for Non-Strange Baryons

New Neutron Electromagnetic Couplings

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- SAID for Baryon Spectroscopy.
- Pion photoproduction off the neutron.
- **FSI** for $\gamma n \rightarrow \pi^- p$.
- **R**ecent **GW**–**ITEP** FSI for $\gamma n \rightarrow \pi^- p$.
- New neutron EM couplings.
- Polarized Measurements.











PWA for non-strange Baryons & <u>SAID</u> Database

0 de fi m	riginally: PWA arose as the etermine amplitude of the re tting scattering data which i athematical problem	technology to action via 's a non-trivial	CINS	_		
[2	<i>Solution of ill-posed problem</i> – Hadamard,	Tikhonov, <i>et al</i> j	Data Analysis Center	<u>Be</u>	elow 4	<u>GeV</u>
R	esonances appeared as a by-	product	Partial-Wave Analyses at GW			
[E	Bound states objects with de	finite quantum	[See Instructions]			
r	numbers, mass, lifetime, <i>etc</i>	31,402	Pion-Nucleon	[\	/ = 1320 to :	1930 MeV]
			Pion-Pion-Nucleon		241,214	evts
	That is the strategy of	the 5,267 -	Kaon-Nucleon			J
	GW/VPI πN PWA since 1	987	Nucleon-Nucleon		38,162	
		27,265	Pion Photoproduction			_
			Pion Electroproduction		113,900	1
		9,086	Kaon Photoproduction			
			Eta Photoproduction		6,235	
		1,030	Eta-Prime Photoproduction			
			Pion-Deuteron (elastic)		1,914	
		6,083	Pion-Deuteron to Proton+Proto	n		
			For $\pi \rightarrow 2\pi$, we use log-likelihood while for the rest – least-squares technologies			
AN ANA	9/22/2013	PWA7 Workshop, Camogli	, Italy, Sept. 2013	lgor Str	akovskv	3
				5	/	

\mathcal{N}^* and Δ^*S tates coupled to $\pi \mathcal{N}$ [SAID: http://gwdac.phys.gwu.edu/]



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Status of Non-strange Resonances

GW SAID Contribution

l=1/2		Status			s	tatus	as s	een ii	n —	 Most of States field fille work to do: Most of OCD models predict more states to 			c + b) <u></u>	bcorvod						
									QUL		ueis	predict		e 5	late	5 U		Juserveu.			
Particle J^P	overa	ll π.N	γN	$N\eta$	Nσ	$N\omega$	ΛK	ΣK	Νρ	$\Delta \pi$	• Where a	are n	nissir	ng re	sonance	es?					
N 1/2 ⁺	****										1-2/2				s	tatus	as se	en in	_		=
$N(1440) 1/2^+$	****	****	****		***				*	***	1-5/2					tatus	asse	LII III			l
$N(1520) 3/2^{-}$	****	****	****	***					***	***			Statu	s							
$N(1535) 1/2^{-1}$	****	****	****	****					**	*	Particle J^{P}	overa	$11 \pi N$	γN	Νη Νσ	$N\omega$	ΛK	ΣK	Νρ	$\Delta \pi$	L
$N(1650) 1/2^{-}$	****	****	***	***			***	**	**	***	$\Delta(1232) 3/2^+$	****	****	****	F						-
$N(1675) 5/2^{-1}$	****	****	***	*			*		*	***	$\Delta(1600) 3/2^+$	***	***	***	0				*	***	
$N(1680) 5/2^+$	****	****	****	*	**				***	***	$\Delta(1620) 1/2^{-1}$	****	****	***	r				***	***	
N(1685) ? [?]	*										$\Delta(1700) 3/2^{-1}$	****	****	****	ь				**	***	
$N(1700) 3/2^{-1}$	***	***	**	*			*	*	*	***	$\Delta(1750) 1/2^+$;						
$N(1710) 1/2^+$	***	***	***	***		**	***	**	*	**	$\Delta(1000) 1/2^{-1}$	- 	**	**		d		**	**	**	
$N(1720) 3/2^+$	****	****	***	***			**	**	**	*	$\Delta(1900) 1/2$ $\Delta(1905) 5/2^+$		**	**		u a		**		**	
$N(1860) 5/2^+$	**	**							*	*	$\Delta(1903) 3/2$ $\Delta(1010) 1/2^{+}$	****	****	****		u		***		**	
$N(1875) 3/2^{-1}$	***	*	***			**	***	**		***	$\Delta(1910) 1/2^{+}$	****	****	**		e		*	*	**	
$N(1880) 1/2^+$	**	*	*		**		*				$\Delta(1920) 3/2^{-1}$	***	***	**			n	***		**	
$N(1895) 1/2^{-1}$	**	*	**	**			**	*			$\Delta(1930) 5/2$	***	***		-						
$N(1900) 3/2^+$	***	**	***	**		**	***	**	*	**	$\Delta(1940) 3/2$	**	*	**	F			(see	n in	$\Delta \eta$)	
$N(1990) 7/2^+$	**	**	**					*			$\Delta(1950) 7/2^+$	****	****	****	0			***	*	***	
$N(2000) 5/2^+$	**	*	**	**			**	*	**		$\Delta(2000) 5/2^+$	**			r					**	
$N(2040) 3/2^+$	*										$\Delta(2150) 1/2^{-1}$	*	*		b						L
$N(2060) 5/2^{-1}$	**	**	**	*				**			$\Delta(2200) 7/2^{-}$	SAI	D:								
$N(2100) 1/2^+$	*										$\Delta(2300) 9/2^+$	To	nds (b	w con	struction)	to mi	cc n	arro	•• NI*	c with	
$N(2150) 3/2^{-}$	**	**	**				**			**	$\Delta(2350) 5/2^{-1}$				vide N*e I		33 III			5 WILL	
$N(2190) 7/2^{-}$	****	****	***			*	**		*		$\Delta(2390) 7/2^+$	• Re	veals	only v	vide iv S, I	Jut no	οι το	W U	ue		
$N(2220) 9/2^+$	****	****									$\Delta(2400) 9/2^{-1}$				and po	ssess	ing i	not t	00 SI	mall B	K [BK > 0.04]
$N(2250) 9/2^{-}$	****	****									$\Delta(2420) 11/2^+$	****	****	*							
$N(2600) 11/2^{-}$	***	***									$\Delta(2750) 13/2^{-1}$	**	**								
N(2700) 13/2 ⁺	**	**									$\Delta(2950) 15/2^+$	**	**								

• More than half of states have poor evidence.

k to do.





Where We Are Now

Some of the N* baryons [N(1675)D₁₅, for instance] have stronger EM coupling to the neutron than to the proton but parameters are very uncertain.

$$N(1675) 5/2^{-1} I(J^P) = \frac{1}{2}(\frac{5}{2})$$
 Status: ****

• **PDG** estimate for the $A_{1/2} \& A_{3/2}$ decay amplitudes of the N(1720)P₁₃ state are consistent with *zero*, while the recent SAID determination gives small but non-vanishing values.

$$N(1720) \ 3/2^{+} I(J^{P}) = \frac{1}{2}(\frac{3}{2}^{+}) \ \text{Status:} \ * * * \mathbf{P_{13}} \qquad N(1720) \ 3/2^{+} \rightarrow p\gamma \ , \ \text{helicity-1/2 ampl, A1/2. -0.01 to 0.11} \ + 0.095 \pm 0.002 \ N(1720) \ 3/2^{+} \rightarrow p\gamma \ , \ \text{helicity-3/2 ampl, A3/8: -0.019 \pm 0.029} \ - 0.048 \pm 0.002 \ - 0.048 \pm 0.00$$

• Other unresolved issues relate to the second P_{11} , $N(1710)P_{11}$, that are not seen in the recent $\pi N PWA$, contrary to other PWAs used by the PDG12.

$$N(1710) \ 1/2^{+} \ I(J^{P}) = \frac{1}{2}(\frac{1}{2}^{+}) \ \text{Status:} \ * * * P_{11} \ The latest GWU analysis (ARNDT 06) finds no evidence for this resonance.$$



SAID for $\pi \mathcal{N} \rightarrow \pi \mathcal{N} \ll \pi^- p \rightarrow \eta n$

[R. Arndt, W. Briscoe, IS, R. Workman, Phys Rev C 74, 045205 (2006)]

- Energy dependent SP06/WI08 and associated SES
 T = 0 2600 MeV [W = 1078 2460 MeV]
 4-channel Chew-Mandelstam K-matrix parameterization [πN, πΔ, ρN, ηN]
 3 mapping variables: g²/4π, a[π-p], Eth
 PWs = 30 πN {15 [I=1/2] + 15 [I=3/2]} + 4 ηN [I < 9]
 Prms = 99 [I=1/2] + 89 [I=3/2]
- <u>1st generation</u> ('57–'79)
 Used by CMB79 and KH84 analyses:
 <u>10k</u> π[±]p each & <u>1.5k</u> CXS.
 17% data is polarized.
- <u>2nd generation</u> ('80–'06) SAID fits:
 13k π[±]p each, 3k CXS & 0.3k π⁻p→ηn 25% data is polarized.
 Meson Factories [LAMPF, TRIUMF, & PSI] are the main source of new measurements.
 There is no discrimination against data
- <u>3rd generation</u> (07'+) New data may come from J-PARC, HADES, EPECUR, *etc*



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Reaction	Data	χ ²
π⁺p→π⁺p	13,354	27,136
π⁻p→ π⁻p	11,978	22,632
π⁻ p→ π ⁰ n	3,115	6,068
π- p→ η n	257	650
DR constraint	2,775	671
Total	31,479	57,157

DRs have been derived from the *first principles*.



27 σ^{tot} & **37** P data

[0 – 2600 MeV] → 10 data/MeV

[550 – 800 MeV] → 1 data/MeV

above 800 MeV \rightarrow 0.03 data/MeV

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Single **Pion** Photoproduction

- An accurate evaluation of the EM couplings $N^*(\Delta^*) \rightarrow \gamma N$ from meson photoproduction data remains a paramount task in hadron physics.
- Only with good data on both the proton and neutron targets, one can hope to disentangle the isoscalar & isovector EM couplings of various N*& Δ* resonances,
 K.M. Watson, Phys Rev 95, 228 (1954); R.L. Walker, Phys Rev 182, 1729 (1969) as well as the isospin properties of non-resonant background amplitudes.



 The lack of the γn→π⁻p & γn→π⁰n data does not allow us to be as confident about the determination of neutron couplings relative to those of the • proton.

• The radiative decay width of neutral baryons may be extracted from $\pi^- \& \pi^0$ photoproduction off the neutron, which involves a bound neutron target and needs the use of model-dependent nuclear (FSI) corrections.

A.B. Migdal, JETP 1, 2 (1955); K.M. Watson, Phys Rev 95, 228 (1954)







SAID for Pion Photoproduction

[W. Chen et al, Phys Rev C 86, 015206 (2012)]



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Recent SAID Progress in PionPR

DU13:	included recent CLAS π^{0} p & π^{+} n Σ	[M. Dugger et al, arXiv:1308.4028 [nucl-ex]	١.
GB12/GZ12:	included recent CLAS π⁻p d o /dΩ	[W. Chen <i>et al,</i> Phys Rev C 86 , 015206 (2012)]	
CM12:	CM parameterization for $T_{\pi N}$	[R. Workman et al, Phys Rev C 86, 015202 (2012)]	Γ
SN11/SK11:	included recent GRAAL π⁻p & π ⁰ n Σ	Ξ	
	<i>LEPS</i> π ⁰ p dσ/dΩ	[R. Workman <i>et al,</i> Phys Rev C 85 , 025201 (2012)]	
M	$= (Born + A)(1 + iT_{\pi N}) + BT_{\pi N} + (C + iC)$	$D)(ImT_{\pi N} - T_{\pi N} ^2)$	
SP09:	included recent CLAS π^+ n d σ /d Ω	[M. Dugger <i>et al,</i> Phys Rev C 79 , 065206 (2009)]	
M	= (Born + α_R)(1 + iT $_{\pi N}$) + $\alpha_R T_{\pi N}$ + high	ier terms	/

Solution	Energy Limit (MeV)	χ^2/N_{Data}	N _{Data}
DU13	2700	2.23	27,265
GB12	2700	2.09	26,179
CM12	2700	2.01	25,814
SN11	2700	2.08	25,553
SP09	2700	2.05	24,912
FA06	3000	2.18	25,524
SM02	2000	2.01	17,571
SM95	2000	2.37	13,415

- The overall SAID χ² has remained stable against the growing database, which has increased by a factor of 2 since 1995.
- Most of this increase coming from photon-tagging facilities.







World Neutral and Charged PionPR Data

[SAID: http://gwdac.phys.gwu.edu/]

Future exp activity will fill empty spots specifically for n-target.



 $\gamma \pi \rightarrow \pi \rho Experiment$

- The existing $\gamma n \rightarrow \pi^- p$ database contains mainly differential cross sections (17%) of which are from polarized measurements.)
- Many of these are old **bremsstrahlung** measurements with limited **angular** (θ = 40 - 140°) coverage and large **energy** binning (E_{γ} = 100 - 200 MeV.) In several cases, the **systematic** uncertainties have not been given.

- At lower energies (E_{γ} < 700 MeV,) there are data sets for the inverse π^- photoproduction reaction: $\pi^-p \rightarrow \gamma n$. This process is free from complications associated with a deuteron target.
- However, the disadvantage of using $\pi^- p \rightarrow \gamma n$ is the large background because of the **5** to **500** times larger cross section for $\pi^- p \rightarrow \pi^0 n \rightarrow \gamma \gamma n$.

FSI and $\gamma d \rightarrow \pi \bar{\rho} \rho \longrightarrow \gamma \pi \rightarrow \pi \bar{\rho}$ [V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, IS, Phys Rev C 84, 035203 (2011)]

• FSI plays a critical role in the state-of-the-art analysis of $\gamma n \rightarrow \pi N$ data. • For $\gamma n \rightarrow \pi^- p$ the effect: 5% - 60%. It depends on (E, θ).

$\gamma d \rightarrow \pi^- pp$ from **DESY** Bubble Chamber [V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, IS, Phys Rev C 84, 035203 (2011)]

..... IA

--- IA + NN_{fsi}

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FSI for $\gamma n \rightarrow \pi^- p$

[V. Tarasov, A. Kudryavtsev, W. Briscoe, H. Gao, IS, Phys Rev C 84, 035203 (2011)]

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CLAS for $\gamma \pi \rightarrow \pi \overline{\rho}$ above 1 GeV

[W. Chen et al, Phys Rev C 86, 015206 (2012); Phys Rev Lett 103, 012301 (2009)]

The new CLAS cross sections have quadrupled the world database for $\gamma n \rightarrow \pi^- p$ above $E_{\gamma} = 1 \text{ GeV}.$

> Systematics: Exp: 6-9% FSI: 2-3%

$\chi^2/dp = 4$	45636/ <mark>626</mark> = 72.9	[<mark>SN11</mark> – no fit]
χ²/dp =	1580/626 = 2.5	[GB12 – fit]

• CLAS data appear to have fewer angular structures than the earlier fits.

MAMI-B for $\gamma n \rightarrow \pi^- p$ around the Δ

[W.J. Briscoe et al, Phys Rev C 86, 065207 (2012)]

• MAMI-B data for $\gamma n \rightarrow \pi^- p$ (including FSI corrections) and previous hadronic data for $\pi^- p \rightarrow n\gamma$ appear to **agree well**.

Neutron Multipoles from SAID GB12 [W. Chen et al, Phys Rev C 86, 015206 (2012)]

CLAS Data Impact for Neutron $S = 0 \ll I = \frac{1}{2}$ Couplings [W. Chen et al, Phys Rev C 86, 015206 (2012)]

 BnGa13 and SAID GB12 used the same (almost) data to fit them while BnGa13 has several new Ad Hoc resonances.

Status of Non-strange Resonances

• More than half of states have poor evidence.

GW SAID Contribution

 Most of states need more work to do. Status as seen in — | = 1/2 Most of QCD models predict more states than observed. Status • Where are missing resonances? Particle J^P overall πN Νη Νσ Νω ΛΚ ΣΚ Νρ Δπ γN N $1/2^{+}$ **** 1 = 3/2Status as seen in — $N(1440) 1/2^+$ $N(1520) 3/2^{-1}$ Status Particle J^P $N(1535) 1/2^{-1}$ overall $\pi N \gamma N$ Νη Νσ Νω ΛΚ ΣΚ Νρ Δπ $N(1650) 1/2^{-1}$ $\Delta(1232) 3/2^+$ F **** **** $N(1675) 5/2^{-1}$ *** $\Delta(1600) 3/2^+$ 0 $N(1680) 5/2^+$ **** *** $\Delta(1620) 1/2^{-1}$ N(1685) ?? $\Delta(1700) 3/2^{-1}$ ь $N(1700) 3/2^{-}$ *** *** $\Delta(1750) 1/2^+$ $N(1710) 1/2^+$ *** $\Delta(1900) 1/2^{-1}$ d $N(1720) 3/2^+$ **** $\Delta(1905) 5/2^+$ $N(1860) 5/2^+$ ** $\Delta(1910) 1/2^+$ $N(1875) 3/2^{-}$ *** *** *** $\Delta(1920) 3/2^+$ ** $N(1880) 1/2^+$ * $\Delta(1930) 5/2^{-1}$ *** $N(1895) 1/2^{-}$ *** ** * ** $\Delta(1940) 3/2^{-}$ F ** (seen in $\Delta \eta$) $N(1900) 3/2^+$ *** ** *** $\Delta(1950) 7/2^+$ $N(1990) 7/2^+$ **** *** ** ** ** $\Delta(2000) 5/2^+$ $N(2000) 5/2^+$ ** ** $\Delta(2150) 1/2^{-1}$ $N(2040) 3/2^+$ * $\Delta(2200) 7/2^{-}$ i $N(2060) 5/2^{-}$ ** ** $N(2100) 1/2^+$ $\Delta(2300) 9/2^+$ d ** $N(2150) 3/2^{-}$ $\Delta(2350) 5/2^{-1}$ ** **Bnda Additional** $N(2190) 7/2^{-}$ $\Delta(2390) 7/2^+$ $N(2220) 9/2^+$ $\Delta(2400) 9/2^{-}$ ** \mathbf{n} $N(2250) 9/2^{-1}$ $\Delta(2420) 11/2^+ ****$ $N(2600) 11/2^{-}$ $\Delta(2750) \ 13/2^{-} **$ N(2700) 13/2⁺ ** ** $\Delta(2950) 15/2^+ **$ **

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Summary for Neutron Couplings

• The differential cross section for the processes $\gamma n \rightarrow \pi^- p$ was extracted from new **CLAS** and **MAMI-B** measurements accounting for **Fermi** motion effects in the **IA** as well as **NN**- and πN -**FSI** effects beyond the **IA**.

Consequential calculations of the **FSI** corrections, as developed by the **GW-ITEP** Collaboration, was applied.

- New cross sections departed significantly from our predictions, at the higher energies, and greatly modified the fit result.
- New γn→π⁻p and γn→π⁰n data will provide a critical constraint on the determination of the multipoles and couplings of low-lying baryon resonances using the PWA and coupled channel techniques.

 Polarized measurements [JLab, CB@MAMI, Spring-8, CB-ELSA...] will help to bring more physics in. FSI corrections need to apply.

Recent GRAAL $\sum_{i=1}^{\infty}$ for $\vec{\gamma}_{i} \rightarrow \pi^{0}_{i}$

[R. Di Salvo et al, Eur Phys J A 42, 151 (2009)]

• The difference between previous Pion Prod and new GRAAL measurements may result in significant changes in the neutron couplings.

• 216 GRAAL Σ s are 60% of the World π^0 n data

Recent GRAAL 2 for yn > T

[G. Mandaglio et al, Phys Rev C 82, 045209 (2010)]

• Previous $\gamma n \rightarrow \pi^- p$ measurements provided a better constraint vs. $\gamma n \rightarrow \pi^0 n$ case.

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 Σ for $\vec{\gamma}$ n \rightarrow p π^-

2-beam Asymmetry & FSI [M.I. Levchuk et al, Phys Rev C 74, 014004 (2006)]

PWA7 Workshop, Camogli, Italy, Sept. 2013

 $\Theta_{\pi}^{\mathsf{lab}}$

[deg]

 $\Theta_{\pi}^{\text{lab}}$ [deg]

Meson Production off the Deuteron at CB@MAMI [Spokespersons: W.J. Briscoe and IS, MAMI-A2-02/12]

- We propose to perform a precision measurement of $d\sigma/d\Omega$ in the reactions $\gamma d \rightarrow \pi^- pp$ and $\gamma d \rightarrow \pi^0 np$ in the tagged-photon energy region from **threshold** to **800** MeV.
- The $d\sigma/d\Omega$ for the processes $\gamma n \rightarrow \pi^- p$ and $\gamma n \rightarrow \pi^0 n$ will be extracted from these **CB@MAMI-C** measurements accounting for Fermi motion effects in **IA** as well as NNand πN -**FSI** effects beyond the IA.

- Data were taken in March 2013 and analysis in progress.
- Consequential calculations of the FSI corrections, as developed by the GW-ITEP Collab, will be applied. We have π⁻p FSI already while π⁰n FSI in progress.

 New CB@MAMI data will provide a critical constraint on the determination of the multipoles and couplings of low-laying baryon resonances using the PWA techniques developed by the SAID group.

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Watson's Theorem as Sample of FSI in πN Rescattering

 Connection between scattering and decay processes provides a solid theoretical ground for describing some hadronic effects – ``Watson's theorem."

PHYSICAL REVIEW

VOLUME 95, NUMBER 1

Some General Relations between the Photoproduction and Scattering of π Mesons

KENNETH M. WATSON Department of Physics, University of Wisconsin, Madison, Wisconsin (Received February 19, 1954)

M. Gell-Mann & K.M. Watson, Ann Rev Nucl Sci 4, 219 (1954)

 For pion photoproduction, isospin amplitudes have to satisfy Watson's theorem below the 2π-threshold allowing for a smooth departure from constraint at high energies.

$$\operatorname{atan}(\frac{\operatorname{Im}A_{I}}{\operatorname{Re}A_{I}}) = \delta_{I}(\pi N)$$

• Above 2π -threshold, the rule may still be true, if **inelasticity** of the corresponding π N-elastic amplitude is small (as, *eg*, for P₃₃.)

Evaluation of PionPR Amplitudes below 2π -Threshold

- SAID uses πN PWA results as a constraint for analysis of PionPR data \longrightarrow [see pg 9]
- Most of **PionPR** analyses are doing the same and uses **SAID** π **N** outcome or its modification as input.
- Let us evaluate several **PionPR** analyses such as **SAID**, **MAID**, **EBAC**, **Giessen**, & **BnGa** and compare πN phases coming from πN and **PionPR amplitudes** on proton target.
- SAID database has all these amplitudes which came from authors or we used WebSites.

	πN	PionPR —	
SAID	R. Arndt <i>, et al,</i> PRC 74 , 045205(2006).	R. Workman, et al, PC 2012.	Thresholds
MAID	L. Tiator <i>, et al</i> , PC 2008.	D. Drechsel <i>, et al,</i> EPJA 34 , 69(2007).	React W(GeV) E _y (MeV)
EBAC	B. Julia-Diaz <i>, et al</i> , PRC 76 , 065201(2007).	T. Sato <i>, et al</i> , PRC 54 , 2660(1996).	
Giessen	V. Shklyar, <i>et al</i> , PRC 71 , 055206 (2005).	V. Shklyar, et al, PRC 71 , 055206 (2005).	$\gamma p \rightarrow \pi^+ n$ 1.079 151.4
BnGa	A. Sarantsev, et al, PC 2009.	A. Anisovich <i>, et al,</i> EPJA 48 , 15(2012).	

3/2 Isospin Amplitudes for $E \in M$

1/2 Isospin Amplitudes for $E \in M$

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3/2 I 1/2 Isospin Amplitudes for E I M

Summary for Watson's Theorem for Proton Amplitudes

- Phases coming from πN amplitudes of different analyses are *consistent*.
- Some phases coming from different PionPR analyses are *consistent* to each other and phases coming from πN amplitudes:
 3/2 Isospin Amplitudes: E0+, E1+, & M1+.
- Some phases coming from different PionPR analyses are *inconsistent* to each other and phases coming from πN amplitudes:
 3/2 Isospin Amplitudes: M1-, E2-, M2-, E2+, & M2+.
 1/2 Isospin Amplitudes: E0+, M1-, E1+, M1+, E2-, M2-, E2+, & M2+.
- Some phases coming from E & M multipoles are *inconsistent* to each other and phases coming from πN amplitudes.

