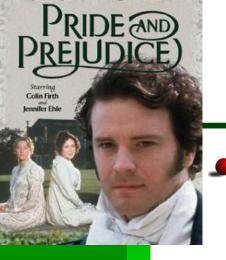


Spectrum of excited states and transition form factors provide unique information about long-range interaction between light-quarks and distribution of hadron's characterising properties amongst its QCD constituents.









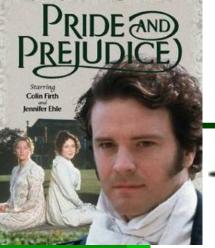
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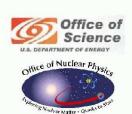






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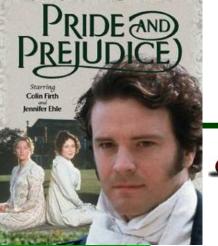
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- p. 2/40



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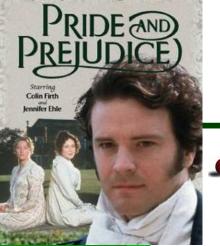
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- p. 2/40

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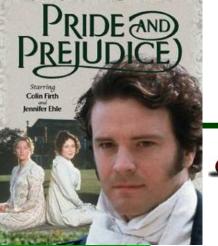






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Back



Spectrum of excited states and transition form factors provide unique information about long-range interaction between light-quarks and distribution of hadron's characterising properties amongst its QCD constituents.

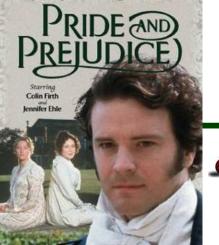
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Challenge: understand relationship between parton properties on the light-front and rest frame structure of hadrons.

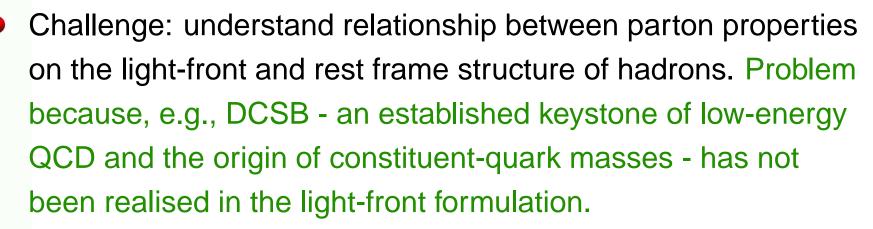






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- **Quark and Gluon Confinement** 
  - No matter how hard one strikes the proton, one cannot liberate an individual quark or gluon









- Quark and Gluon Confinement
  - No matter how hard one strikes the proton, one cannot liberate an individual quark or gluon
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  - Very unnatural pattern of bound state masses
    - e.g., Lagrangian (pQCD) quark mass is small but ... no degeneracy between  $J^{P=+}$  and  $J^{P=-}$









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### Understand Emergent Phenomena

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- Neither of these phenomena is apparent in QCD's Lagrangian yet they are the dominant determining characteristics of real-world QCD.
- QCD Complex behaviour – – arises from apparently simple rules









#### - Goldstone Mode and Bound state







- p. 4/40



#### Goldstone Mode and Bound state

How does one make an almost massless particle ..... from two massive constituent-quarks?









#### Goldstone Mode and Bound state

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Current Algebra ... 1968







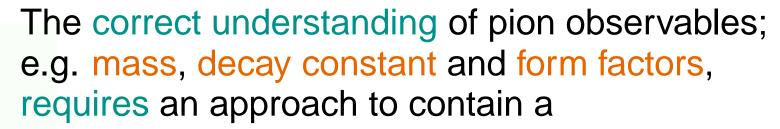
- p. 4/40



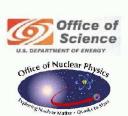
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The correct understanding of pion observables; e.g. mass, decay constant and form factors, requires an approach to contain a

- well-defined and valid chiral limit;
- and an accurate realisation of dynamical chiral symmetry breaking.

**Highly Nontrivial** 



















- Minimal requirements
  - detailed understanding of connection between Current-quark and Constituent-quark masses;
  - and systematic, symmetry preserving means of realising this connection in bound-states.







- p. 5/40

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- p. 5/40

## What's the Problem? Relativistic QFT!

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- p. 5/40

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  - Interaction between quarks the Interquark "Potential" unknown throughout > 98% of a hadron's volume







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#### Intranucleon Interaction



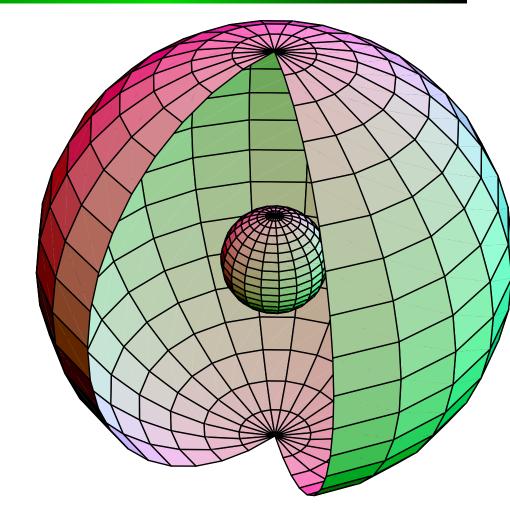






Conclusion

#### Intranucleon Interaction



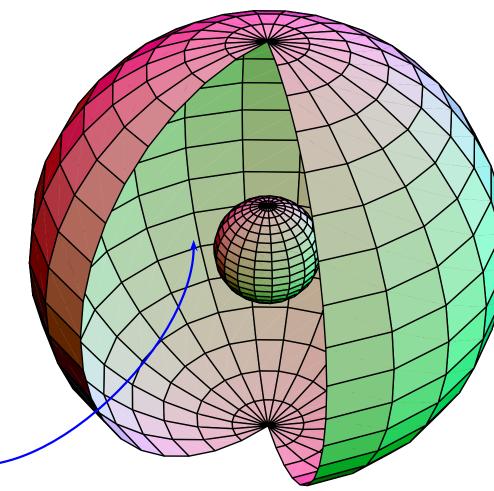






#### Intranucleon Interaction





98% of the volume



## What is the Intranucleon Interaction?

The question must be rigorously defined, and the answer mapped out using experiment and theory.







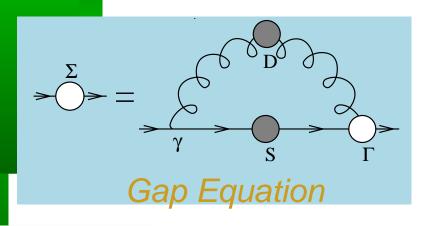
## Frontiers of Nuclear Science: A Long Range Plan (2007)







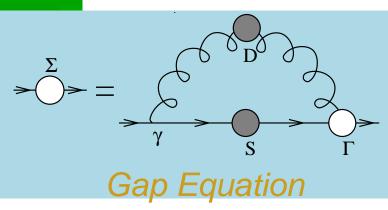
Conclusion

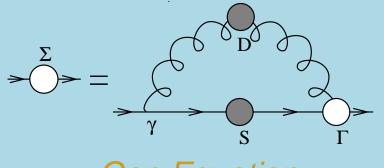


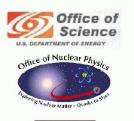








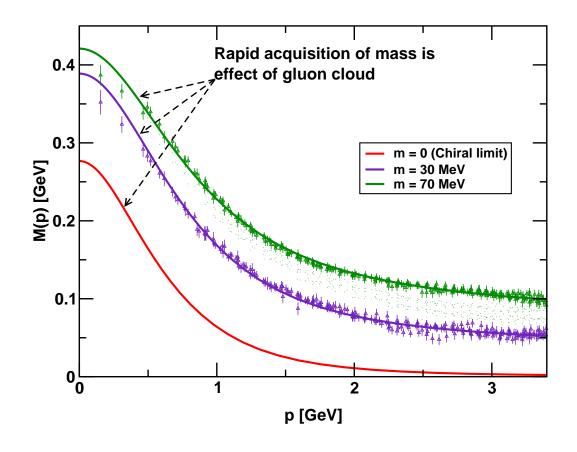






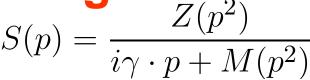


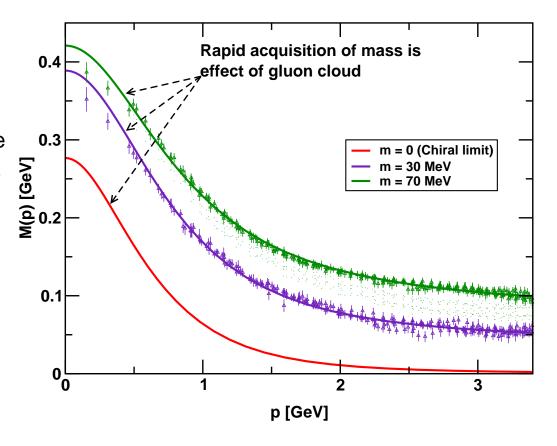
$$S(p) = \frac{Z(p^2)}{i\gamma \cdot p + M(p^2)}$$



Mass from nothing

In QCD a quark's effective mass depends on its momentum. The function describing this can be calculated and is depicted here. Numerical simulations of lattice QCD (data, at two different bare masses) have confirmed model predictions (solid curves) that the vast bulk of the constituent mass of a light quark comes from a cloud of gluons that are dragged along by the quark as it propagates. In this way, a quark that appears to be absolutely massless at high energies (m=0, red curve) acquires a large constituent mass at low energies.







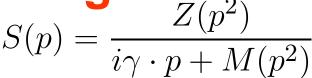


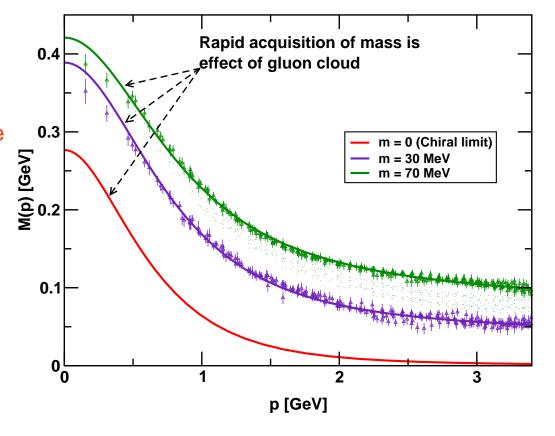




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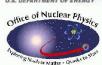


 Established understanding of two- and three-point functions

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- Established understanding of two- and three-point functions
- What about bound states?

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 Without bound states, Comparison with experiment is impossible







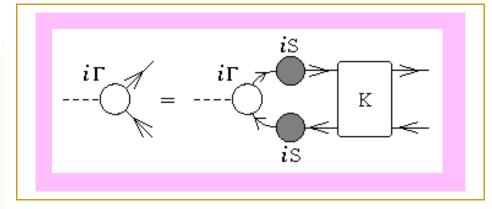
- Without bound states, Comparison with experiment is impossible
- They appear as pole contributions to  $n \ge 3$ -point colour-singlet Schwinger functions







- Without bound states, Comparison with experiment is impossible
- Bethe-Salpeter Equation



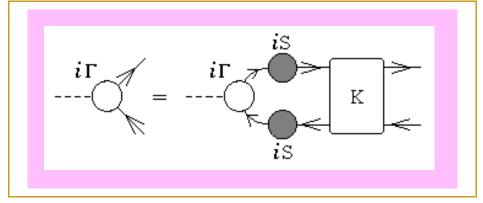
QFT Generalisation of Lippmann-Schwinger Equation.







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QFT Generalisation of Lippmann-Schwinger Equation.

What is the kernel, K?

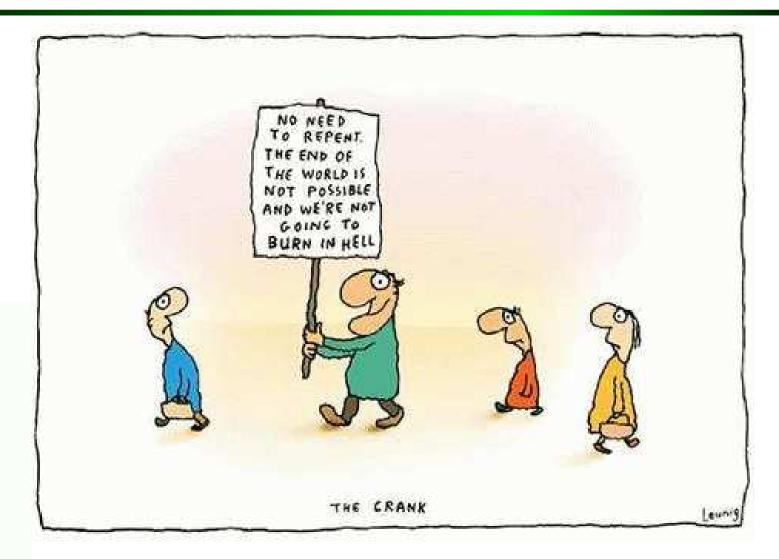
or What is the long-range potential in QCD?







# What is the light-quark Long-Range Potential?









# What is the light-quark Long-Range Potential?



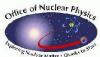






Potential between static (infinitely heavy) quarks measured in simulations of lattice-QCD is not related in any simple way to the light-quark interaction.



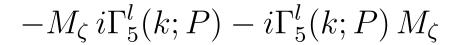






Axial-vector Ward-Takahashi identity

$$P_{\mu} \Gamma^{l}_{5\mu}(k;P) = \mathcal{S}^{-1}(k_{+}) \frac{1}{2} \lambda^{l}_{f} i \gamma_{5} + \frac{1}{2} \lambda^{l}_{f} i \gamma_{5} \mathcal{S}^{-1}(k_{-})$$









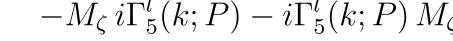


**QFT Statement of Chiral Symmetry** 

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$$-M_{\zeta} i\Gamma_5^l(k;P) - i\Gamma_5^l(k;P) M_{\zeta}$$





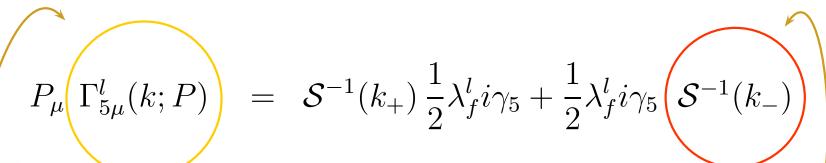








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**Satisfies BSE** 

Satisfies DSE

Kernels very different but must be intimately related

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Satisfies DSE

- →Kernels very different but must be intimately related
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- Nontrivial constraint

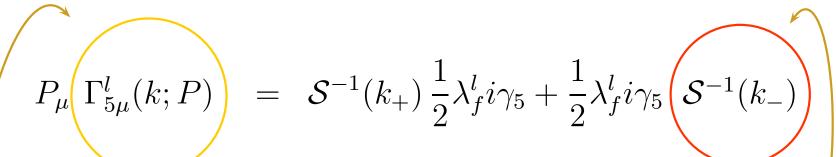








#### Axial-vector Ward-Takahashi identity



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#### Satisfies BSE

Satisfies DSE

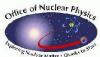


Kernels very different but must be intimately related



- Relation must be preserved by truncation
- Failure ⇒ Explicit Violation of QCD's Chiral Symmetry



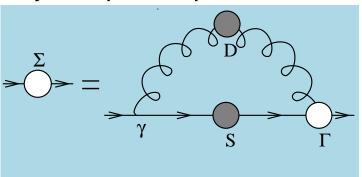








Infinitely Many Coupled Equations











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$$\sum_{\gamma} = \sum_{\gamma} \sum_{S} \sum_{\Gamma}$$

Coupling between equations necessitates truncation









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Infinitely Many Coupled Equations

$$\sum_{\gamma} = \sum_{\gamma} \sum_{S} \sum_{\Gamma}$$

- Coupling between equations necessitates truncation
  - Weak coupling expansion ⇒ Perturbation Theory Not useful for the nonperturbative problems in which we're interested







- p. 11/40



- Infinitely Many Coupled Equations
- There is at least one systematic nonperturbative, symmetry-preserving truncation scheme H.J. Munczek Phys. Rev. D 52 (1995) 4736 Dynamical chiral symmetry breaking, Goldstone's theorem and the consistency of the Schwinger-Dyson and Bethe-Salpeter Equations

A. Bender, C. D. Roberts and L. von Smekal, Phys.

Lett. B **380** (1996) 7

Goldstone Theorem and Diquark Confinement Beyond Rainbow Ladder Approximation









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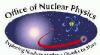






# Radial Excitations & Chiral Symmetry









## & Chiral Symmetry

(Maris, Roberts, Tandy nu-th/9707003)

$$f_H$$
  $m_H^2 = - \rho_{\zeta}^H$   $\mathcal{M}_H$ 



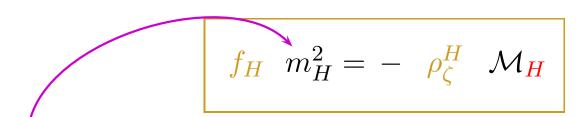




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# & Chiral Symmetry

(Maris, Roberts, Tandy nu-th/9707003)



Mass<sup>2</sup> of pseudoscalar hadron







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#### & Chiral Symmetry

(Maris, Roberts, Tandy nu-th/9707003)

$$f_{H} \quad m_{H}^{2} = - \quad \rho_{\zeta}^{H} \quad \mathcal{M}_{H}$$

$$\mathcal{M}_{H} := \operatorname{tr}_{\text{flavour}} \left[ M_{(\mu)} \left\{ T^{H}, \left( T^{H} \right)^{\text{t}} \right\} \right] = m_{q_{1}} + m_{q_{2}}$$

Sum of constituents' current-quark masses

$$ullet$$
 e.g.,  $T^{K^+}=rac{1}{2}\left(\lambda^4+i\lambda^5
ight)$ 

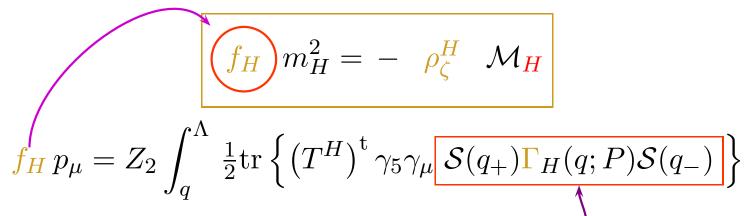






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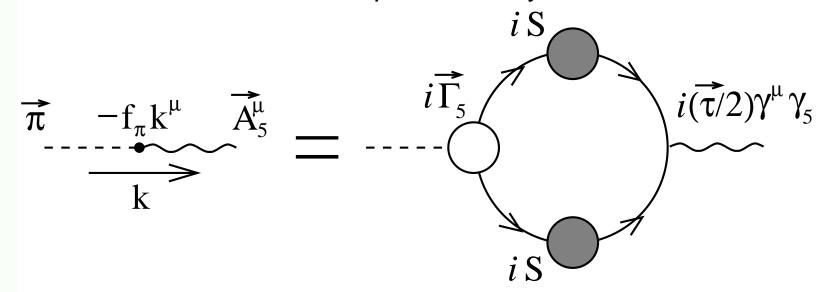


- Pseudovector projection of BS wave function at x=0
- Pseudoscalar meson's leptonic decay constant





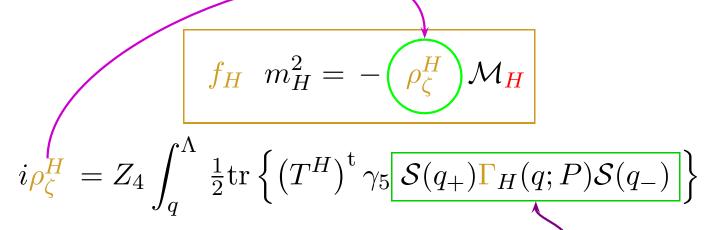




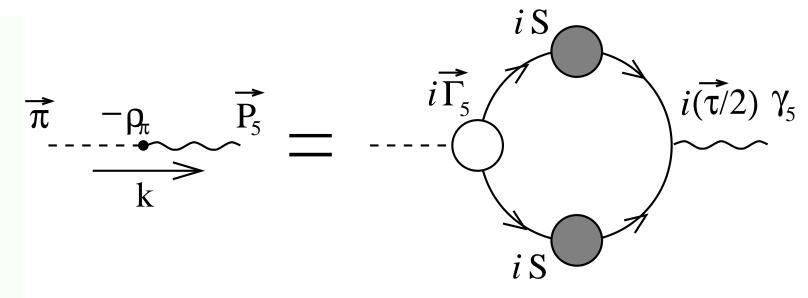
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& Chiral Symmetry

(Maris, Roberts, Tandy nu-th/9707003



• Pseudoscalar projection of BS wave function at x=0









## & Chiral Symmetry

(Maris, Roberts, Tandy nu-th/9707003)

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Light-quarks; i.e.,  $m_q \sim 0$ 

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$$f_H o f_H^0$$
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Hence 
$$m_H^2 = rac{-\langle ar q q 
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 ...GMOR relation, a corollary







## & Chiral Symmetry

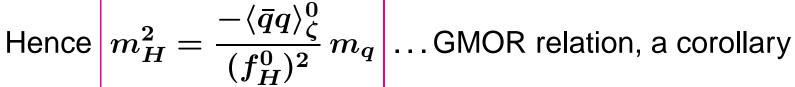
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Heavy-quark + light-quark

$$\Rightarrow f_H \propto rac{1}{\sqrt{m_H}}$$
 and  $ho_{\zeta}^H \propto \sqrt{m_H}$ 

Hence, 
$$m_H \propto m_q$$







QCD Proof of Potential Model result

CLAS12 European Workshop ... 23 transparencies

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# New Challenges









#### New Challenges

Next Steps ... Applications to excited states and axial-vector mesons, e.g., will improve understanding of confinement interaction between light-quarks.







#### New Challenges

- Next Steps ... Applications to excited states and axial-vector mesons, e.g., will improve understanding of confinement interaction between light-quarks.
- Move on to the problem of a symmetry preserving treatment of hybrids and exotics.

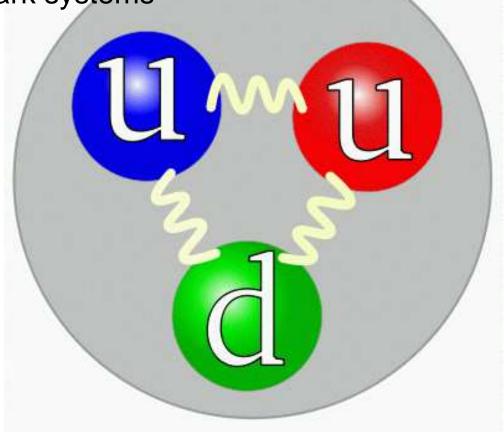






Another Direction ... Also want/need information about

three-quark systems



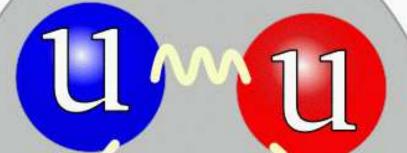






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Another Direction . . . Also want/need information about three-quark systems



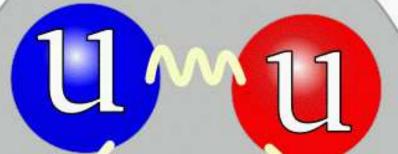
With this problem ... most wide-ranging studies employ expertise familiar from meson applications circa  $\sim$  1995.







Another Direction ... Also want/need information about three-quark systems



• With this problem ... most wide-ranging studies employ expertise familiar from meson applications circa  $\sim$  1995.

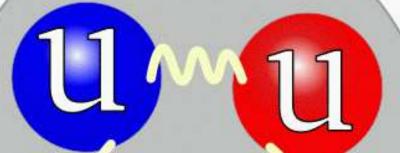






Namely ... Model-building and Phenomenology, constrained by the DSE results outlined already.

Another Direction . . . Also want/need information about three-quark systems



- With this problem ... most wide-ranging studies employ expertise familiar from meson applications circa  $\sim$  1995.
  - However, that is beginning to change ...









# **Unifying Study** of Mesons and Baryons









# Unifying Study of Mesons and Baryons

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# **Unifying Study** of Mesons and Baryons

How does one incorporate dressed-quark mass function,  $M(p^2)$ , in study of baryons? Behaviour of  $M(p^2)$  is essentially a quantum field theoretical effect.

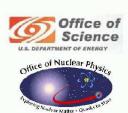






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  - Residue is proportional to nucleon's Faddeev amplitude
  - Poincaré covariant Faddeev equation sums all possible exchanges and interactions that can take place between three dressed-quarks
  - Tractable equation is founded on observation that an interaction which describes colour-singlet mesons also generates quark-quark (diquark) correlations in the colour-3 (antitriplet) channel







# Faddeev equation

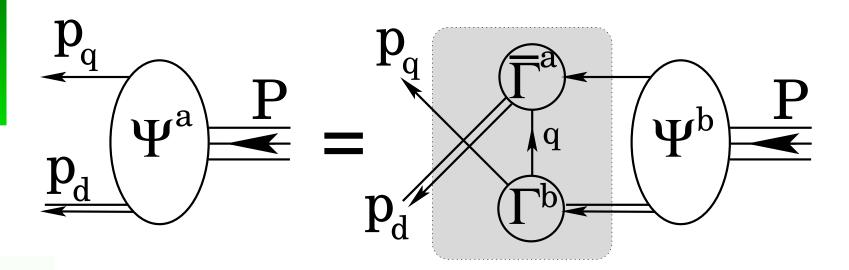




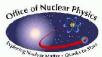




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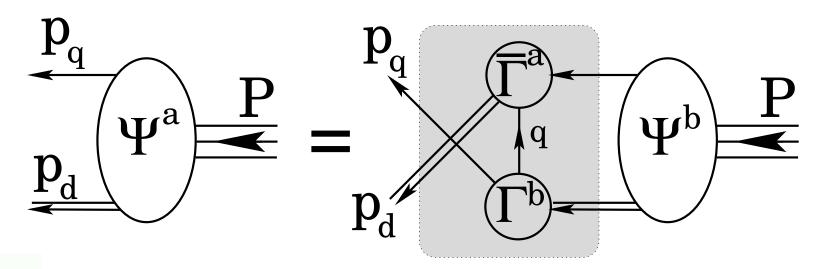




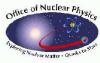




## Faddeev equation





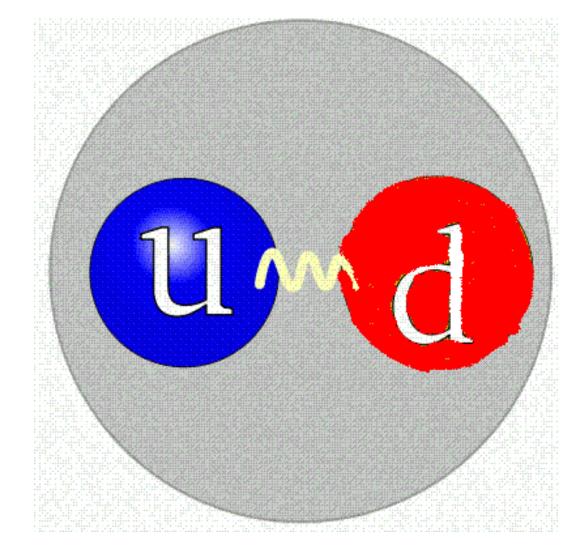






- Linear, Homogeneous Matrix equation
- Yields wave function (Poincaré Covariant Faddeev Amplitude) that describes quark-diquark relative motion within the nucleon
- Scalar and Axial-Vector Diquarks ... In Nucleon's Rest Frame Amplitude has ... s-, p- & d-wave correlations

# Diquark correlations









### **QUARK-QUARK**

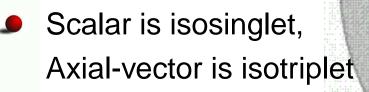
### Same interaction that

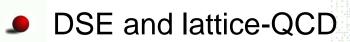
# Diquark correlations

describes mesons also generates three coloured quark-quark correlations:

blue-red, blue-green, green-red

Confined ... Does not escape from within baryon





$$m_{[ud]_{0^+}} = 0.74 - 0.82$$

$$m_{(uu)_{1^{+}}} = m_{(ud)_{1^{+}}} = m_{(dd)_{1^{+}}} = 0.95 - 1.02$$







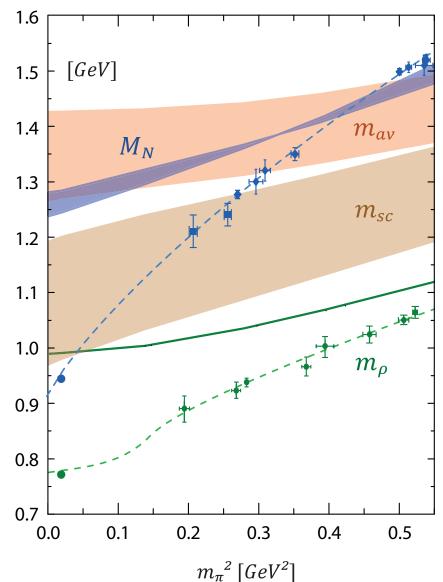




- arXiv:0802.1948 [nucl-th]
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# **Ab-initio study** of mesons & nucleons

Bethe-Salpeter & Faddeev equations built from same RG-improved rainbow-ladder interaction





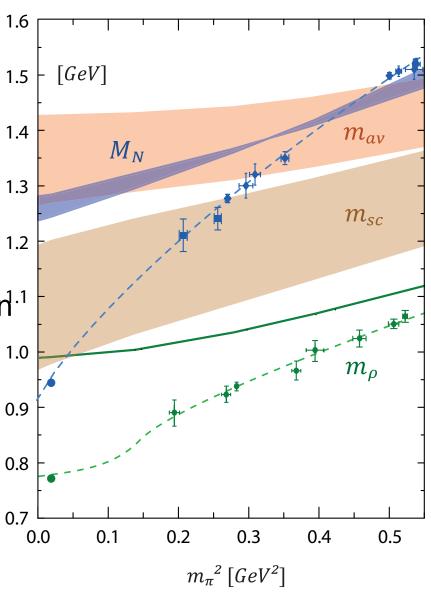




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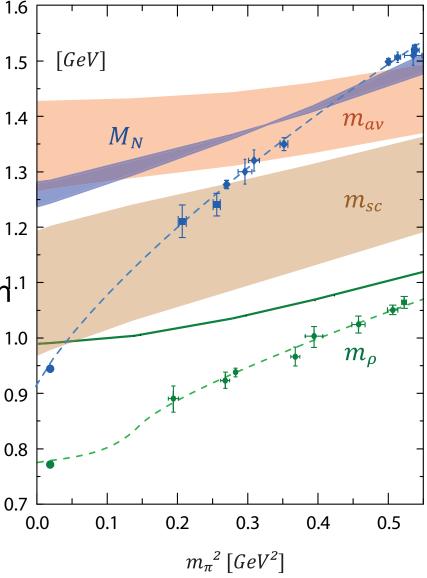
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Continuum prediction for evolution of  $m_{
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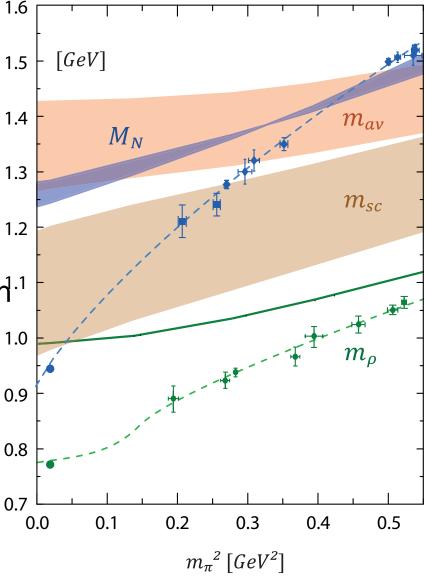
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Continuum prediction for evolution of  $m_{
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Systematically improvable



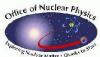






### **Nucleon-Photon Vertex**









M. Oettel, M. Pichowsky and L. von Smekal, nu-th/9909082 6 terms ...

### **Nucleon-Photon Vertex**

constructed systematically ... current conserved automatically for on-shell nucleons described by Faddeev Amplitude







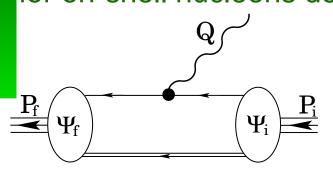
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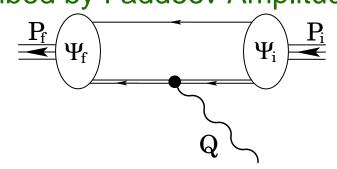
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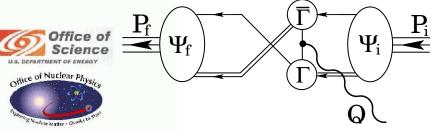
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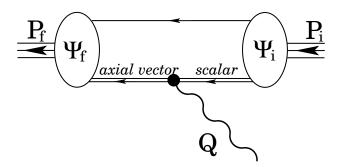
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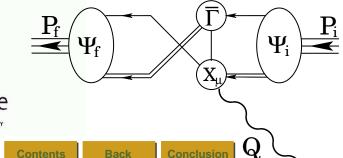


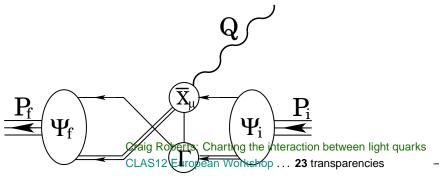














# **DSE-based** Faddeev Equation











- arXiv:0710.2059 [nucl-th]
- arXiv:0710.5746 [nucl-th]
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# DSE-based Faddeev Equation

Faddeev equation input – algebraic parametrisations of DSE results, constrained by  $\pi$  and K observables







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Faddeev equation input – algebraic parametrisations of DSE results, constrained by  $\pi$  and K observables

Two parameters

$$-M_{0^+} = 0.8\,{
m GeV},$$

$$M_{
m 1+}=0.9\,{
m GeV}$$

chosen to give

$$M_N = 1.18, M_{\Delta} = 1.33$$

allow for pseudoscalar meson contributions







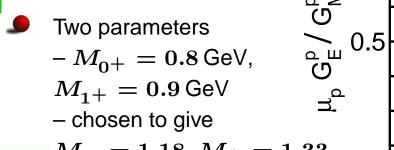


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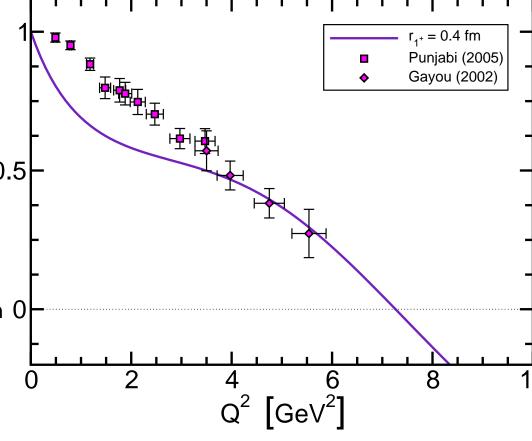
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 $M_N=1.18,\,M_\Delta=1.33$  – allow for pseudoscalar meson 0 contributions

Sensitivity to details of the current – expressed through diquark radius









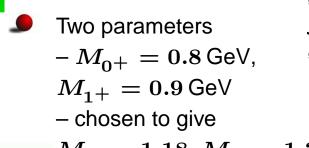


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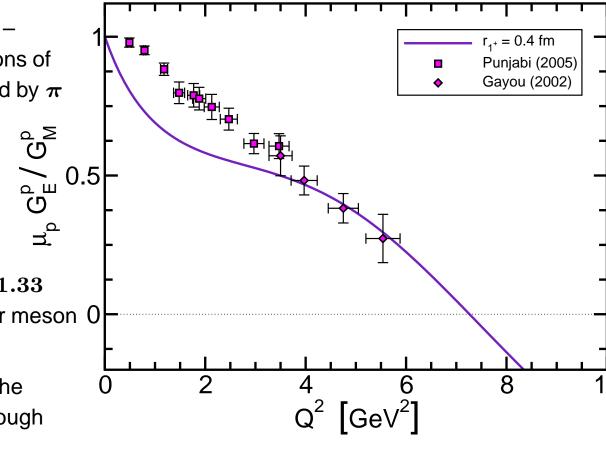
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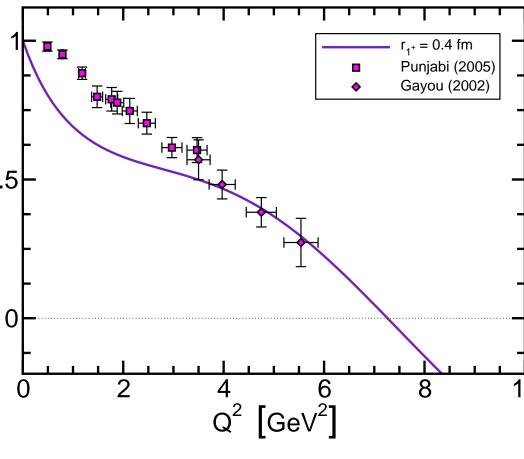
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Sensitivity to details of the current – expressed through diquark radius







- On  $Q^2 \lesssim 4\,{\rm GeV^2}$  result lies below experiment. This can be attributed to omission of pseudoscalar-meson-cloud contributions
- Always a zero but position depends on details of current



# ab initio Faddeev Equation









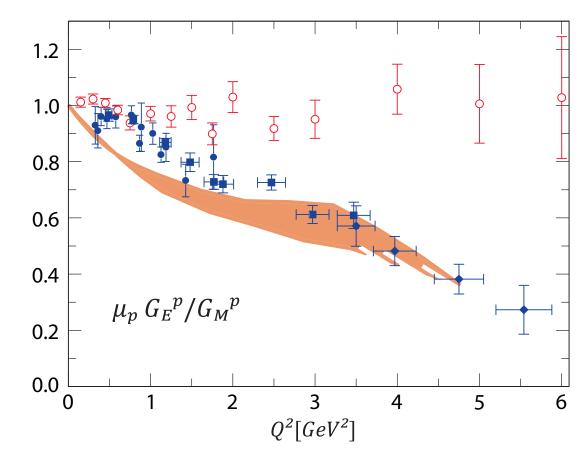


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### ab initio

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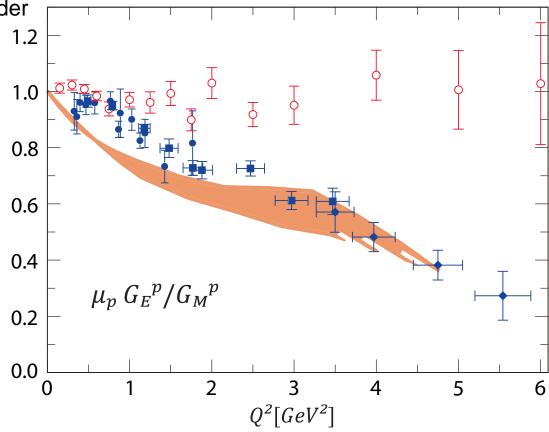


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### Faddeev Equation

Parameter-free rainbow-ladder Faddeev equation – result 1.3 qualitatively identical and in semiquantitative agreement 1.5









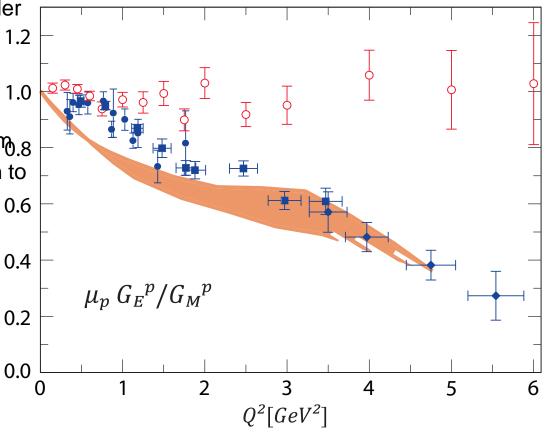
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Parameter-free rainbow-ladder
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Improved numerical algorithm  $_{.8}$  needed to extend calculation to larger  $Q^2$  0.6









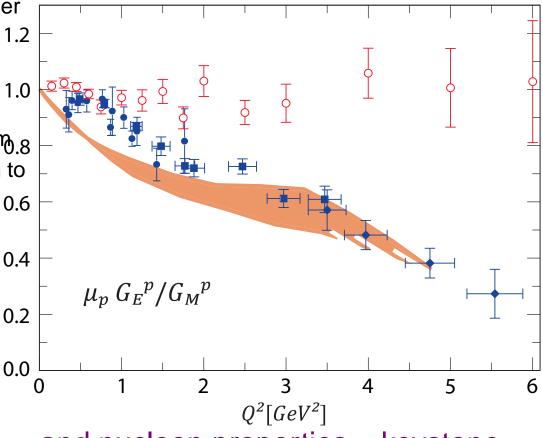
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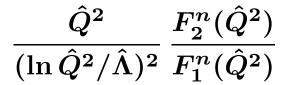






Calculation unifies  $\pi$ ,  $\rho$  and nucleon properties – keystone is behaviour of dressed-quark mass function and hence veracious description of QCD's Goldstone mode

# Ratio of Neutron Pauli & Dirac Form Factors



$$\hat{\Lambda}=\Lambda/M_N=0.44$$

Ensures proton ratio constant for  $\hat{Q}^2 \geq 4$ 





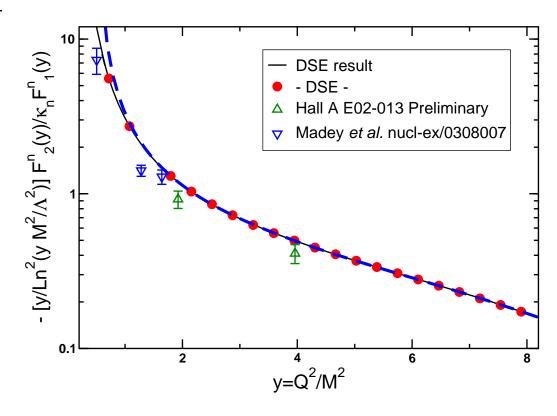


### Ratio of Neutron Pauli & Dirac Form Factors

$$rac{\hat{Q}^2}{(\ln \hat{Q}^2/\hat{\Lambda})^2} \, rac{F_2^n(\hat{Q}^2)}{F_1^n(\hat{Q}^2)}$$

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## Ratio of Neutron Pauli & Dirac Form Factors

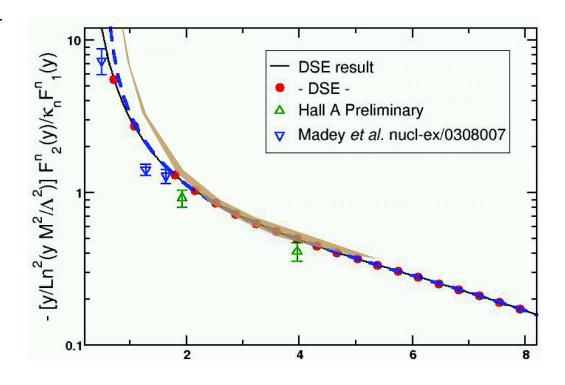
$$\frac{\hat{Q}^2}{(\ln \hat{Q}^2/\hat{\Lambda})^2}\,\frac{F_2^n(\hat{Q}^2)}{F_1^n(\hat{Q}^2)}$$

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ab initio RL result









#### **Pion Cloud**









## **Pion Cloud** F2 - neutron







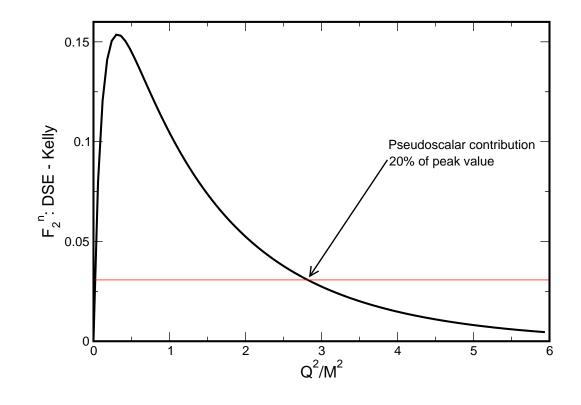


Conclusion

### **Pion Cloud** F2 – neutron

Comparison between Faddeev equation result and Kelly's parametrisation

**Faddeev** equation set-up to describe dressed-quark core







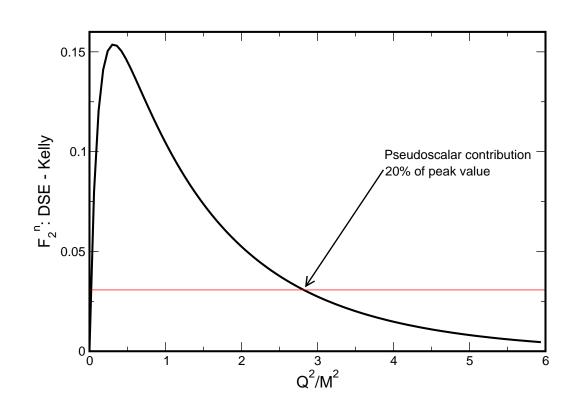




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Pseudoscalar meson cloud (and related effects) significant for  $Q^2 \lesssim 3-4\,M_N^2$ 



## **Epilogue**









**First** 



## **Epilogue**









**First** 



DCSB exists in QCD.









**First** 

Contents



DCSB exists in QCD.



It is manifest in dressed propagators and vertices









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#### DCSB exists in QCD.



- It is manifest in dressed propagators and vertices
- It predicts, amongst other things, that
  - light current-quarks become heavy constituent-quarks: 4 → 400 MeV
  - pseudoscalar mesons are unnaturally light:  $m_{\rho}=770$  cf.  $m_{\pi}=140\,\mathrm{MeV}$
  - pseudoscalar mesons couple unnaturally strongly to light-quarks:  $g_{\pi \bar{q} q} \approx 4.3$
  - pseudscalar mesons couple unnaturally strongly to the lightest baryons

$$g_{\pi \bar{N}N} \approx 12.8 \approx 3g_{\pi \bar{q}q}$$



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$$g_{\pi \bar{N}N} \approx 12.8 \approx 3g_{\pi \bar{q}q}$$

It impacts dramatically upon observables.



### **Epilogue**

- Dyson-Schwinger Equations
  - Poincaré covariant unification of meson and baryon observables







Conclusion



#### **Epilogue**

- **Dyson-Schwinger Equations** 
  - Poincaré covariant unification of meson and baryon observables
  - All global and pointwise corollaries of DCSB are manifested naturally without fine-tuning











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  - Poincaré covariant unification of meson and baryon observables
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Conclusion

- Mesons already being studied
- Baryons are within practical reach











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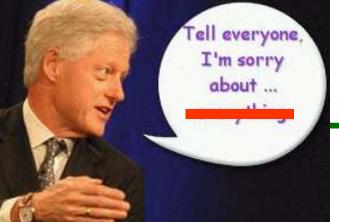
Conclusion

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#### nothing!





- Poincaré covariant unification of meson and baryon observables
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- Excited states:
  - Mesons already being studied
  - Baryons are within practical reach
- ullet Ab-initio study of  $N o \Delta$  transition underway
- Tool enabling insight to be drawn from experiment into long-range piece of interaction between light-quarks









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- 4. Dressed-Quark Propagator
- 5. Frontiers of Nuclear Science
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- 8. Bethe-Salpeter Kernel
- 9. Persistent Challenge
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- 15. Nucleon Challenge
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- 17. Faddeev equation
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- 19. Ab-initio study of mesons & nucleons

20. 
$$r_{\pi} f_{\pi}$$

- 21. Nucleon-Photon Vertex
- 22. DSE-based Faddeev Equation
- 23. Ratio of Neutron FFs
- 24. Pion Cloud







### **Dyson-Schwinger Equations**







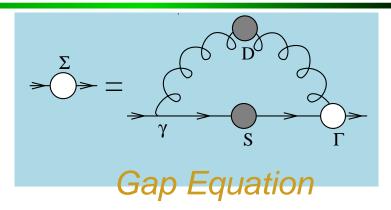








$$S(p) = \frac{Z(p^2)}{i\gamma \cdot p + M(p^2)}$$

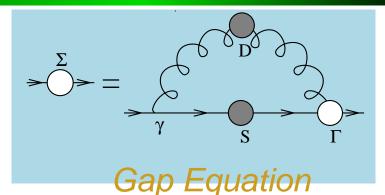








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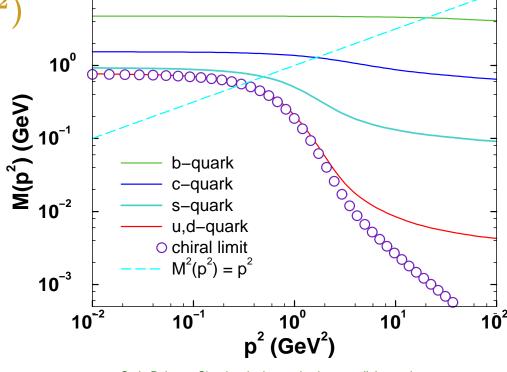
Gap Equation's Kernel Enhanced on IR domain

 $\Rightarrow$  IR Enhancement of  $M(p^2)$ 



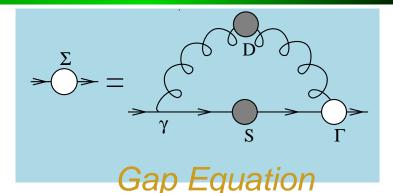






Craig Roberts: Charting the interaction between light quarks CLAS12 European Workshop ... 13 transparencies

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Gap Equation's Kernel Enhanced on IR domain





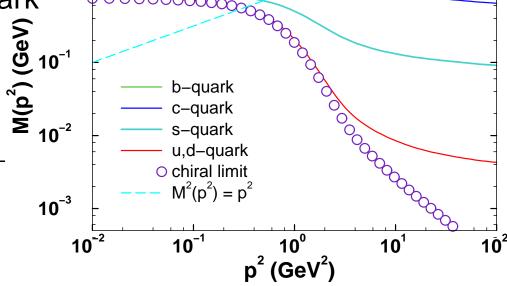


Mass:  $M_f^E$ :  $p^2 = M(p^2)^2$ 



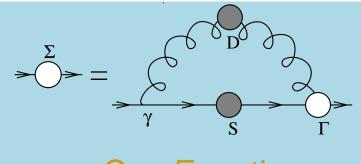
flavour	u/d	s	c	b
$rac{M^E}{m_{\mathcal{L}}}$	$\sim 10^2$	$\sim 10$	$\sim 1.5$	$\sim 1.1$





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$$S(p) = \frac{Z(p^2)}{i\gamma \cdot p + M(p^2)}$$



Gap Equation

b-quark c-quark

Gap Equation's Kernel Enhanced on IR domain







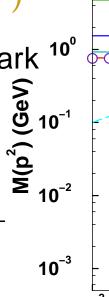
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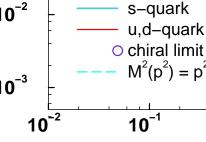


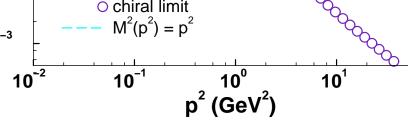
flavour	u/d	s	c	b
$\frac{M^E}{m_{\zeta}}$	$\sim 10^2$	~ 10	$\sim 1.5$	~ 1.1



Predictions confirmed in numerical simulations of lattice-QCD







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10

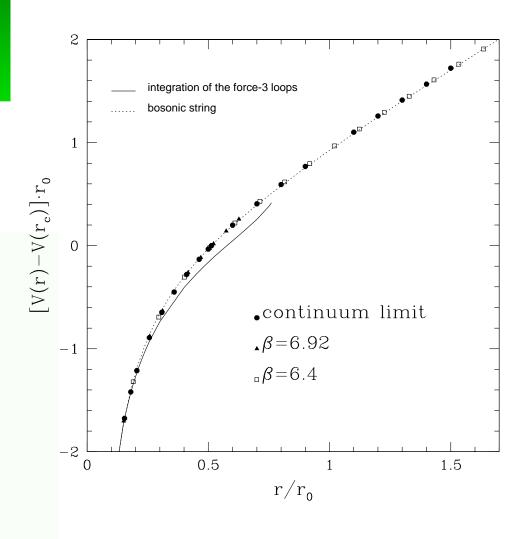








#### Infinitely Heavy Quarks ... Picture in Quantum Mechanics

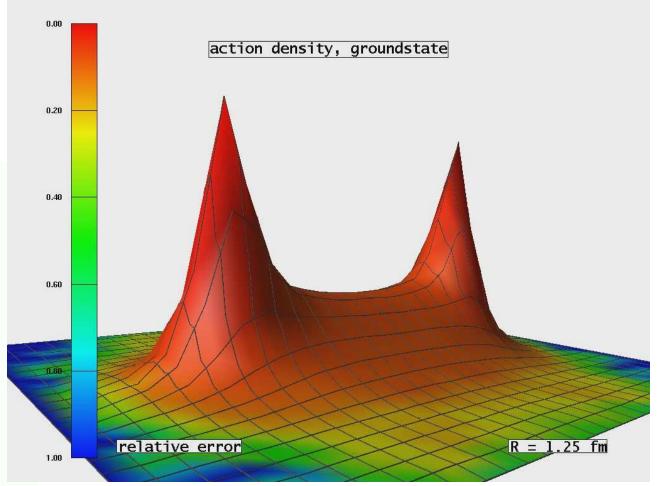


$$V(r) = \sigma\,r - rac{\pi}{12}\,rac{1}{r}$$

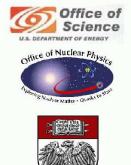
 $\sigma \sim 470\,\mathrm{MeV}$ 

Necco & Sommer he-la/0108008

 $m extcolor{ll}$  Illustrate this in terms of the action density . . . analogous to plotting the Force  $=F_{ar QQ}(r)=\sigma+rac{\pi}{12}rac{1}{r^2}$ 

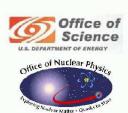


Bali, *et al*. he-la/0512018





What happens in the real world; namely, in the presence of light-quarks?

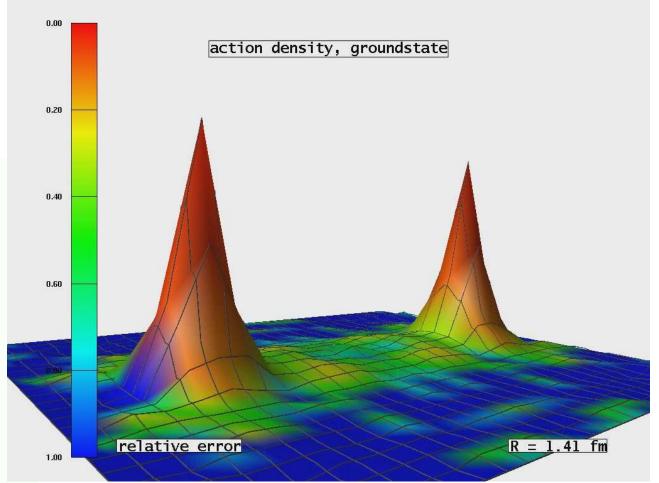






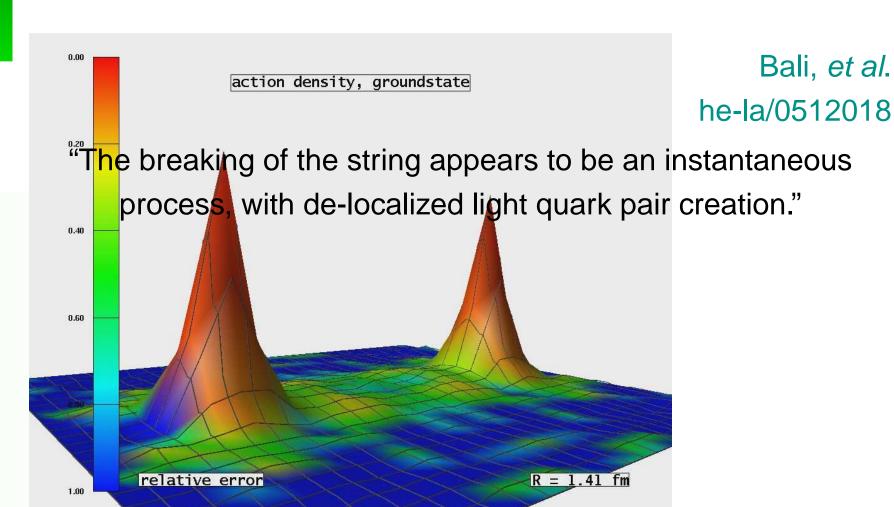
Conclusion

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Bali, *et al.* he-la/0512018

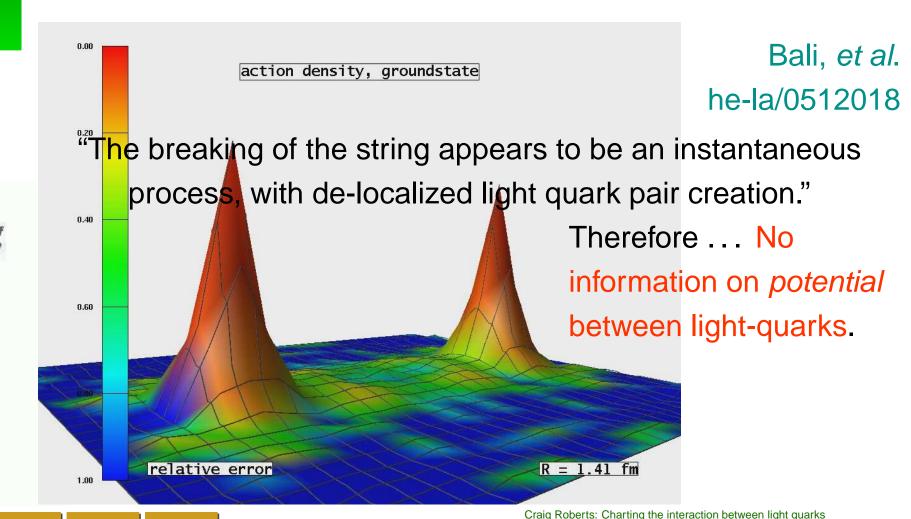
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## Radial Excitations

& Chiral Symmetry

Höll, Krassnigg, Roberts

nu-th/0406030

$$f_H$$
  $m_H^2 = - \rho_{\zeta}^H$   $\mathcal{M}_H$ 

Valid for ALL Pseudoscalar mesons







## Höll, Krassnigg, Roberts nu-th/0406030

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- p. 29/40

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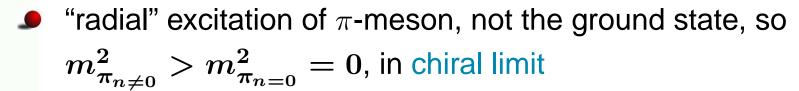
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Dynamical Chiral Symmetry Breaking
 Goldstone's Theorem -









McNeile and Michael he-la/0607032

# Radial Excitations & Lattice-QCD







Conclusion

#### McNeile and Michael he-la/0607032

# Radial Excitations & Lattice-QCD

When we first heard about [this result] our first reaction was a combination of "that is remarkable" and "unbelievable".







-p.30/40

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- $m ext{CLEO: } au o \pi(1300) + 
  u_{ au} \ o f_{\pi_1} < 8.4 \, ext{MeV} \ ext{Diehl & Hiller} \ ext{he-ph/0105194}$





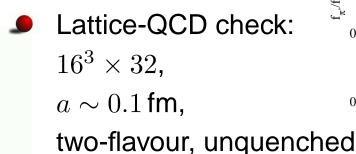


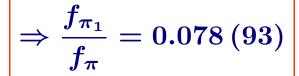
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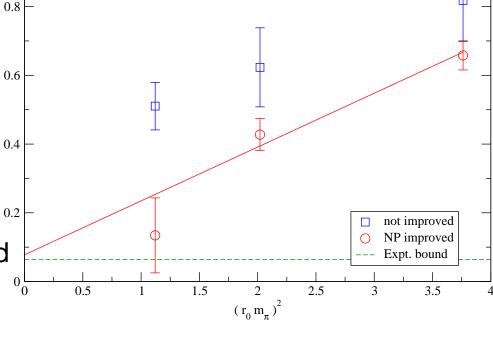
#### & Lattice-QCD

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# Radial Excitations

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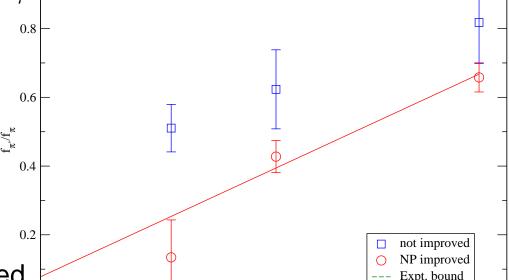
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Lattice-QCD check:

 $16^3 \times 32$ 

 $a \sim 0.1$  fm,









two-flavour, unquenched

$$\Rightarrow rac{f_{\pi_1}}{f_{\pi}} = 0.078 \, (93)$$

Full ALPHA formulation is required to see suppression, because PCAC relation is at the heart of the conditions imposed for improvement (determining coefficients of irrelevant operators)

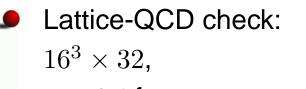
 $(r_0 m_{-})^2$ 

# Radial Excitations

#### & Lattice-QCD

#### McNeile and Michael he-la/0607032

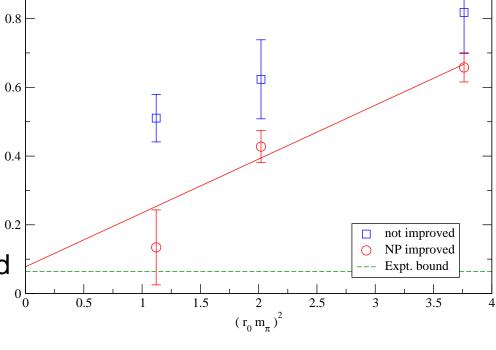
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The suppression of  $f_{\pi_1}$  is a useful benchmark that can be used to tune and validate lattice QCD techniques that try to determine the properties of excited states mesons craig Roberts: Charting the interaction between light quarks

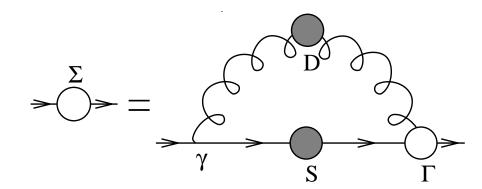
#### Procedure Now Straightforward







- Solve Gap Equation
  - $\Rightarrow$  Dressed-Quark Propagator, S(p)

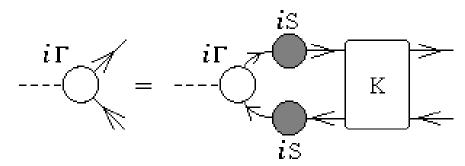








- Use that to Complete Bethe Salpeter Kernel, K
- Solve Homogeneous Bethe-Salpeter Equation for Pion Bethe-Salpeter Amplitude,  $\Gamma_{\pi}$

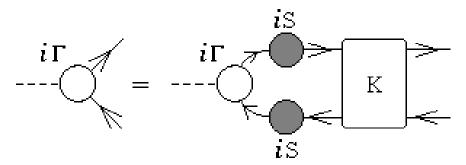








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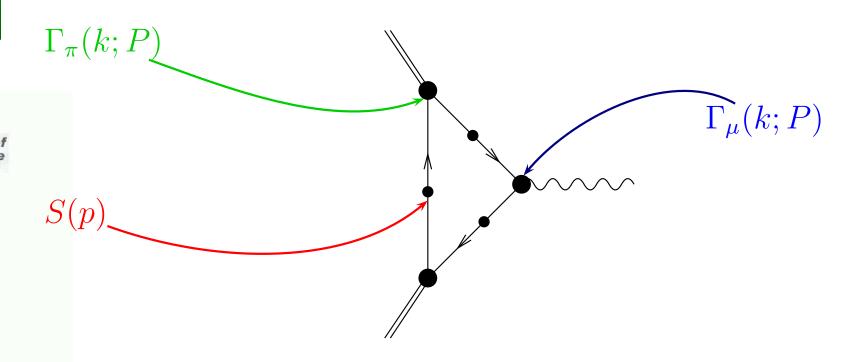




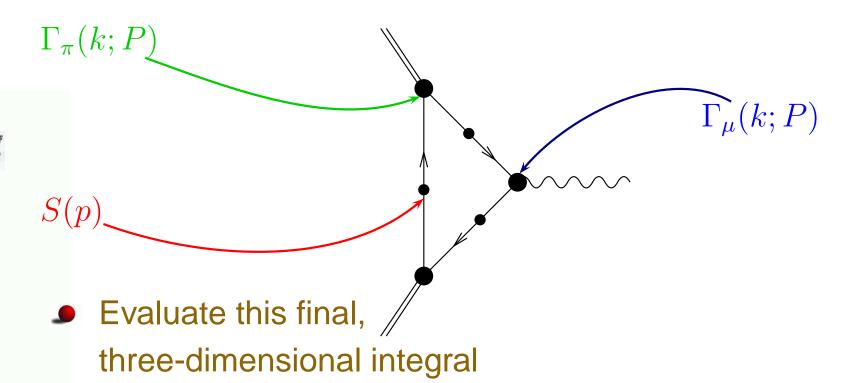
Solve Inhomogeneous Bethe-Salpeter Equation for Dressed-Quark-Photon Vertex,  $\Gamma_{\mu}$ 

Conclusion

Now have all elements for Impulse Approximation to Electromagnetic Pion Form factor

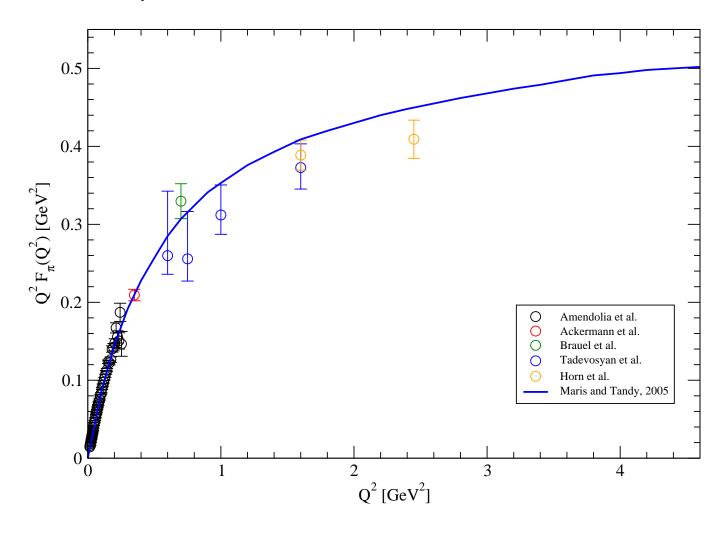


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#### Calculated Pion Form Factor

Calculation first published in 1999; No Parameters Varied Numerical method improved in 2005









**Back** 

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#### Calculated Pion Form Factor

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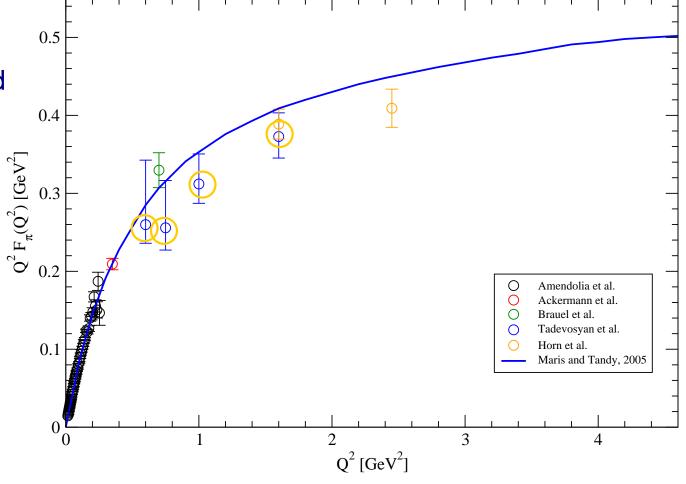
Data published in 2001.

Subsequently revised





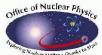






#### **Timelike Pion Form Factor**











Ab initio calculation into timelike region. Deeper than ground-state  $\rho$ -meson pole timelike Pion Form Factor







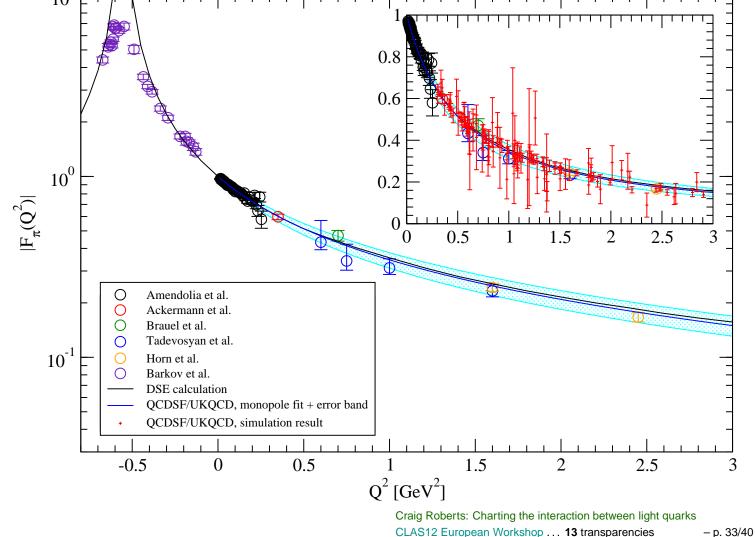


Ab initio calculation into timelike region. Deeper than ground-state  $\rho$ -meson pole Timelike Pion Form Factor







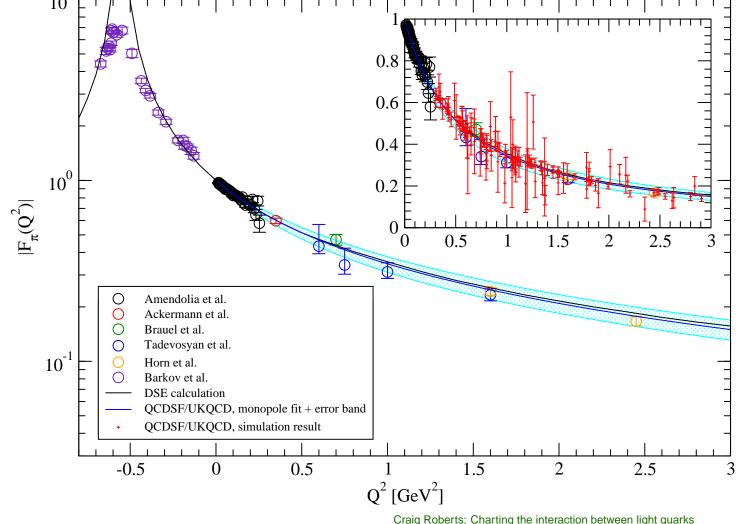




Ab initio calculation into timelike region. Deeper than

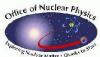
ground-state  $\rho$ -meson pole Timelike Pion Form Factor

ρ-meson not put in "by hand" – generated dynamically as a boundstate of dressed-quark and dressed-antiquark













Conclusion

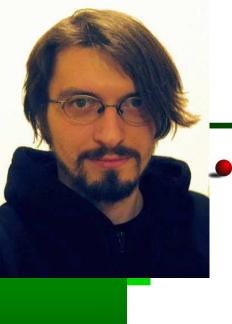












Improved rainbow-ladder interaction









Improved rainbow-ladder interaction

Repeating  $F_{\pi}(Q^2)$  calculation









Improved rainbow-ladder interaction

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Experimentally:  $r_{\pi}f_{\pi}=0.315\pm0.005$ 





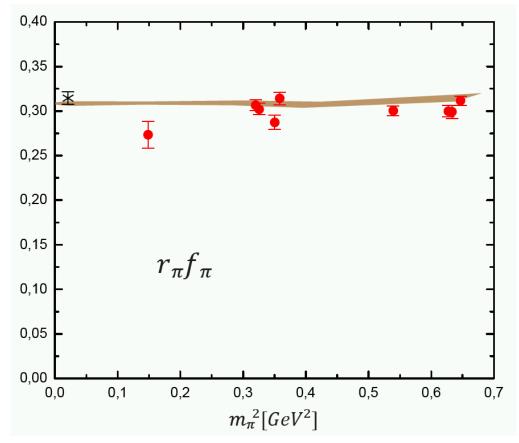


- Improved rainbow-ladder interaction
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- DSE prediction



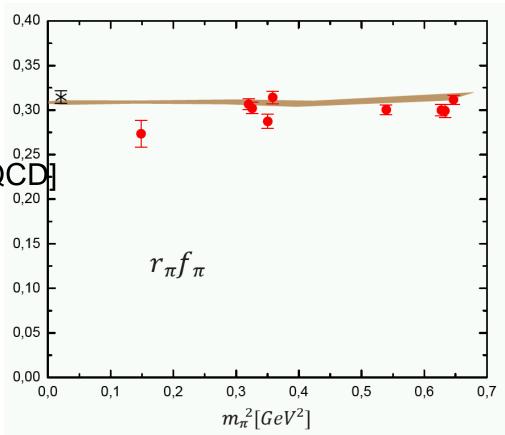






- p. 34/40

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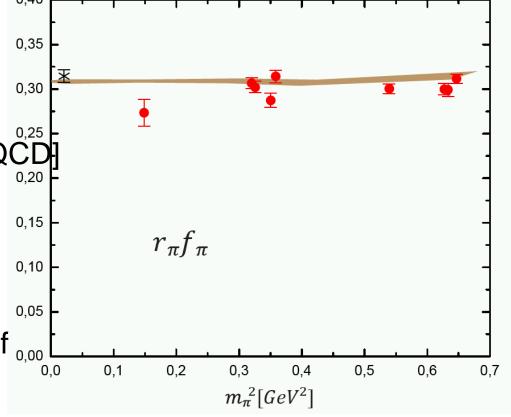


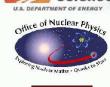






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**DSE** prediction

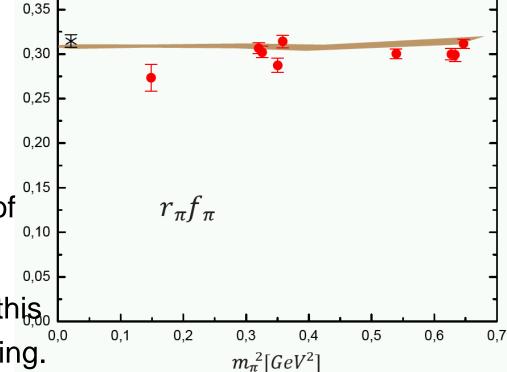
Fascinating result:

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We have understood this Implications far-reaching.

Conclusion









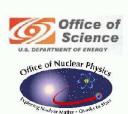
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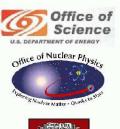






## Infidelity without simultaneity

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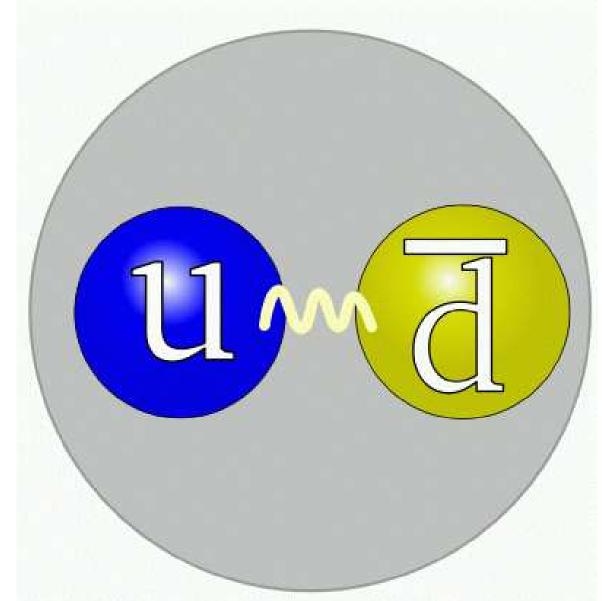
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- Must similarly require prediction of  $\gamma^*\pi \to \pi\pi$  and all other anomalous processes

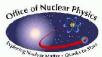














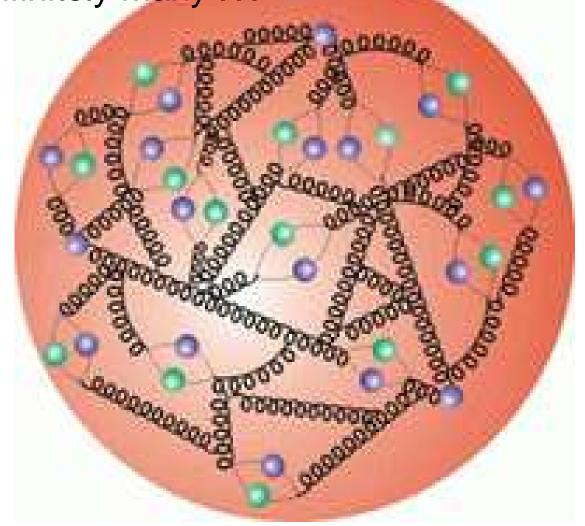




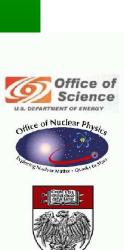




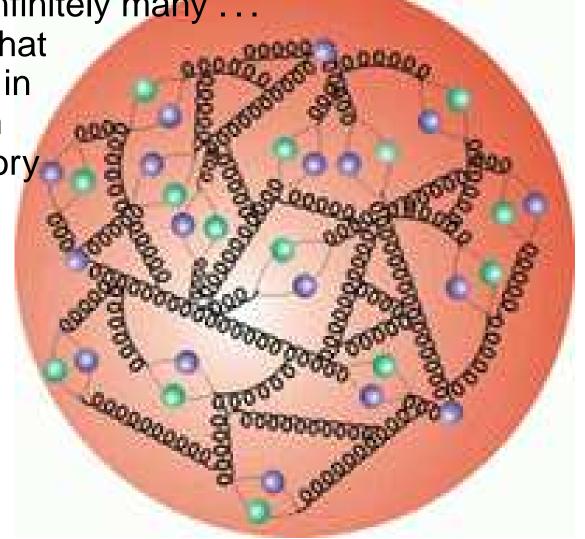


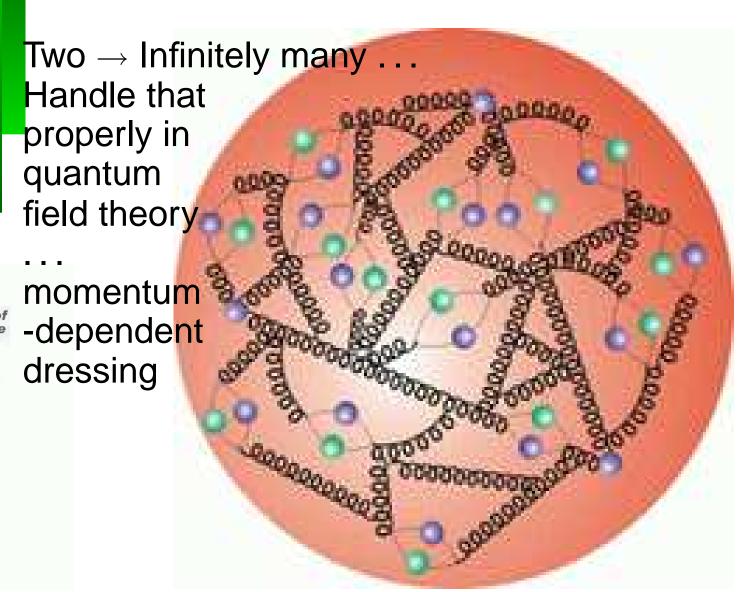


Two  $\rightarrow$  Infinitely many .... Handle that properly in quantum field theory



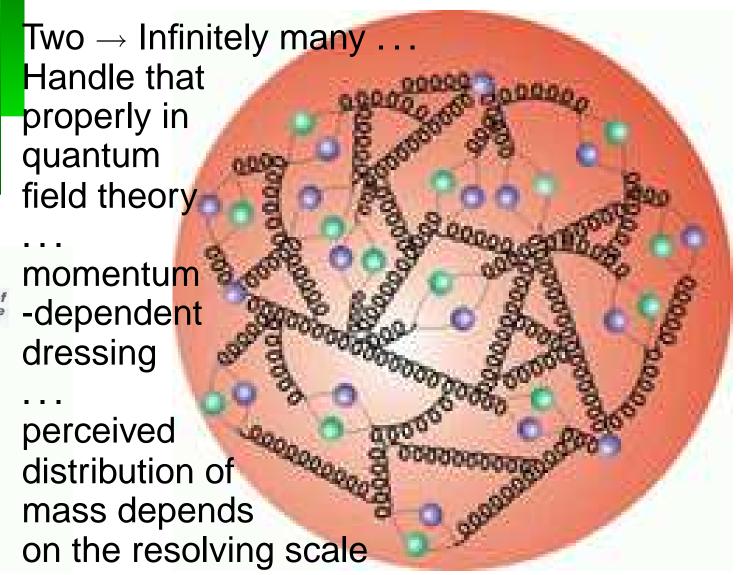








### Answer for the pion







### **Ab-initio study** of mesons & nucleons







Conclusion



- arXiv:0802.1948 [nucl-th]

– arXiv:0810.1222 [nucl-th]

## Ab-initio study of mesons & nucleons







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Conclusion

- arXiv:0802.1948 [nucl-th]
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## Ab-initio study of mesons & nucleons

Leading-order truncation of DSEs – rainbow-ladder







-p.37/40

- arXiv:0802.1948 [nucl-th]
- arXiv:0810.1222 [nucl-th]

- Leading-order truncation of DSEs rainbow-ladder
- Corrections vanish with increasing current-quark mass
  - ⇒ rainbow-ladder exact in heavy-quark limit







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  - Use this knowledge to constrain interaction in infrared
  - Interaction in ultraviolet predicted by perturbative expansion of DSEs







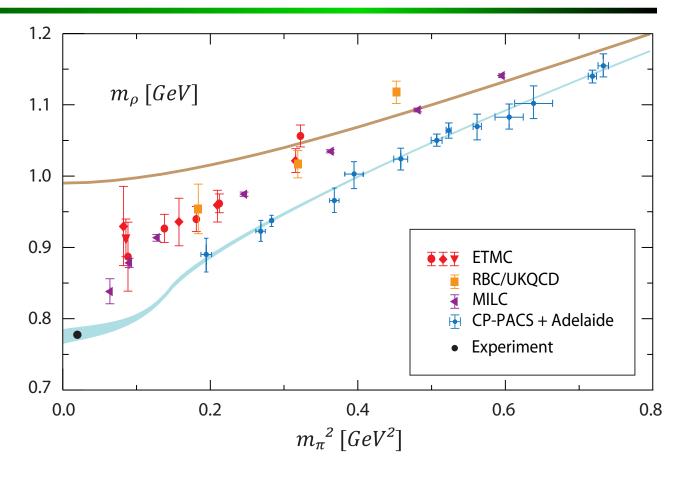




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## Ab-initio study of mesons & nucleons





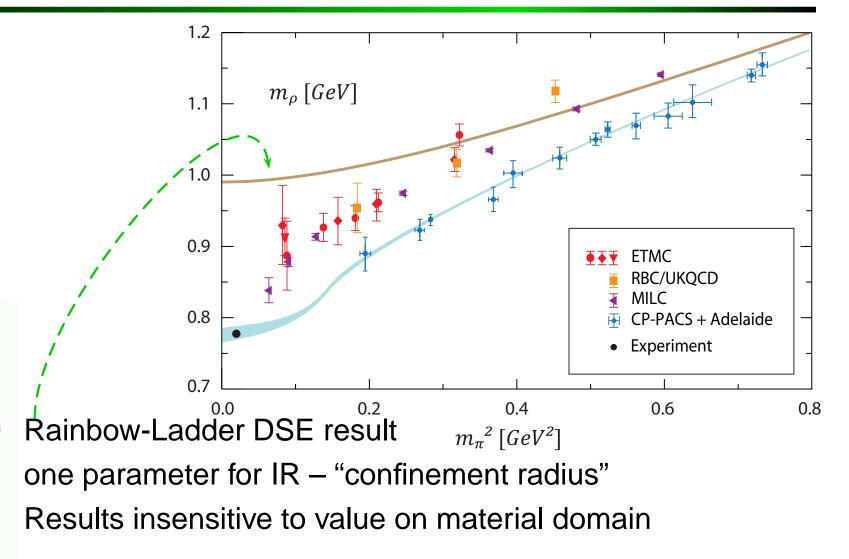




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Conclusion

- arXiv:0802.1948 [nucl-th]
- arXiv:0810.1222 [nucl-th]



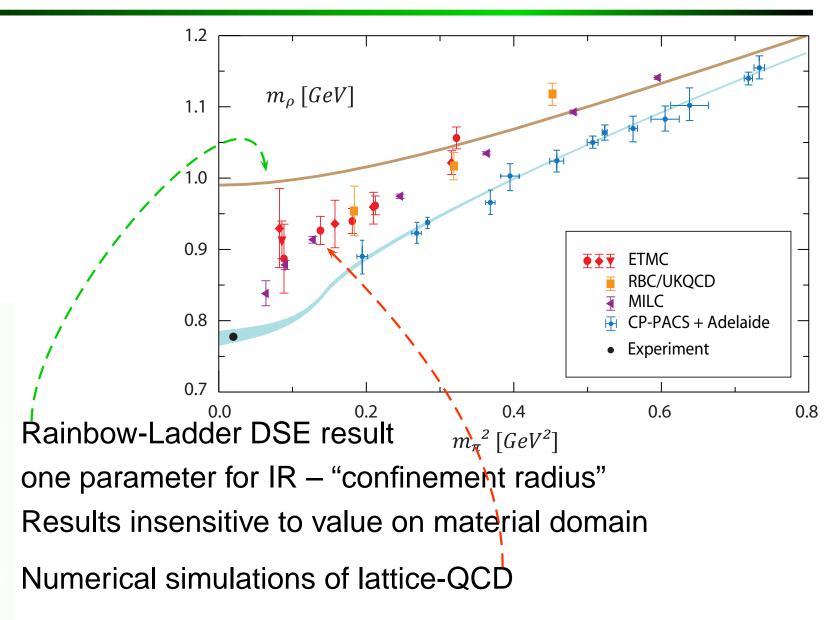








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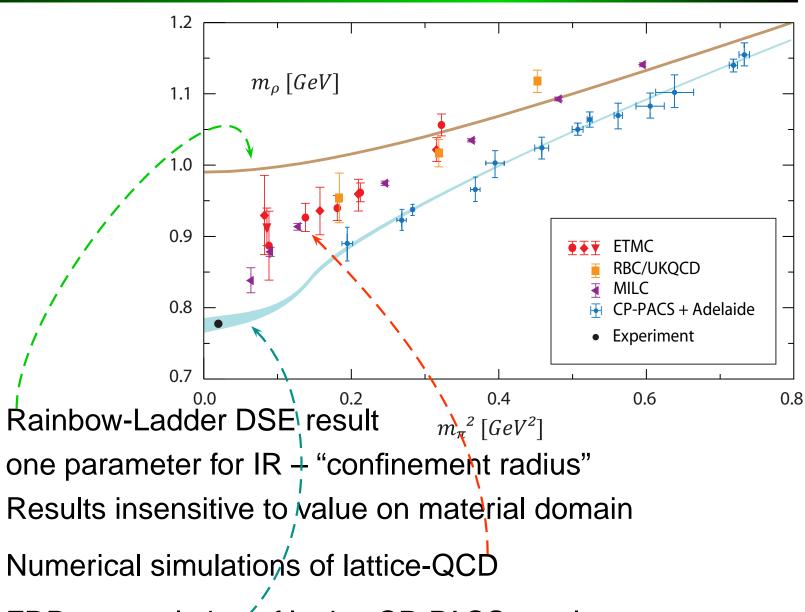








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# Ab-initio study of mesons & nucleons

Precisely the same interaction



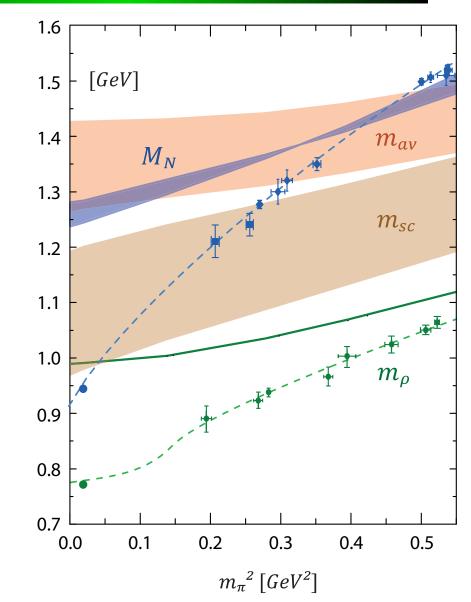




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## Ab-initio study of mesons & nucleons

- Precisely the same interaction
- Same  $\rho$ -meson curve







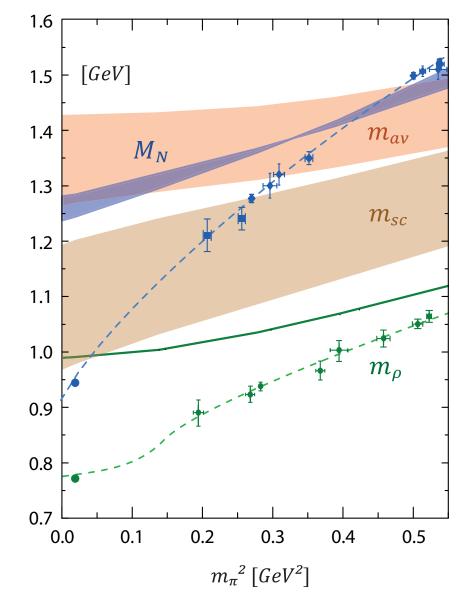


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- arXiv:0802.1948 [nucl-th]
- arXiv:0810.1222 [nucl-th]

- Precisely the same interaction
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  - "unobservable" show marked sensitivity to single model parameter; viz., confinement radius







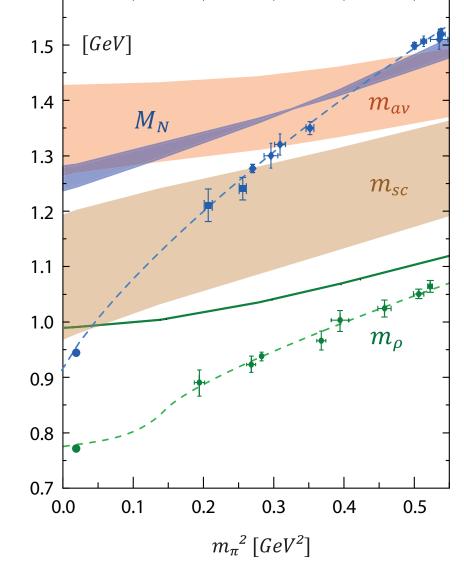


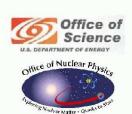
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## Ab-initio study of mesons & nucleons

1.6

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- But ...  $[m_{av}-m_{sc}], m_{
  ho}$  &  $M_N$  ... are independent of that parameter





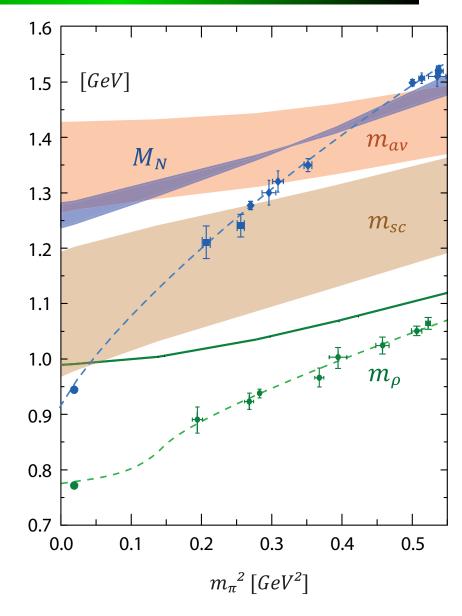




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# Ab-initio study of mesons & nucleons

Parameter-independent RL-DSE predictions, with veracious description of Goldstone mode





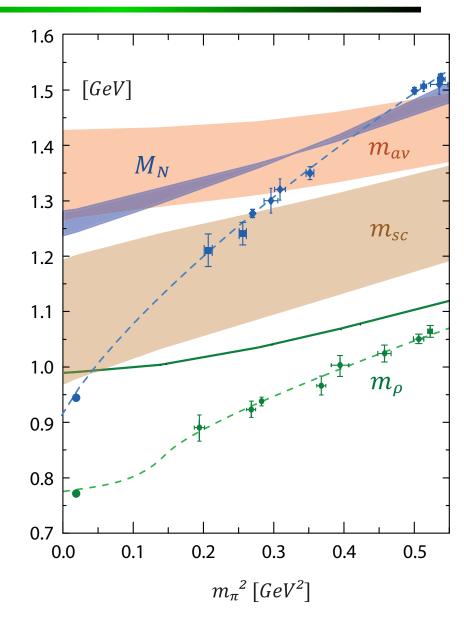




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## Ab-initio study of mesons & nucleons

- Parameter-independent RL-DSE predictions, with veracious description of Goldstone mode
- DSE and lattice agree on heavy-quark domain









-p.40/40

- arXiv:0802.1948 [nucl-th]
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### **Ab-initio** study

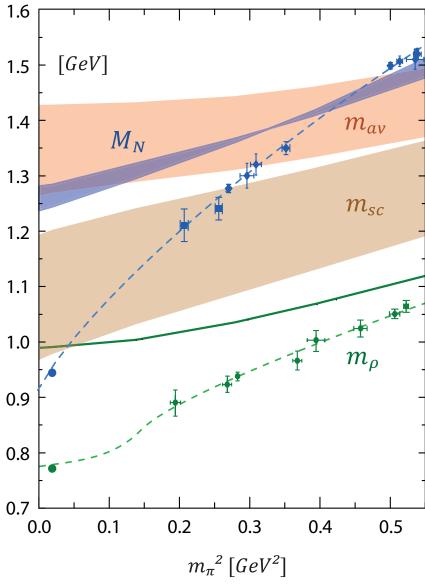
### of mesons & nucleons

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Prediction: at physical  $m_{\pi}^2$ ,  $M_N^{
m quark-core}=1.26(2)\,{
m GeV}$ cf. FRR+lattice-QCD,  $= 1.27(2) \, {\rm GeV}_{_{0.8}}$ 



⇒ subleading corrections, including 0<sup>-</sup>-meson loops,  $\delta M_N = -320 \, {\rm MeV}$ ,



Conclusion  $\delta m_
ho = -220\,{
m MeV}$