(PDG) Review of Scalar Mesons Below 2 GeV

#### Simon Eidelman

#### Budker Institute of Nuclear Physics SB RAS, and Novosibirsk State University, Novosibirsk, Russia

#### Outline

- 1. The I = 1/2 states
- 2. The I = 1 states
- 3. The I = 0 states
- 4. Conclusions

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## Introduction

- Scalars are difficult to resolve because large decay widths cause a strong overlap between resonances and background
- Several decay channels can open up within a short mass range (e.g., at the  $K\bar{K}$  and  $\eta\eta$  thresholds) producing cusps
- One expects non- $q\bar{q}$  objects like glueballs and multiquark states
- Special sophisticated analyses like PWA and Dalitz plot needed
- We briefly discuss I = 1/2, I = 1 and I = 0 states since the  $(I = 2) \pi \pi$  and  $(I = 3/2) K \pi$  phase shifts do not exhibit any resonant behavior

# The I = 1/2 States – I

- The whole story began from the LASS experiment, which gave rise to numerous  $K^*$  states by studying  $K^-p \rightarrow \bar{K}^0 \pi^+ \pi^- n$  and  $K^-p \rightarrow K^- \pi^+ n$  at 11 GeV
- High-statistics measurements of *D* decays with CLEO, FOCUS and E791 under favorable conditions of small combinatorial background
- Experiments at the B factories (BaBar and Belle) have to account for  $K^*$ 's although under rather ambiguous conditions
- $\tau$  decays into  $K\pi$  and  $K\pi\pi$  are a clean laboratory without background and final state interactions

# The I = 1/2 States – II

State	$K_0^*(800)$	$K_0^*(1430)$	$K_0^*(1950)$
Mass, MeV	$682 \pm 29$	$1425 \pm 50$	$1945 \pm 10 \pm 20$
Width, MeV	$547 \pm 24$	$270\pm80$	$201 \pm 34 \pm 79$

#### Some comments:

- $K_0^*(1430)$  is a difficult state to study with  $K^*(1410)$   $(J^P = 1^-)$  and  $K_2^*(1430)$   $(J^P = 2^+)$  nearby
- $\mathcal{B}(K\pi) = (93 \pm 10)\%$  is based on a single LASS measurement: D. Aston et al., Nucl. Phys. 296 (1988) 493;  $K^-p \to K^-\pi^+n$  at 11 GeV
- $K_0^*(1950)$  is solely based on the same LASS analysis,  $\mathcal{B}(K\pi) = (52 \pm 14)\%$ , confirmed by reanalysis of A. Anisovich (1997) and Z.Y. Zhou (2006)

The 
$$I = 1/2$$
 States – III

- The  $K^*(800)$  or  $\kappa$  is still very controversial
- BES sees it in  $J/\psi \to \bar{K}^* K \pi$ , but at higher mass ~ 850 MeV
- J.R. Pelaez, Mod. Phys. Lett. A19 (2004) 2879, obtains  $M = 753 \pm 52$  MeV,  $\Gamma = 470 \pm 66$  MeV from  $K\pi \to K\pi$  and S-matrix pole
- S. Descotes-Genon, B. Moussallam, Eur. Phys. J. C 48 (2006) 553 obtain M = 753 ± 52 MeV, Γ = 557 ± 24 MeV from Kπ → Kπ, S-matrix pole and applying Roy-Steiner equations
- D. Bugg, Phys. Rev. D 81 (2010) 014002 obtains  $M = 663 \pm 35$  MeV,  $\Gamma = 658 \pm 45$  MeV accounting for s-dependent width, couplings to  $K\pi$  and  $K\eta'(958)$  and Adler zero





D. Epifanov et al. (Belle), Phys. Lett. B 648 (2007) 341

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## The I = 1 States -I

- Two isovector states are established:  $a_0(980)$  and  $a_0(1450)$
- a<sub>0</sub>(980) decays into ηπ and KK̄. The KK̄ component is large, but the resonance is just below the threshold generating a cusp. Mass and width of the a<sub>0</sub>(980) can be strongly distorted, the results are strongly model-dependent
- $a_0(1450)$  is well above thresholds with various decays reported:  $\eta \pi, K\bar{K}, \eta'(958)\pi, \omega \pi \pi, a_0(980)\pi \pi$

#### The I = 1 States – II

State	$a_0(980)$	$a_0(1450)$
Mass, MeV	$980\pm20$	$1474 \pm 19$
Width, MeV	50 - 100	$265 \pm 13$

- Large spread of values for the  $a_0(980)$ , strong model dependence; data for the  $K\bar{K}$  are scarce, might be different from the  $\eta\pi$ ?
- The a<sub>0</sub>(1450) parameters are based on Crystal Barrel pp̄ data at rest, other analyses: a big spread from 1290 to 1565 MeV for mass, 80 to 270 MeV for width

# The I = 0 States – I (General)

- The I = 0,  $J^{PC} = 0^{++}$  sector is the most complex
- Data exist in the  $\pi\pi$ ,  $K\bar{K}$ ,  $\eta\eta$ ,  $4\pi$ ,  $\omega\omega$ , and  $\eta\eta'(958)$  in S-wave
- Various production processes:  $\pi N$ ,  $p\bar{p}$ ,  $J/\psi$ , D, B decays
- Five poles are needed:  $f_0(500)$  or  $\sigma$ ,  $f_0(980)$ ,  $f_0(1370)$ ,  $f_0(1500)$ , and  $f_0(1710)$

## The I = 0 States – II $(\sigma)$

- Long and controversial history
- $\pi\pi$  S-wave phase shifts  $\delta_J^I = \delta_0^0$ : old information from  $\pi N$  scattering and new data near threshold from  $K_{e4}$  decay
- An enhancement in  $M(\pi^+\pi^-)$  is observed in  $D^+ \to \pi^+\pi^-\pi^+, J/\psi \to \omega\pi^+\pi^- \text{ and } \psi(2S) \to J/\psi\pi^+\pi^-$
- The  $\pi\pi$  scattering amplitude has unusual energy dependence due to crossed channel exchanges and  $f_0(1370)$
- Analyses of  $\pi\pi$  data that include unitarity, threshold behavior, constraints from Adler zeroes, scattering lengths and  $K_{e4}$  data find a light scalar with a pole near (500 - *i*250) MeV

### The I = 0 States – III ( $\sigma$ )

• Precise pole positions with uncertainties below 20 MeV were extracted using Roy equations and minimal data by the Bern group, B. Moussalam and R. Garcia-Martin et al.

• 
$$\sqrt{s_{\text{Pole}}^{\sigma}} = (400 - 550) - i(200 - 350) \text{ MeV}$$

- A more radical approach is to average the advanced dispersion analyses:  $\sqrt{s_{\text{Pole}}^{\sigma}} = (446 \pm 6) - i(276 \pm 5) \text{ MeV}$
- The  $\gamma\gamma$  data are inconclusive due to its large width
- There are theoretical indications that the  $f_0(500)$  is not a  $q\bar{q}$





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## The I = 0 States – V $(f_0(980))$

- Seen very clearly as a narrow signal in many processes
- Overlaps strongly with the background from  $f_0(500)$  and  $f_0(1370)$ , that can result in a dip in the  $M(\pi\pi)$  spectrum at the  $K\bar{K}$  threshold
- A dip changes to the  $M(\pi^0\pi^0)$  peak in  $\pi^- p \to \pi^0\pi^0 n$
- Seen in  $\phi \to \pi^+ \pi^- \gamma$ ,  $\pi^0 \pi^0 \gamma$  by KLOE, SND and CMD-2
- Also seen in various charmonium and B meson decays
- A recent dispersive analysis of B. Moussallam and R. Garcia-Martin et al.!
- $M(f_0(980)) = 990 \pm 20 \text{ MeV}$

## The I = 0 States – VI ( $f_0(980)$ mass)



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#### The I = 0 States – VII (Above 1 GeV)

- $f_0(1500)$  is well established and seen by Crystal Barrel in  $p\bar{p}$  in five decay modes:  $\pi\pi$ ,  $K\bar{K}$ ,  $\eta\eta$ ,  $\eta\eta'(958)$ ,  $4\pi$ ; also seen in central production
- $p\bar{p}$ ,  $p\bar{n}/n\bar{p}$  show a single enhancement at 1400 MeV in  $M(4\pi)$  resolved in  $f_0(1370)$ and  $f_0(1500)$ , the  $4\pi$  mode is about half of the  $f_0(1500)$  and dominant for the  $f_0(1370)$ , other branchings rather uncertain
- Both Belle and BaBar observe scalars in B and D decays reporting broad or narrow structures between 1 and 1.6 GeV in  $\pi^+\pi^-$  and  $K^+K^-$
- $f_0(1710)$  seen in  $J/\psi$  radiative decays and central production mostly decays into  $K\bar{K}$  and  $\eta\eta$ ;  $\pi\pi$  and  $\omega\omega$  possibly seen
- Belle studied  $\gamma \gamma \to \pi^0 \pi^0$ ,  $(\eta \eta)$  and found an *S*-wave state with mass 1470 (1262) MeV and width 90 (484) MeV and huge errors

### The I = 0 States – VIII (Summary)

State	$\sqrt{s_{\text{pole}}},  \text{MeV}$	$M_{\rm BW},  {\rm MeV}$	$\Gamma_{\rm BW},  {\rm MeV}$
$f_0(500)$	(400-550)-i(200-350)	400-1200	600-1000
$f_0(980)$	_	$990\pm20$	40-100
$f_0(1370)$	(1200-1500)-i(150-250)	1200-1500	200-500
$f_0(1500)$	_	$1505\pm6$	$109\pm7$
$f_0(1710)$	_	$1720\pm 6$	$135\pm8$

The  $<<\,$  BW >> values are a mixture of Breit-Wigner and S-matrix poles Averaging hardly makes sense!

#### Interpretation of the I = 0 States -I (Below 1 GeV)

- $f_0(980)$  and  $a_0(980)$  are 4-quark states or molecules:  $\gamma\gamma$  is inconclusive,  $f_0(980)$  has a large  $s\bar{s}$  from  $J/\psi$  and  $D_s$  decays,  $\phi \to f_0(980)\gamma$  and  $\phi \to a_0(980)\gamma$  show a role of kaon loops
- $f_0(980), a_0(980), f_0(500)$  and  $K_0^*(800)$  form a low-mass scalar nonet of predominantly four-quark states
- Models with chiral Lagrangians predict  $f_0(500)$  and a nonet with  $f_0(500), f_0(980), a_0(980)$  and  $K_0^*(800)$  (or  $K_0^*(1430)$ )
- The unitarized quark model with coupled q\$\overline{q}\$ and meson-meson channels, light scalars are bare q\$\overline{q}\$ confinement states strongly shifted from 1.3-1.5 GeV due to strong coupling to S-wave two-meson decays.
  \$f\_0(500)\$, \$f\_0(980)\$, \$K\_0^\*(800)\$ and \$a\_0(980)\$ as well as \$f\_0(1370)\$, \$f\_0(1500)\$ (or \$f\_0(1710)\$)\$, \$K\_0^\*(1430)\$ and \$a\_0(1450)\$ form two nonets manifestations of the same bare input states

#### Interpretation of the I = 0 States – II (Above 1 GeV)

- The  $f_0(1370)$  and  $f_0(1500)$  decay mostly into  $2\pi$  and  $4\pi$ , while  $f_0(1710)$  into  $K\bar{K}$
- In the naive quark model f<sub>0</sub>(1370), a<sub>0</sub>(1450) and K<sub>0</sub><sup>\*</sup>(1430) the same SU(3) flavor nonet as (uū + dd), ud and us states, mixing with the light scalars (?), the f<sub>0</sub>(1710) is an ss state: f<sub>0</sub>(1710) (f<sub>2</sub>'(1525)) are produced in pp, but suppressed compared to f<sub>0</sub>(1500) (f<sub>2</sub>(1270)), and f<sub>0</sub>(1500) is also a (uū + dd), but very narrow and too light
- But in  $\gamma\gamma$  the  $f_0(1500)$  is not observed while the  $f_0(1710)$  is, so it is mainly glue (?). The  $f_0(1500)$ 's small width and enhanced production at small  $p_t$  in central collisions favor its non- $q\bar{q}$ . In a mixing scheme of F. Close glue is shared between  $f_0(1370)$ ,  $f_0(1500)$  and  $f_0(1710)$  (mainly  $n\bar{n}$ , glue and  $s\bar{s}$ )
- In other schemes two-gluon coupling to  $n\bar{n}$  is suppressed by chiral symmetry and  $K\bar{K}$  is enhanced  $\Rightarrow f_0(1710)$  is an unmixed glueball
- Broad  $KK\bar{K}$  signal of Belle as a  $f_0(1500)$  glueball, in conflict with BaBar
- $f_0(1370)$  and  $f_0(1710)$  are bound states of two vectors

### Conclusions

- $f_0(500)(\sigma)$  definitely and  $K_0^*(800)(\kappa)$  most probably exist
- Many other scalars observed, probably too many for normal  $q\bar{q}$
- $\gamma\gamma$  collisions can help (?) with much larger data samples, KLOE-2 with taggers, CMD-3/SND, BES-III, Super*B*-factories
- Many interesting observations from the J/ψ and ψ(2S) decays (BES), results from Belle and BaBar suffer from low data hadronic samples, but their τ decay data are very rich and potentially helpful
- Theoretical interpretation is very far from final and new experimental observations badly needed
- Strong model dependence, variety of approaches
- A nightmare for experiment and paradise for theory!

