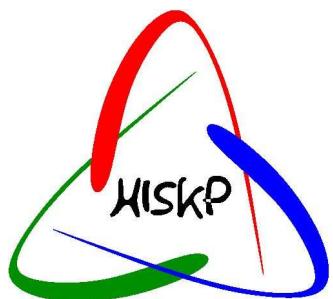


# Hadron Spectroscopy:

## Results and Ideas.



SFB/TR16



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International Workshop on  
NEW PARTIAL WAVE ANALYSIS TOOLS  
FOR NEXT GENERATION HADRON SPECTROSCOPY EXPERIMENTS

June 20-22, 2012 Camogli, Italy

# **Hadron Spectroscopy: results and ideas**

- 1. Why baryon spectroscopy**
- 2. New results and PDG related issues**
- 3. Baryon resonance spectrum: interpretation**

Parity doublets and chiral symmetry

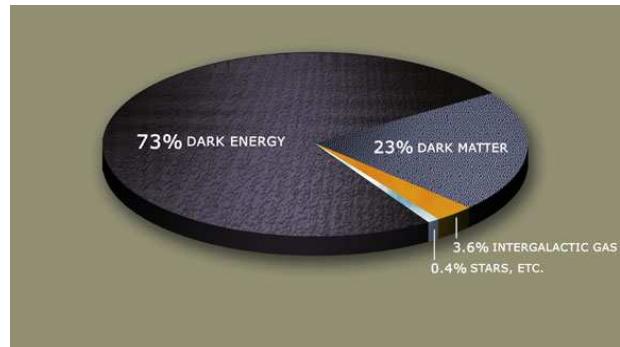
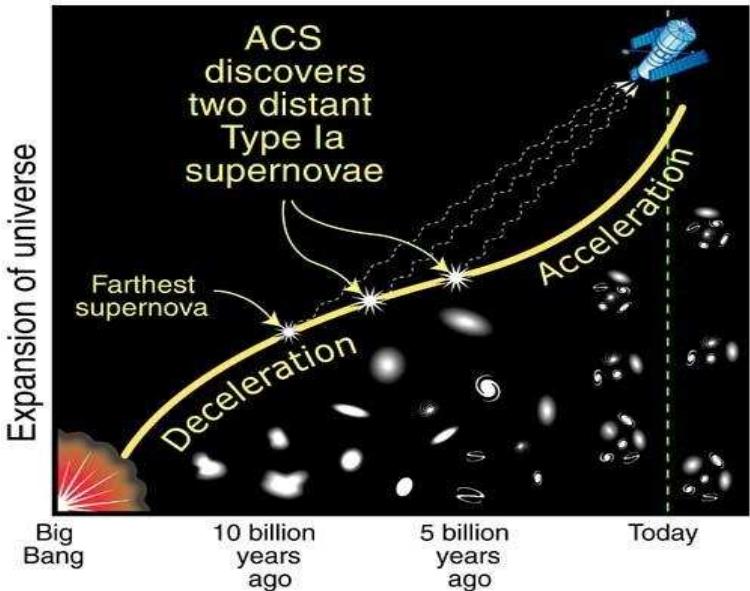
AdS/QCD

$SU(6) \otimes O(3)$

Dynamically generated resonances versus C.D.D. poles

- 4. New results from ELSA**
- 5. Summary**

# 1. Why Baryon Spectroscopy?

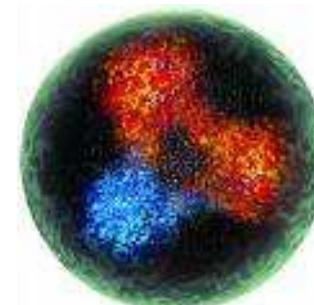


- **Mass of the Universe:**

Dark energy	73%
Dark matter	23%
Intergalactic gas	3.6%
Stars	0.4% } atoms

- **Mass of atoms**

Mass of quarks	1%
Mass of electrons	0.1%
Field energy	99%

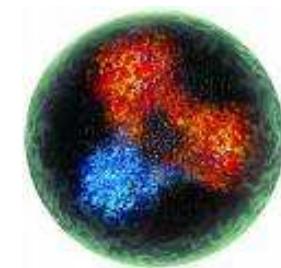


What are ↑ these objects?

- What are constituent quarks?

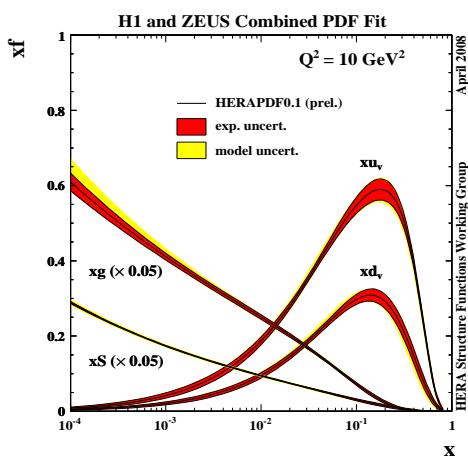
Mass is field energy!

$$M_N = \langle N | -\frac{9\alpha_s}{4\pi} (\mathbf{B}^2 - \mathbf{E}^2) + \sum_{\text{flavors}} m_i \bar{\psi}_i \psi_i | N \rangle .$$



- How to explore constituent quarks?

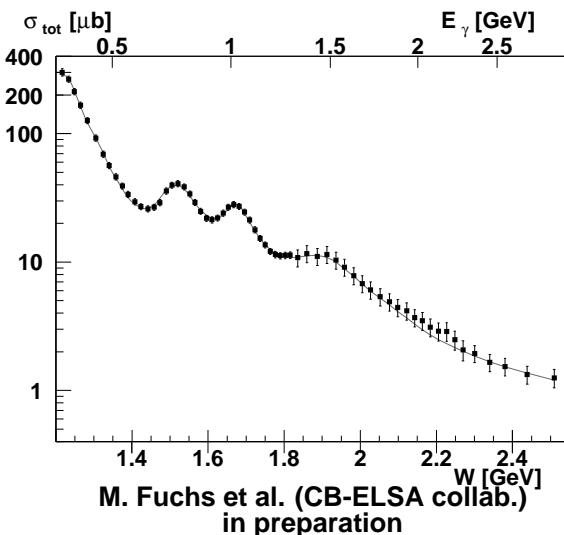
### Nucleon tomography:



**Goal: Distribution of linear and angular momenta of gluons and quarks in space**

Collectivity is lost!

### Nucleon spectroscopy:



Wave length matches the size of constituents

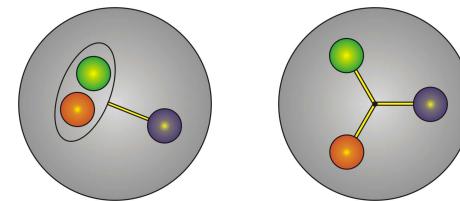
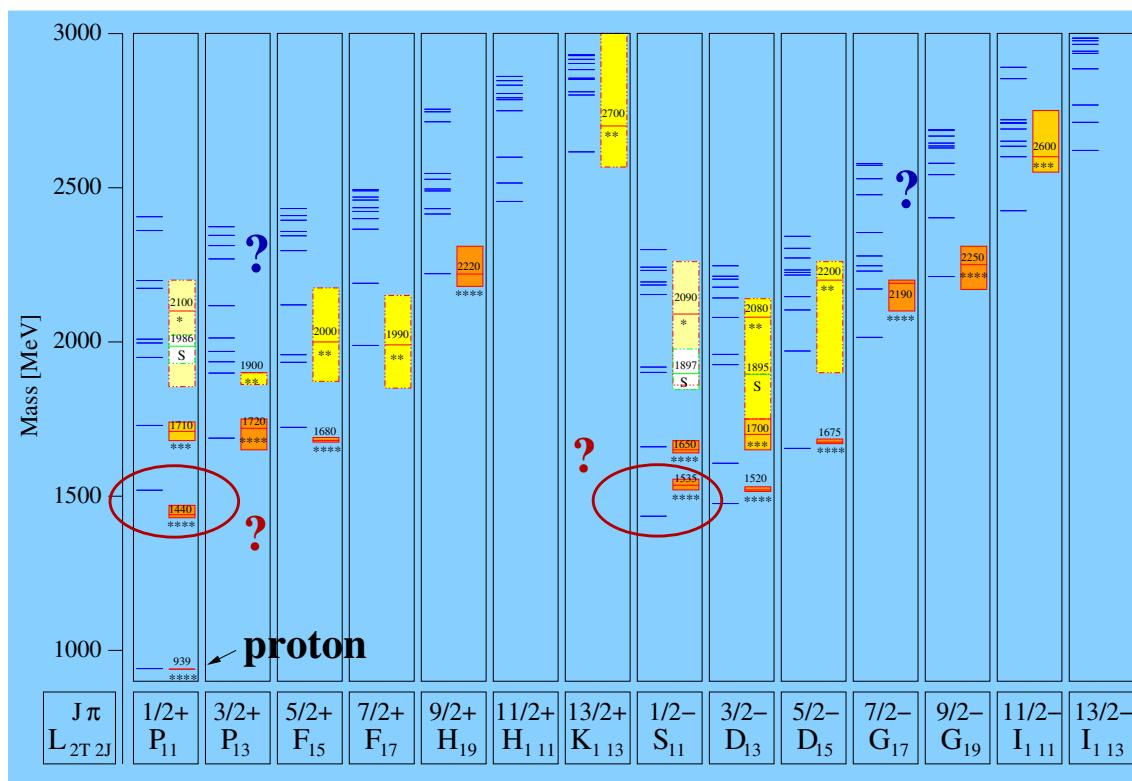
Explores collective response!

## Questions:

- Are constituent quarks the relevant degrees of freedom ?
- What are the effective forces between them ?

## Findings:

- Linear Regge trajectories
- Many more resonances expected in symmetric quark models than observed



Diquark substructure?  
QCD strings?      AdS/QCD?

H. Forkel and E. Klempt,  
Phys. Lett. B 679, 77 (2009).

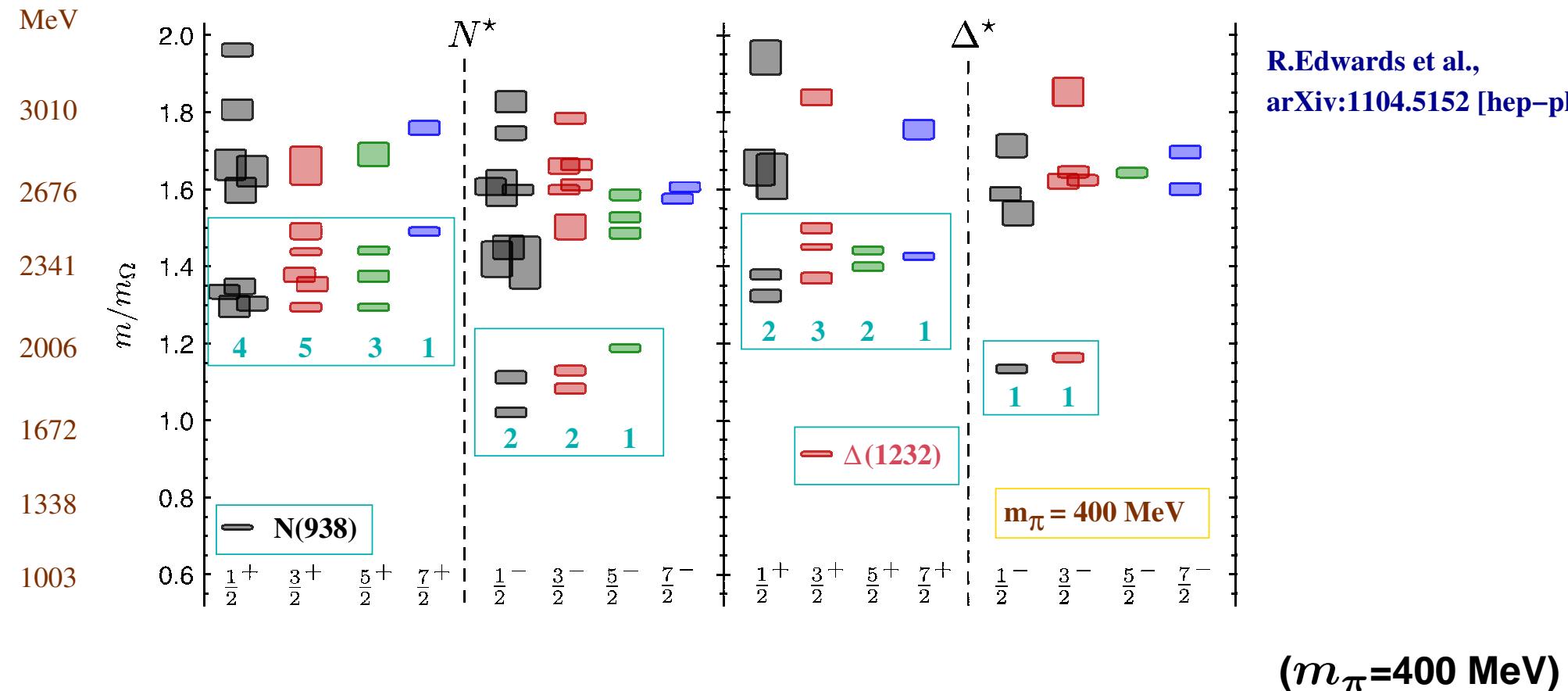
## Non-strange $N^*$ -resonances

U. Löring, B. Metsch, H. Petry, Eur. Phys. J. A 10, 395 (2001).

Constituent quarks  
confinement potential  
+ residual interaction

↔ wrong degrees of freedom ?

# Excited baryons from Lattice QCD



Exhibits the features  $SU(6) \otimes O(3)$ -symmetry

- Counting of levels consistent with non-rel. quark model
- Striking similarity with quark model
- No parity doubling

The lattice seems not to solve the problems!

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## 2. New results and PDG related issues

J. Beringer et al. (Particle Data Group), Phys. Rev. D86, 010001 (2012).

Resonance	Rating	$N_{PP}$	Resonance	Rating	$N_{PP}$	Resonance	Rating	$N_{PP}$
$N(1440)1/2^+$	****	13	$N(1520)3/2^-$	****	17	$N(1535)1/2^-$	****	15
$N(1650)1/2^-$	****	18	$N(1675)5/2^-$	****	14	$N(1680)5/2^+$	****	17
$N(1685)$	*		$N(1700)3/2^-$	***	15	$N(1710)1/2^+$	***	14
$N(1720)3/2^+$	****	17	$N(1860)5/2^+$	**	9	$N(1875)3/2^-$	***	16
$N(1880)1/2^+$	**	20	$N(1895)1/2^-$	**	17	$N(1900)3/2^+$	***	18
$N(1990)7/2^+$	**	9	$N(2000)5/2^+$	**	11	$N(2040)3/2^+$	*	
$N(2060)5/2^-$	**	13	$N(2100)1/2^+$	*		$N(2150)3/2^-$	**	11
$N(2190)7/2^-$	****	11	$N(2220)7/2^-$	****	7	$N(2250)9/2^-$	****	
$N(2600)11/2^-$	***		$N(2700)13/2^+$	**				
$\Delta(1232)$	****	8	$\Delta(1600)3/2^+$	***	12	$\Delta(1620)1/2^-$	****	10
$\Delta(1700)3/2^-$	****	11	$\Delta(1750)1/2^+$	*		$\Delta(1900)1/2^-$	**	13
$\Delta(1905)5/2^+$	****	11	$\Delta(1910)1/2^+$	****	13	$\Delta(1920)3/2^+$	***	21
$\Delta(1930)5/2^-$	***		$\Delta(1940)3/2^-$	*	5	$\Delta(1950)7/2^+$	****	13
$\Delta(2000)5/2^+$	**		$\Delta(2150)1/2^-$	*		$\Delta(2200)7/2^-$	*	
$\Delta(2300)9/2^+$	**		$\Delta(2350)3/2^-$	*		$\Delta(2390)7/2^+$	*	
$\Delta(2420)11/2^+$	****		$\Delta(2400)9/2^-$	****		$\Delta(2750)13/2^-$	**	
$\Delta(2950)15/2^+$	**							

E.g.: V. Kuznetsov et al., Phys. Lett. B 647, 23 (2007); V. Kuznetsov et al., Phys. Rev. C 83, 022201 (2011); I. Jaegle et al., Eur. Phys. J. A 47, 89 (2011).

M. Ablikim et al. [BES Collaboration], Phys. Rev. D 80, 052004 (2009).

A. V. Anisovich, R. Beck, E. Klempf, V. A. Nikonov, A. V. Sarantsev and U. Thoma, Eur. Phys. J. A 48, 15 (2012);

$N_{PP}$  particle properties were determined; 400 in total. Be cautious, there are ambiguities !

Promoted to three-star resonance

# PDG related issues

## 1. Naming scheme:

- $\Delta^{++}(2420)$ : Need to know quantum numbers “by heart”
- $\Delta^{++}(2420)H_{3,11}$ : Refers to  $\pi^+ p$  scattering. Need to know that  $H$  corresponds to  $L = 5$ . Not straightforward to calculate  $N\omega$  partial wave (which is also  $L = 5$ ).
- $\Delta^{++}(2420)11/2^+$  is self explanatory and used in the new PDG edition.

## 2. Baryon part in PDG emphasizes properties of poles rather than of Breit-Wigner parameters

## 3. Branching ratios often not defined

**$a_1(1260)$**  [k]

$I^G(J^{PC}) = 1^-(1^{++})$

Mass  $m = 1230 \pm 40$  MeV [l]

Full width  $\Gamma = 250$  to 600 MeV

<b><math>a_1(1260)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )	$p$ (MeV/c)
$(\rho\pi)_{S-\text{wave}}$	seen	353
$(\rho\pi)_{D-\text{wave}}$	seen	353
$(\rho(1450)\pi)_{S-\text{wave}}$	seen	†
$(\rho(1450)\pi)_{D-\text{wave}}$	seen	†
$\sigma\pi$	seen	—
$f_0(980)\pi$	not seen	179
$f_0(1370)\pi$	seen	†
$f_2(1270)\pi$	seen	†
$K\bar{K}^*(892) + \text{c.c.}$	seen	†
$\pi\gamma$	seen	608

**$N(1900)$  3/2<sup>+</sup>**

$I(J^P) = \frac{1}{2}(\frac{3}{2}^+)$

Breit-Wigner mass  $\approx 1900$  MeV

Breit-Wigner full width  $\sim 250$  MeV

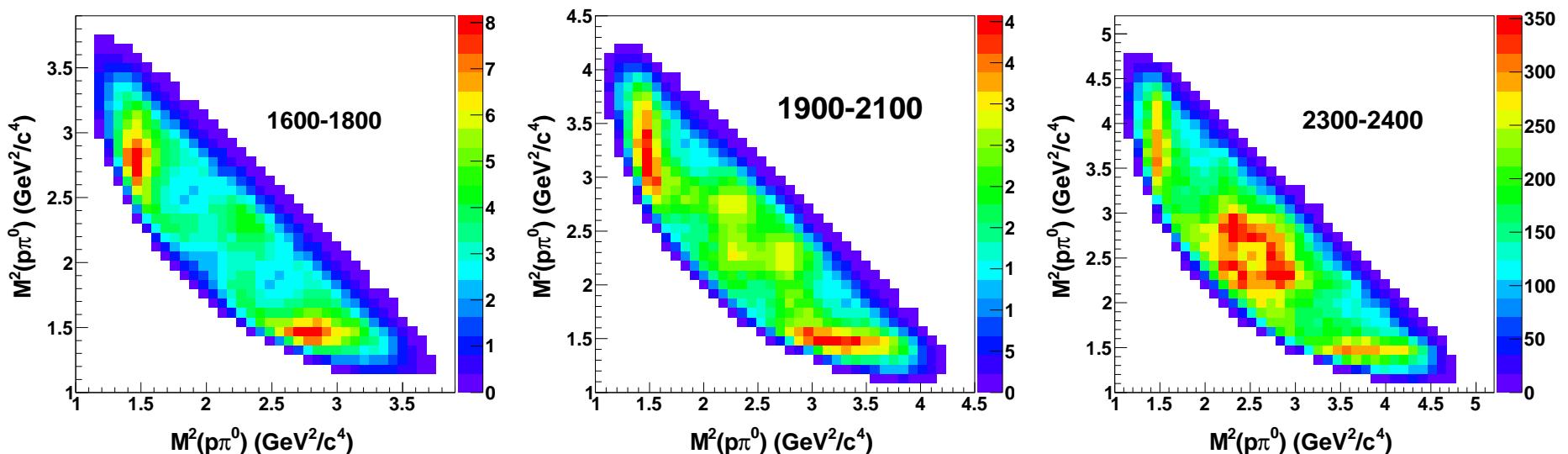
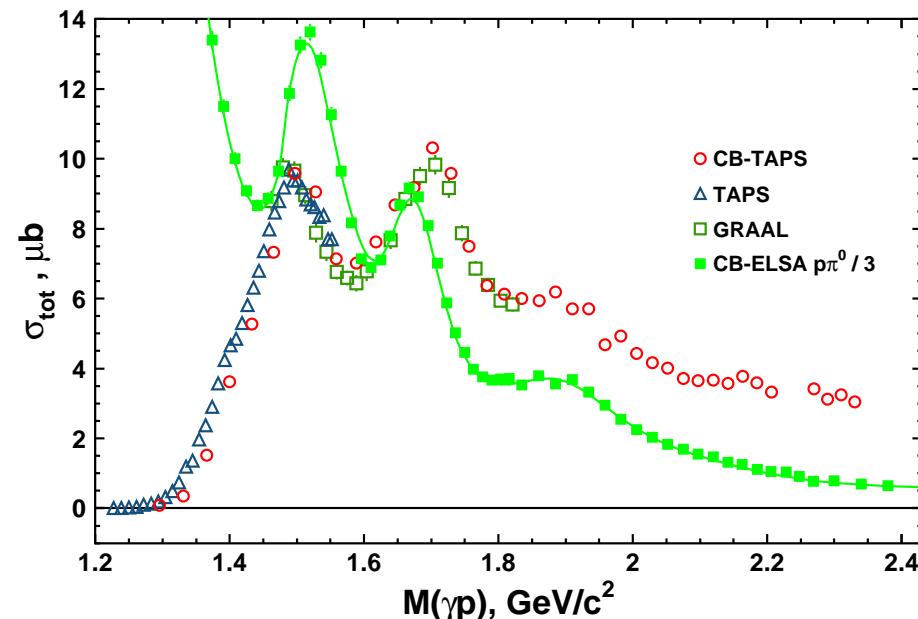
$\text{Re}(\text{pole position}) = 1900 \pm 30$  MeV

$-2\text{Im}(\text{pole position}) = 200^{+100}_{-60}$  MeV

<b><math>N(1900)</math> DECAY MODES</b>	Fraction ( $\Gamma_i/\Gamma$ )	$p$ (MeV/c)
$N\pi$	$\sim 10$ %	710
$N\eta$	$\sim 12$ %	579
$N\omega$	$(39 \pm 9)$ %	401
$\Lambda K$	0–10 %	477
$\Sigma K$	$(5.0 \pm 2.0)$ %	410

$\gamma p \rightarrow p\pi^0$  and  $\gamma p \rightarrow p2\pi^0$

- Cascades via  
 $\Delta(1232)$ ,  $N(1520)3/2^-$ ,  $N(1680)5/2^+$
- Peak position change, pole positions remain



## Three definitions of branching ratios:

$$1. \ BR_a = \Gamma_a / \Gamma_{\text{BW}}, \quad m_{\text{BW}} \Gamma_a = g_a^2 \rho_a(m_{\text{BW}}^2),$$

where the energy dependent partial is defined by

$\sqrt{s} \Gamma_a(s) = g_a^2 \rho_a(s)$ ,  $g_a$  coupling constant,  $\rho_a$  phase space.

$\sum BR_a = 1$  satisfied.

$$2. \ BR' = \int_{\text{threshold}}^{\infty} \frac{ds}{\pi} \frac{g_a^2 \rho(s)}{(m_{\text{BW}}^2 - s)^2 + (\sum_a g_a^2 \rho_a(s))^2}.$$

This definition allows for subthreshold decays. It corresponds to intuition but  $\sum BR'_a = 1$  needs to be imposed.

$$3. \ BR_{\text{pole}}(\text{channel } b) = \frac{|\text{Res}(\pi N \rightarrow b)|^2}{|\text{Res}(\pi N \rightarrow N\pi)| \cdot (\Gamma_{\text{pole}}/2)}.$$

The latter definition uses only pole related quantities.

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Parity doublets and chiral symmetry

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### 3. Baryon resonance spectrum: interpretation

#### 3.1 Parity doublets and chiral symmetry

In the limit of massless quarks:

1. Pseudoscalar mesons become massless Goldstone bosons
2. Parity partners acquire the same mass
3. Spontaneous symmetry breaking leads to hadronic masses

Vector mesons

$\approx 700 \text{ MeV}$

Nucleon mass  $M_N$

$\approx 1000 \text{ MeV}$

Its parity partner  $N(1535)1/2^-$  acquires a mass

$\approx 1500 \text{ MeV}$

$N(1650)1/2^-$	$N(1700)3/2^-$	$N(1675)5/2^-$	$\Delta(2200)7/2^-$ a.o. $\Delta(1905)7/2^+$
$N(1710)1/2^+$	$N(1720)3/2^+$	$N(1680)5/2^+$	
$N(1895)1/2^-$	$N(1880)3/2^-$	$N(2060)5/2^-$	$N(2180)7/2^-$
$N(1870)1/2^+$	$N(1890)3/2^+$	$N(2090)5/2^+$	$N(2105)7/2^+$

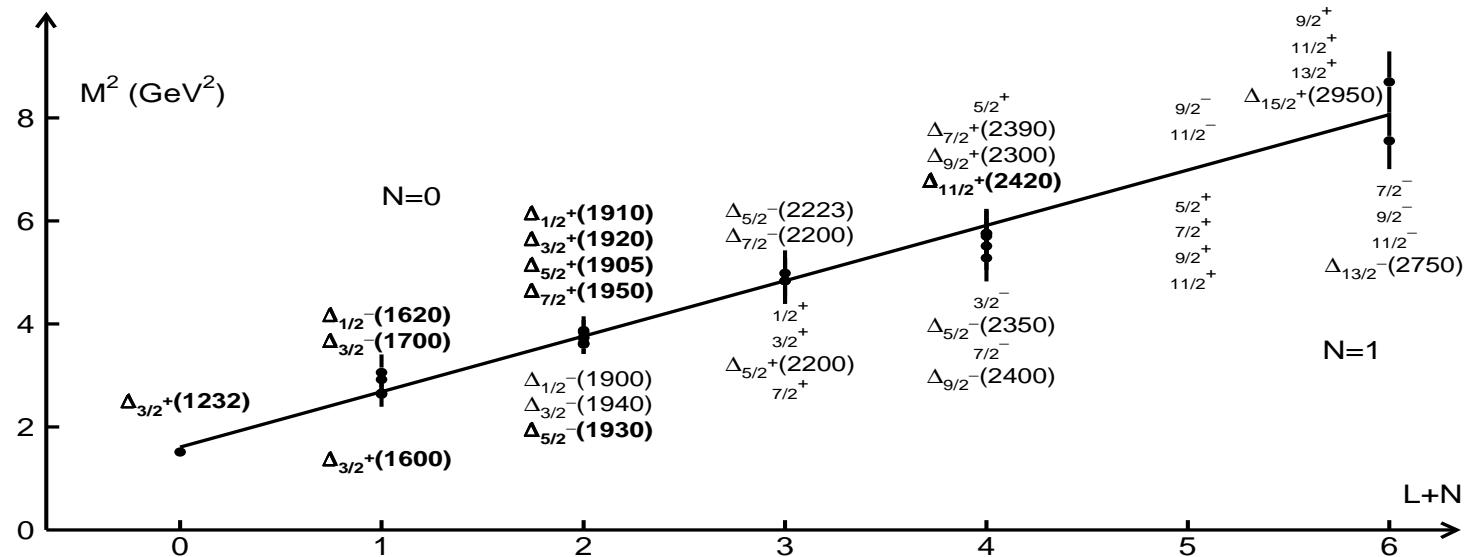
Nucleon resonances organized as parity doublets (solution BnGa2011-02).

Restoration of chiral symmetry?

L. Y. Glozman, Phys. Rept. 444, 1 (2007)

Meson and baryon resonances on the leading Regge trajectory have no parity partner!

### 3.2 AdS/QCD



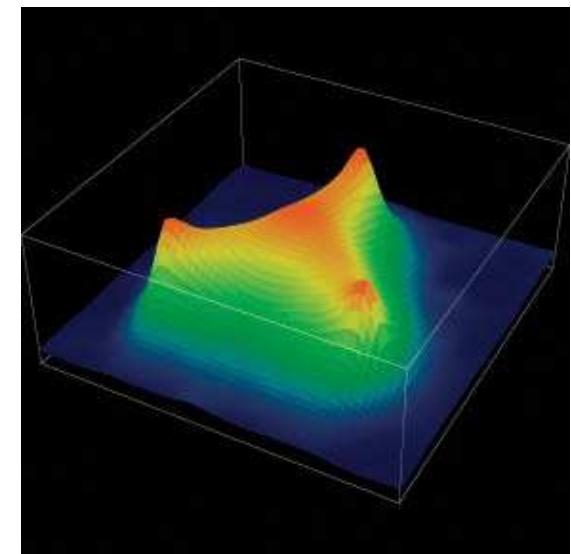
Regge trajectory of  $\Delta$  resonances:  $M^2 \propto L + N$ : non-relativistic concept !

- Mass is stored in strings which expand  $\propto L + N$
- $(l_1, l_2), (n_1, n_2) \implies (L, N)$
- Constituent quark mass is ill-defined
- AdS/QCD:

H. Forkel and E. Klempt,  
Phys. Lett. B 679, 77 (2009).

$$M^2 = a \cdot (L + N + 3/2) - b \alpha_D$$

- Two-parameter fit reproduces mass spectrum 2× better than quark models, Skyrme models (or LQCD) !



### 3.3 $SU(6) \otimes O(3)$

$J^P =$	$\frac{1}{2}^+$	$\frac{3}{2}^+$	$\frac{5}{2}^+$	$\frac{7}{2}^+$
(56, $2_2^+$ ) $S = 3/2; L = 2; N=0$ $S = 1/2; L = 2; N=0$	$\Delta_{1/2}^+(1910)$ $N_{3/2}^+(1720)$	$\Delta_{3/2}^+(1920)$ $N_{5/2}^+(1620)$	$\Delta_{5/2}^+(1905)$	$\Delta_{7/2}^+(1950)$
(70, $2_2^+$ ) $S = 1/2; L = 2; N=0$ $S = 3/2; L = 2; N=0$ $S = 1/2; L = 2; N=0$	$N_{1/2}^+(1880)$	$\Delta_{3/2}^+$ $N_{3/2}^+(1900)$ $N_{3/2}^+$	$\Delta_{5/2}^+$ $N_{5/2}^+(2000)$ $N_{5/2}^+$	$N_{7/2}^+(1990)$
<b>2<sup>nd</sup></b> (20, $1_2^+$ ) $S = 1/2; L = 1; N=0$	$N_{1/2}^+$	$N_{3/2}^+$		
(56, $0_2^+$ ) $S = 3/2; L = 0; N=1$ $S = 1/2; L = 0; N=1$		$\Delta_{3/2}^+(1600)$ $N_{1/2}^+(1440)$		
(70, $0_2^+$ ) $S = 1/2; L = 0; N=1$ $S = 3/2; L = 0; N=1$ $S = 1/2; L = 0; N=1$	$\Delta_{1/2}^+(1750)$ $N_{1/2}^+(1710)$	$\text{does not exist}$ $N_{3/2}^+$ $4^{\text{th}} \text{ shell}$		
1 <sup>st</sup> (70, $1_1^-$ ) $S = 1/2; L = 1; N=0$ $S = 3/2; L = 1; N=0$ $S = 1/2; L = 1; N=0$	$\Delta_{1/2}^-(1620)$ $N_{1/2}^-(1650)$ $N_{1/2}^-(1535)$	$\Delta_{3/2}^-(1700)$ $N_{3/2}^-(1700)$ $N_{3/2}^-(1520)$	$N_{5/2}^-(1675)$	
0 <sup>th</sup> (56, $0_0^+$ ) $S = 3/2; L = 0; N=0$ $S = 1/2; L = 0; N=0$		$\Delta_{3/2}^+(1232)$ $N_{1/2}^+(939)$		

Some multiplets are fully, others completely empty.

Why ?

- There is a spin-doublet  $\mathbf{N}_{1/2^-}$  (1880),  $\mathbf{N}_{3/2^-}$  (1890).

There is no close-by  $\mathbf{N}_{5/2^-}$

There is no close-by spin triplet.

There is  $\Delta_{1/2^-}$  (1900),  $\Delta_{3/2^-}$  (1940),  $\Delta_{5/2^-}$  (1930)

SU(6):  $56 = {}^4\mathbf{10} + {}^2\mathbf{8}$

A symmetric spin-flavor wave function with  $L = 1$  excludes  $N = 0$ .

These resonances are radial excitations belonging to the third shell

- There are four negative-parity resonances

$\mathbf{N}_{3/2^-}$  (2130),  $\mathbf{N}_{5/2^-}$  (2075),  $\mathbf{N}_{7/2^-}$  (2185),  $\mathbf{N}_{9/2^-}$  (2250).

These may form a spin quartet

There is no evidence for a ( $\mathbf{N}_{3/2^-}$ ,  $\mathbf{N}_{5/2^-}$ ) doublet

There is a  $\Delta_{7/2^-}$  (2200); likely,  $\Delta_{5/2^-}$  is missing (GWU: 2223 MeV).

SU(6):  $70 = {}^2\mathbf{10} + {}^4\mathbf{8} + {}^2\mathbf{8} + {}^2\mathbf{1}$

These resonances can be assigned to the third shell  $L = 3$

$3^{\text{rd}}$	$J^P =$	$\frac{1}{2}^-$	$\frac{3}{2}^-$	$\frac{5}{2}^-$	$\frac{7}{2}^-$	$\frac{9}{2}^-$
$(56, 1_3^-)$	$S = 3/2; L = 1; N=1$ $S = 1/2; L = 1; N=1$	$\Delta_{1/2-}(1900)$ $N_{1/2-}(1880)$	$\Delta_{3/2-}(1940)$ $N_{3/2-}(1870)$	$\Delta_{5/2-}(1930)$		
$(70, 3_3^-)$	$S = 1/2; L = 3; N=0$ $S = 3/2; L = 3; N=0$ $S = 1/2; L = 3; N=0$			$\Delta_{5/2-}(2223)$ $N_{3/2-}(2170)$	$\Delta_{7/2-}(2200)$ $N_{5/2-}(2070)$ $N_{7/2-}(2190)$	$N_{9/2-}(2250)$
		$(56, 3_3^-), (20, 3_3^-), (70, 2_3^-), (70, 1_3^-), (70, 1_3^-), (20, 1_3^-) :$ Many states predicted, no candidates known				

There are two multiplets completely filled, six multiplets completely empty.      Why?

- Minimal and maximal moments of inertia ?
- Consequence of string dynamics ?

### 3.4 dynamically generated resonances versus C.D.D. poles

$\sigma; \Lambda(1405); \dots$

$\rho; \Delta(1232); \dots$

- Meson-meson interactions:

Meson resonances can be interpreted as

$q\bar{q}$	$q\bar{q}q\bar{q}$	$q\bar{q}g$	$gg$	$m_1 m_2$
mesons	tetraquarks	hybrids	glueballs	molecules

- Meson-baryon interactions:

Baryon resonances can be interpreted as

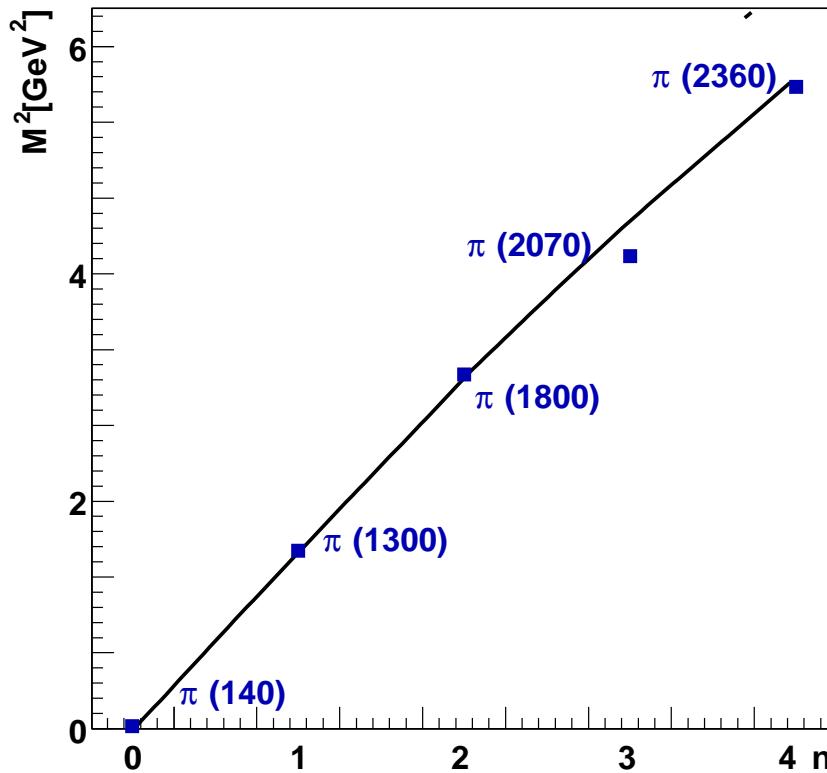
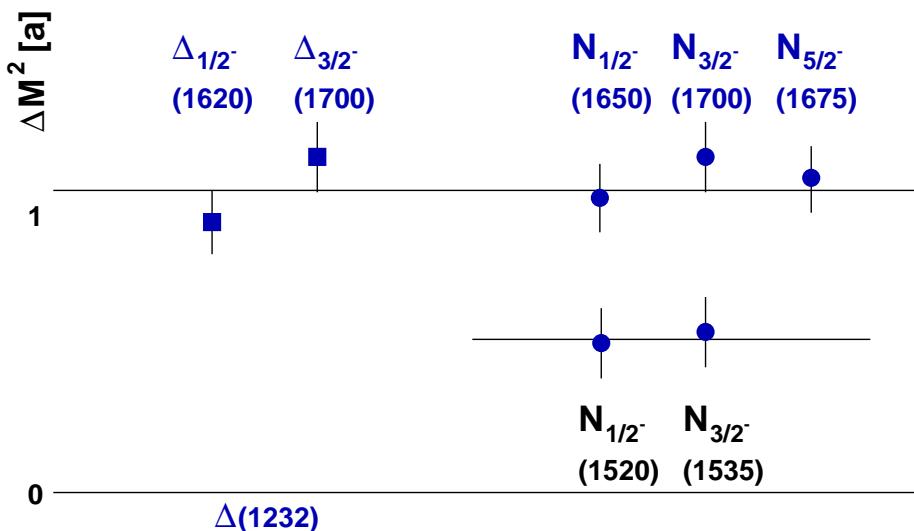
$qqq$	$qqqq\bar{q}$	$qqqg$	$b_1 m_2$
baryons	pentaquarks	hybrids	molecules

- Are all these Fock components realized and mix?

Even the quark model predicts too many baryon states!

## Two examples:

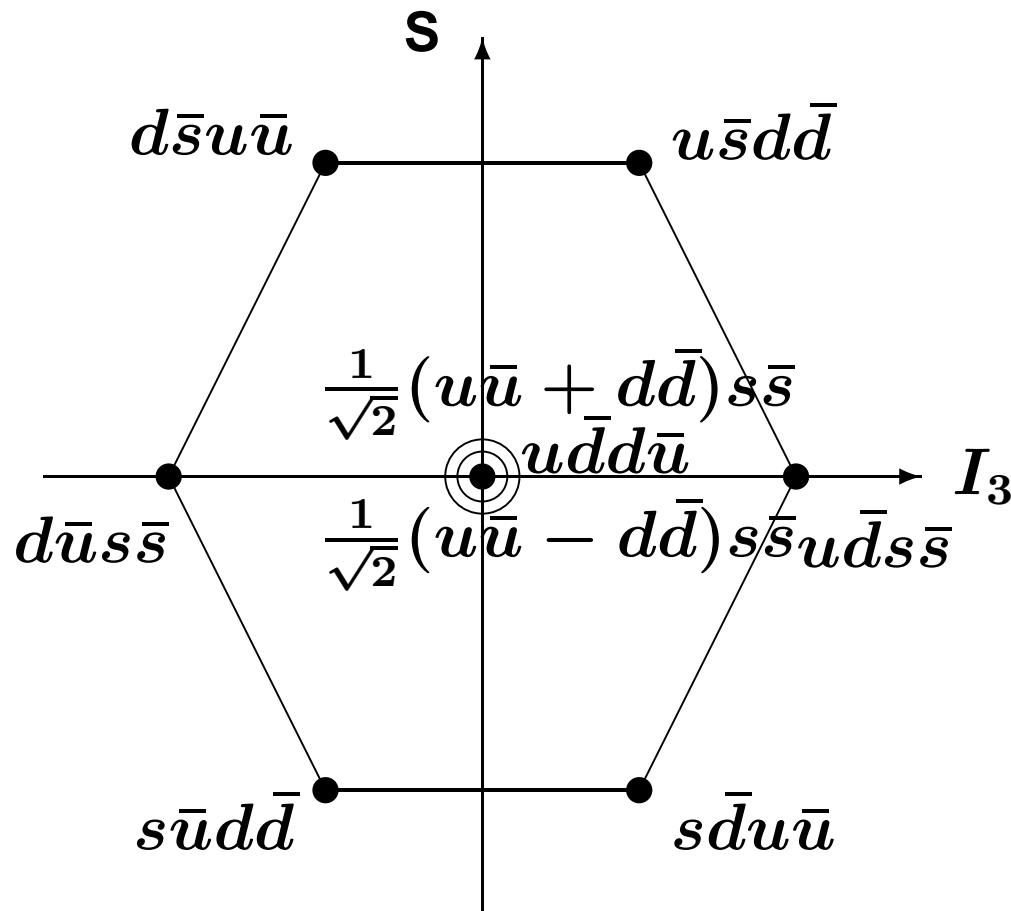
- The quark model predict five negative-parity  $N$  resonances
- $N(1535)1/2^-$  can be generated dynamically from  $N\eta - \Sigma K$  interactions



- $q\bar{q}$  meson resonances fall onto Regge trajectories
- $q\bar{q}g$  hybrids are predicted

## A third example: scalar mesons

Jaffe: Four quarks in S-wave, pairs of quarks in color  $\bar{3}$  (antiquarks in 3), energetically favored over  $q\bar{q}$  with orbital excitation.



$$\underline{f_0, a_0 = n\bar{n}s\bar{s}, \text{ 980 MeV}}$$

$$\underline{\kappa = n\bar{s}n\bar{n}, \text{ 670 MeV}}$$

$$\underline{\sigma = n\bar{n}n\bar{n}, \text{ 460 MeV}}$$

**9  $q\bar{q}$ , 9  $q\bar{q}q\bar{q}$  states  $\Rightarrow \text{SU}(4)$**

## Including *c* quarks:

There are 36 tetraquark states with quark pairs in color  $\bar{3}$ , 16 of them have a  $q\bar{q}$  component, 20 are flavor exotics. There is a subset of 10 tetraquark states with open charm having a symmetry where the two quarks can form a  $\bar{3}$  and two antiquarks a  $3$ ,

6 of them have a  $q\bar{q}$  component. 4 of them are flavor exotics.

$$\begin{array}{ll} c\bar{d}\,u\bar{u} & ; \quad c\bar{d}\,s\bar{s} \quad ; \quad c\bar{d}\,d\bar{s} \quad ; \quad c\bar{d}\,s\bar{u} \quad ; \quad c\bar{d}\,u\bar{s} \\ c\bar{u}\,d\bar{d} & ; \quad c\bar{u}\,s\bar{s} \quad ; \quad c\bar{u}\,u\bar{s} \quad ; \quad c\bar{u}\,s\bar{d} \quad ; \quad c\bar{u}\,d\bar{s} \end{array}$$

$$D_0^*(2400) = c\bar{n}(u\bar{u} + d\bar{d} + s\bar{s})/\sqrt{3}; \quad D_{s0}^*(2317)$$

Flavor exotics have never been observed. So even if the dominant configuration has a four-quark nature, the  $q\bar{q}$  component is essential to form a bound or resonant state!

## Conclusion:

The Fock expansion is realized in the form

$$\begin{aligned} |\psi(Q)\rangle = & \alpha_1 |qqq(Q)\rangle + \alpha_2 |qqqq\bar{q}(Q)\rangle \\ & + \alpha_3 |qqqg(Q)\rangle + \alpha_4 |b_1 m_2(Q)\rangle \end{aligned}$$

and there are no additional states beyond the quark model  
(perhaps exotics, perhaps chemistry of heavy-light systems)

This does not exclude that mesons like  $a_0(980)$  have a large or even dominant  $K\bar{K}$  configuration but the raison d'être is their  $q\bar{q}$  component.

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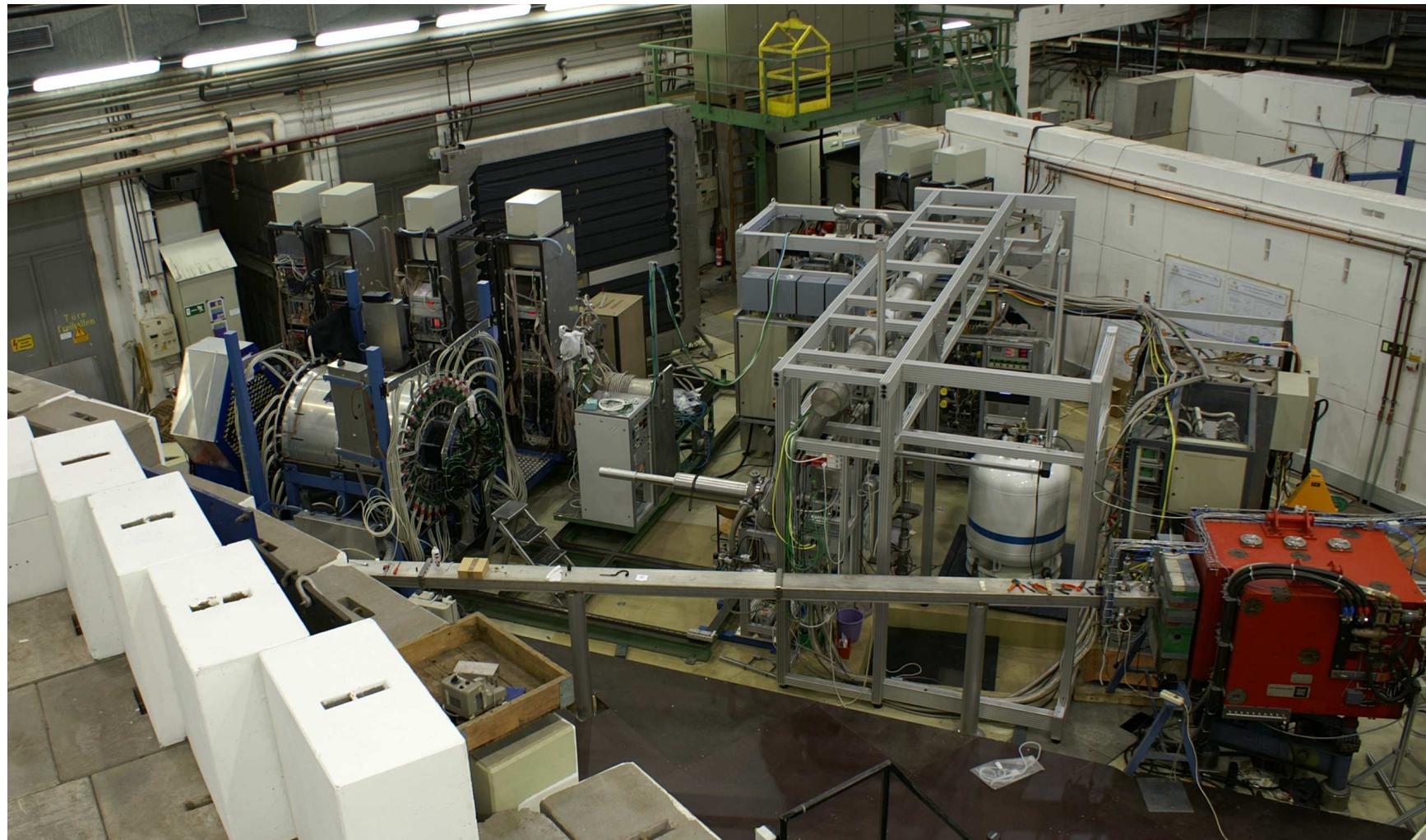
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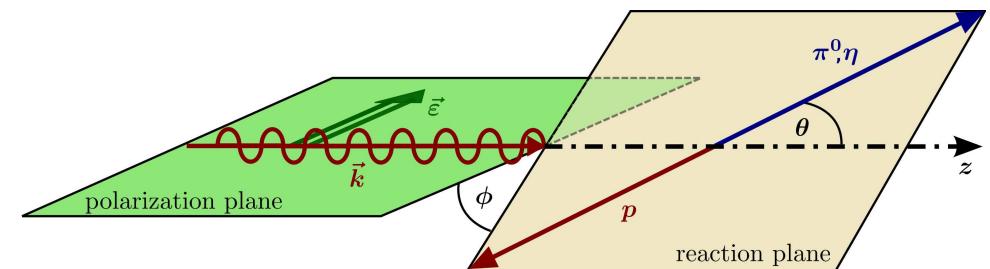
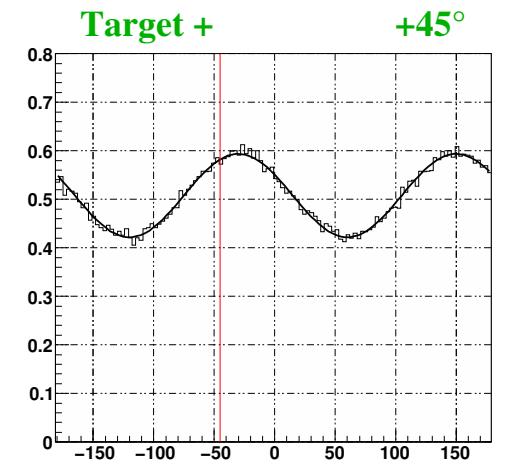
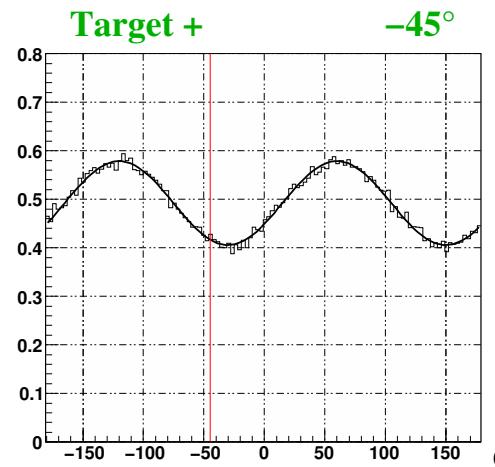
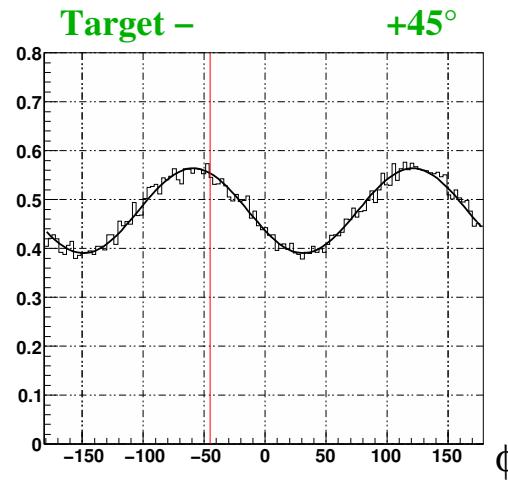
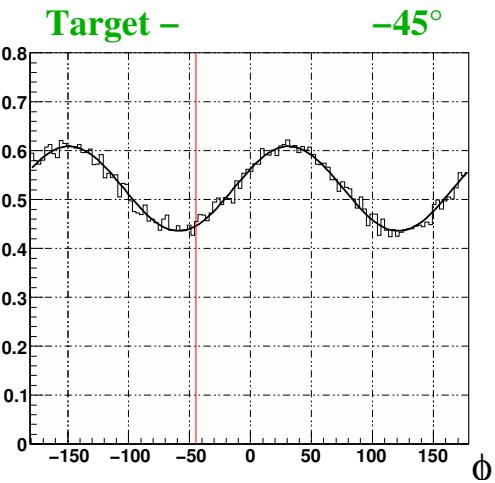
ELSA



A step forward towards a complete experiment

## Linearly polarized photons, longitudinally polarized target

$$\frac{d\sigma}{d\Omega}(\Phi) = \frac{d\sigma}{d\Omega_0} \cdot \left( 1 - P_\gamma^{\text{lin}} \Sigma \cos(2\phi) + P_\gamma^{\text{lin}} P_z G \sin(2\phi) \right)$$

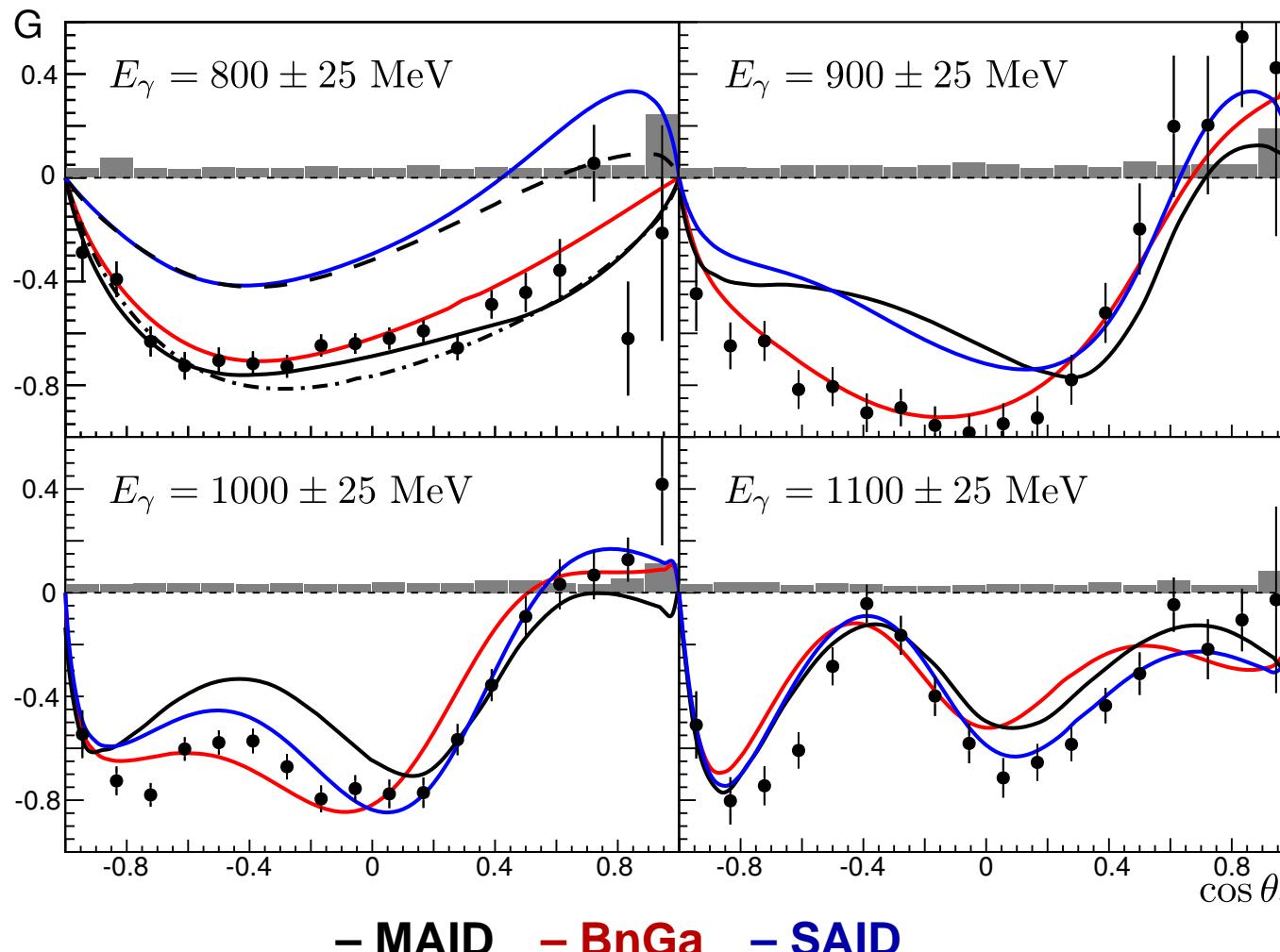


( $E_\gamma = 750-1200$  MeV)

⇒ Clear effect  
from G observed

Linearly polarized photons, longitudinally polarized target: G

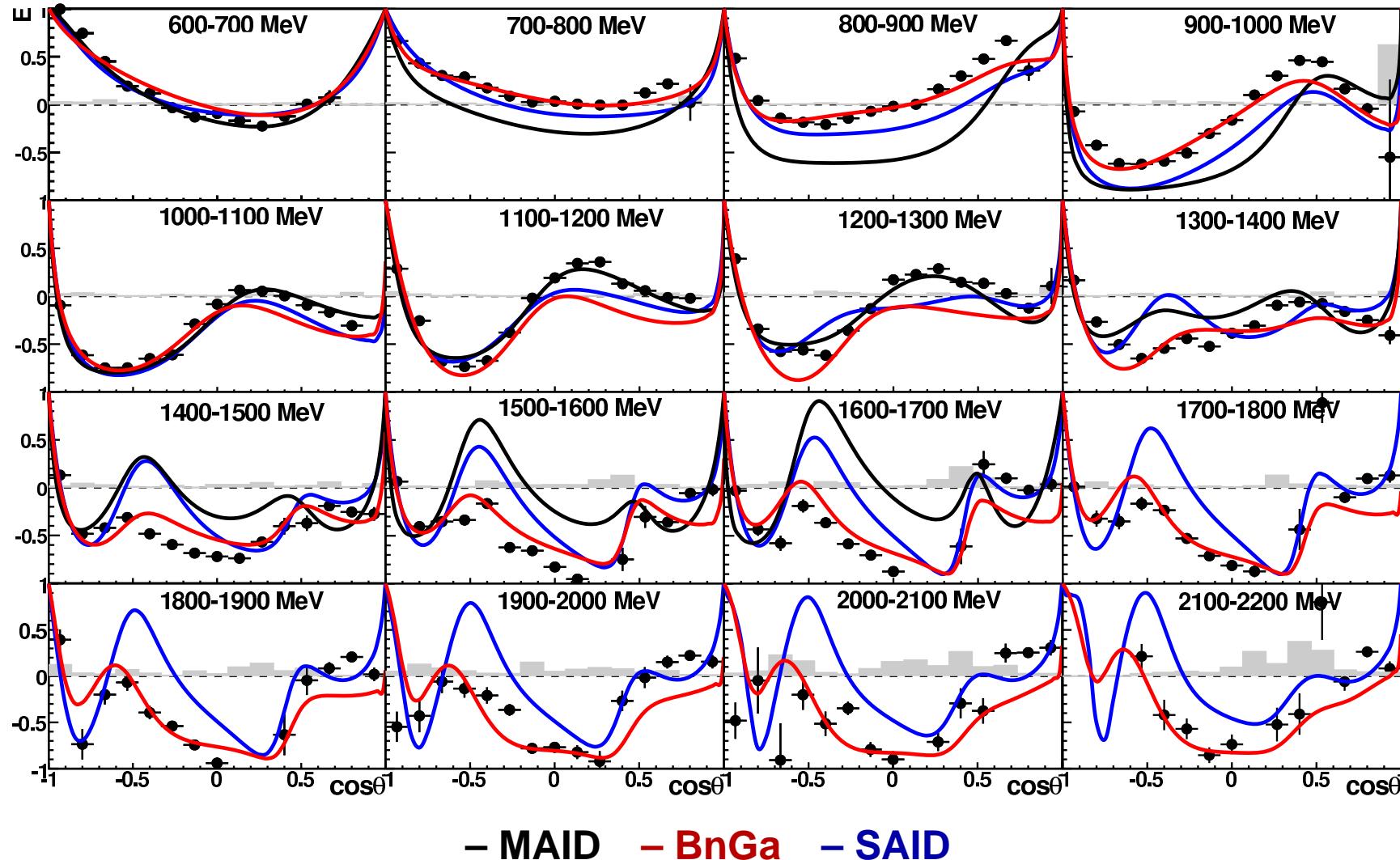
$$\frac{d\sigma}{d\Omega}(\Phi) = \frac{d\sigma}{d\Omega_0} \cdot (1 - P_\gamma^{\text{lin}} \Sigma \cos(2\phi) + P_\gamma^{\text{lin}} P_z G \sin(2\phi))$$



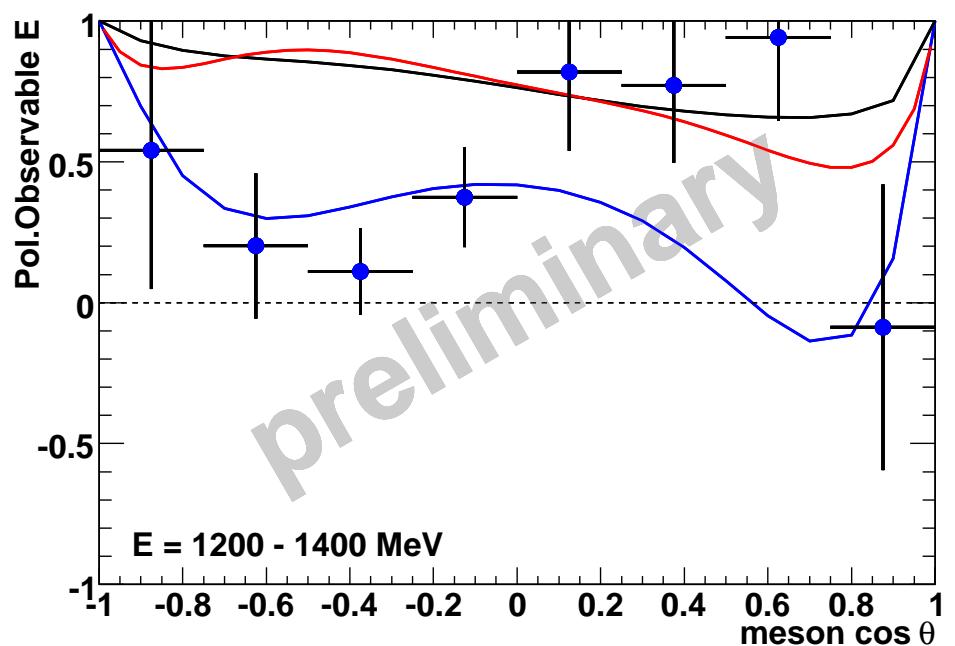
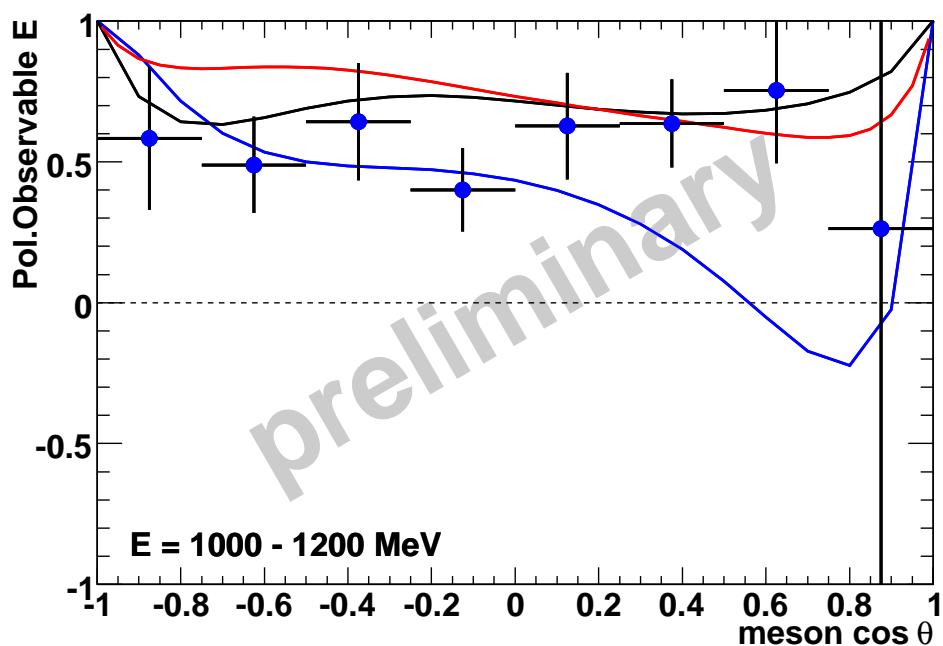
already in the second resonance region:  
significant differences !

related to the  $E_{0+}, E_{2-}$ -multipoles  
 $\leftrightarrow 1/2^-, 3/2^-$ -partial waves

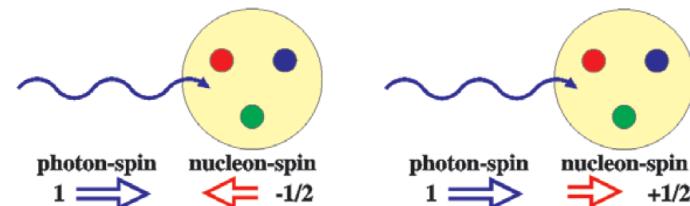
Circularly polarized photons, longitudinally polarized target: E



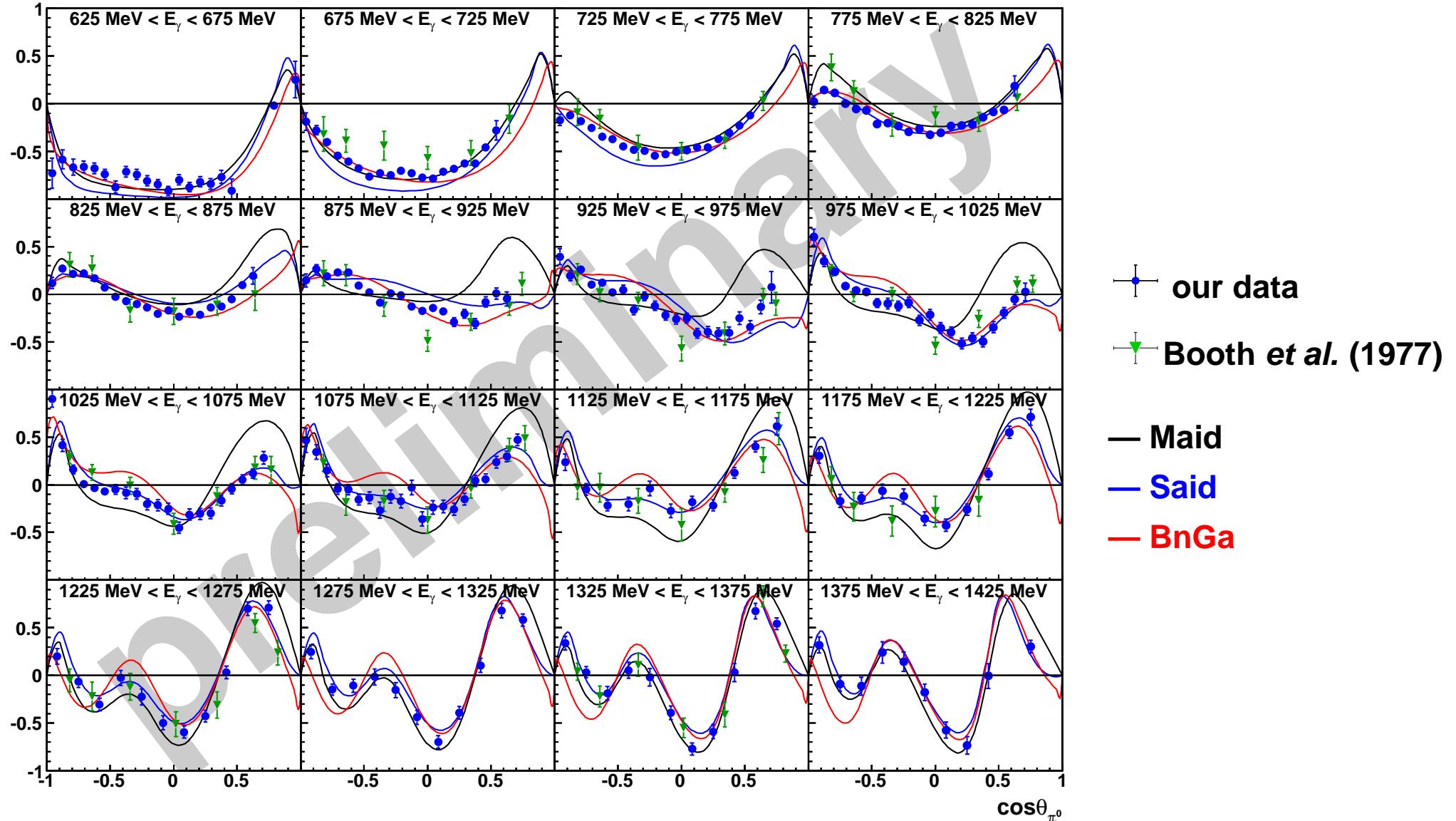
Circularly polarized photons, longitudinally polarized target: E



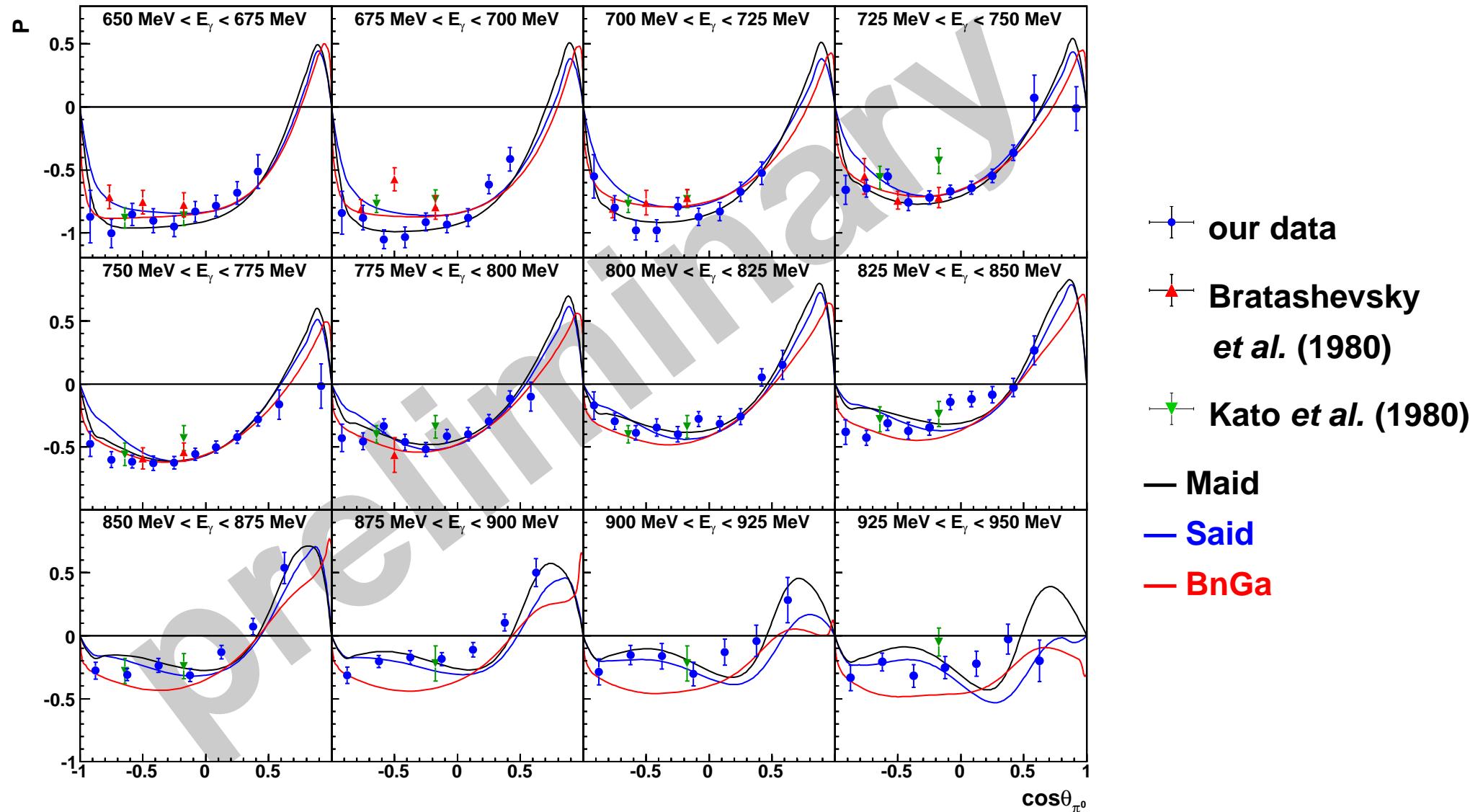
– MAID – BnGa – SAID



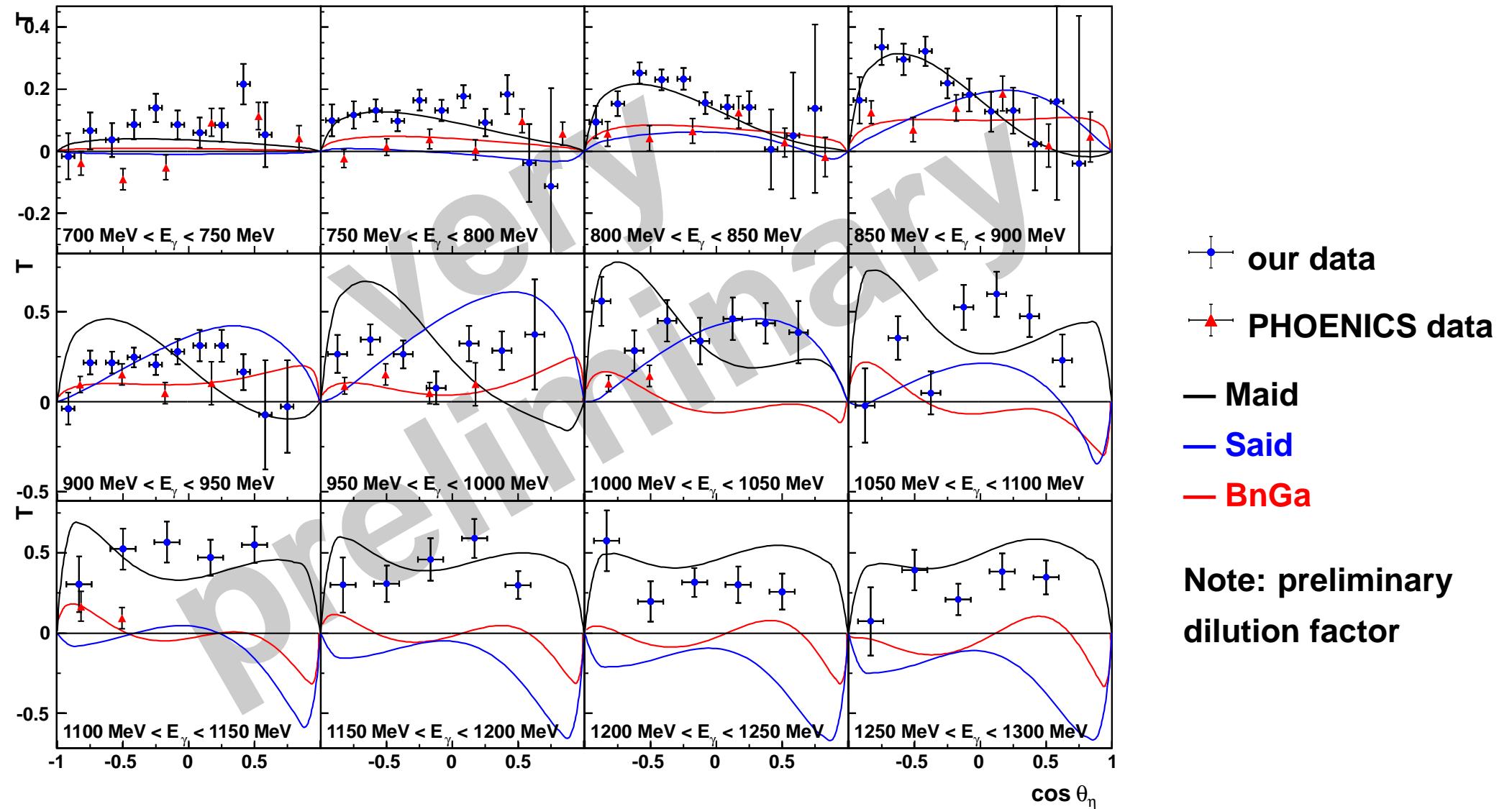
**Linear photon polarization, transversally polarized target: T**



## ⇒ Measurement of the recoil polarization: P



**Linear photon polarization, transversally polarized target: T**



# Hadron Spectroscopy: results and ideas

1. Why baryon spectroscopy

2. New results and PDG related issues

3. Baryon resonance spectrum: interpretation

Parity doublets and chiral symmetry

AdS/QCD

$SU(6) \otimes O(3)$

Dynamically generated resonances versus C.D.D. poles

4. New results from ELSA

5. Summary

## **5. Summary**

- Photoproduction experiments begin to make a significant impact on the spectrum of baryon resonances
- New data are being analyzed at Bonn, Jlab, Mainz
- Excellent prospects for the future
- The problem of missing resonances persists
- The dynamical reason why resonances are missing is unresolved:
  - Chiral symmetry restoration ?
  - Are complete multiplets missing ?
  - Constituent quarks or strings ?