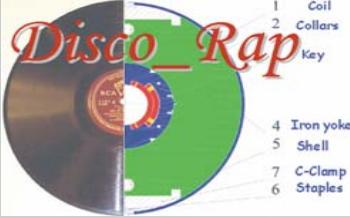


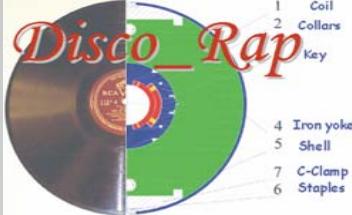
Short resume on magnetic measurements of NbTi wires



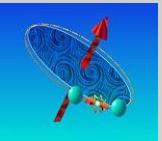
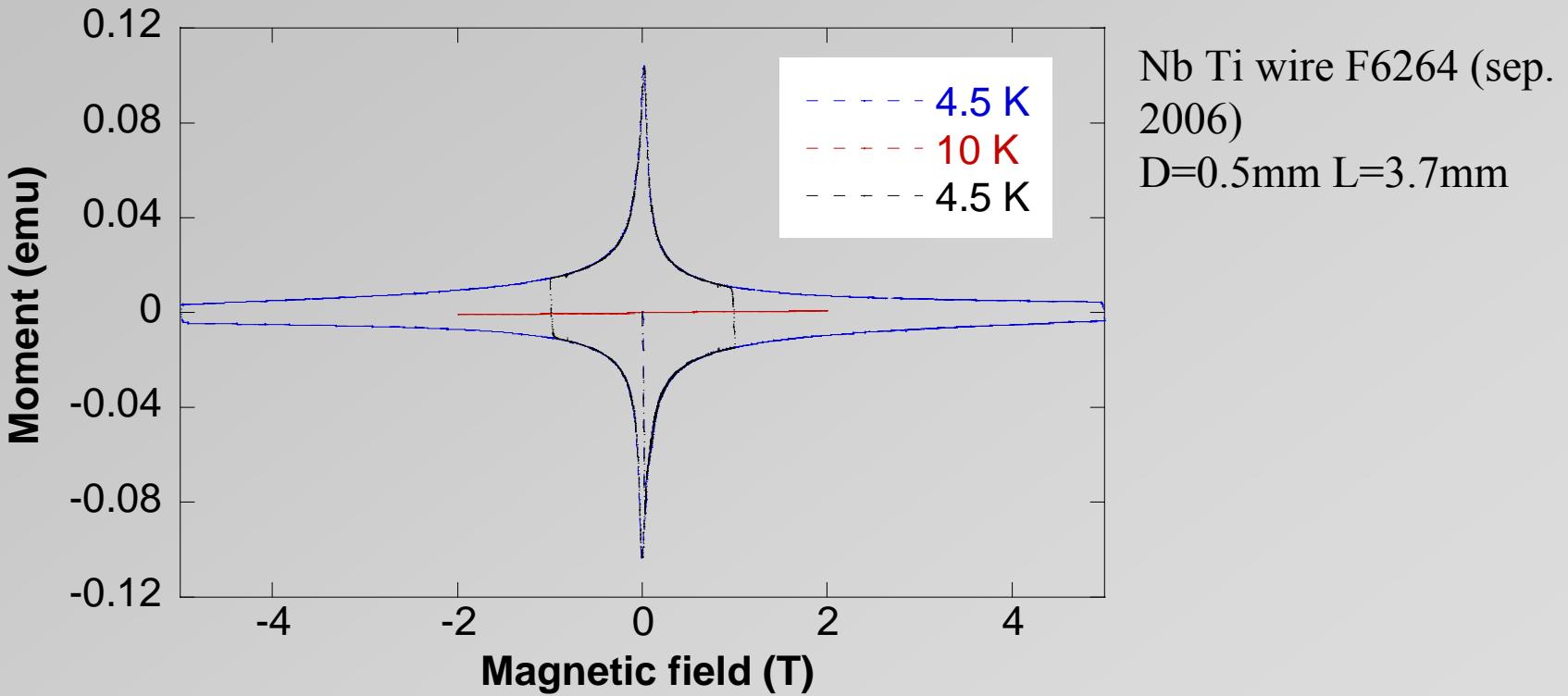
Here we wish to summarize results on NbTi wires made of thin filaments.

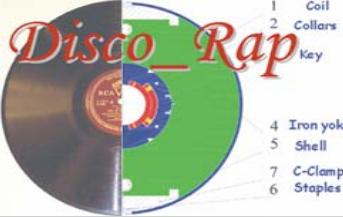
$M(H)$ curves were recorded with different magnetometers (VSM 12T, VSM 16 T, 5T SQUID Magnetometer), which have been cross checked against the calibrations.



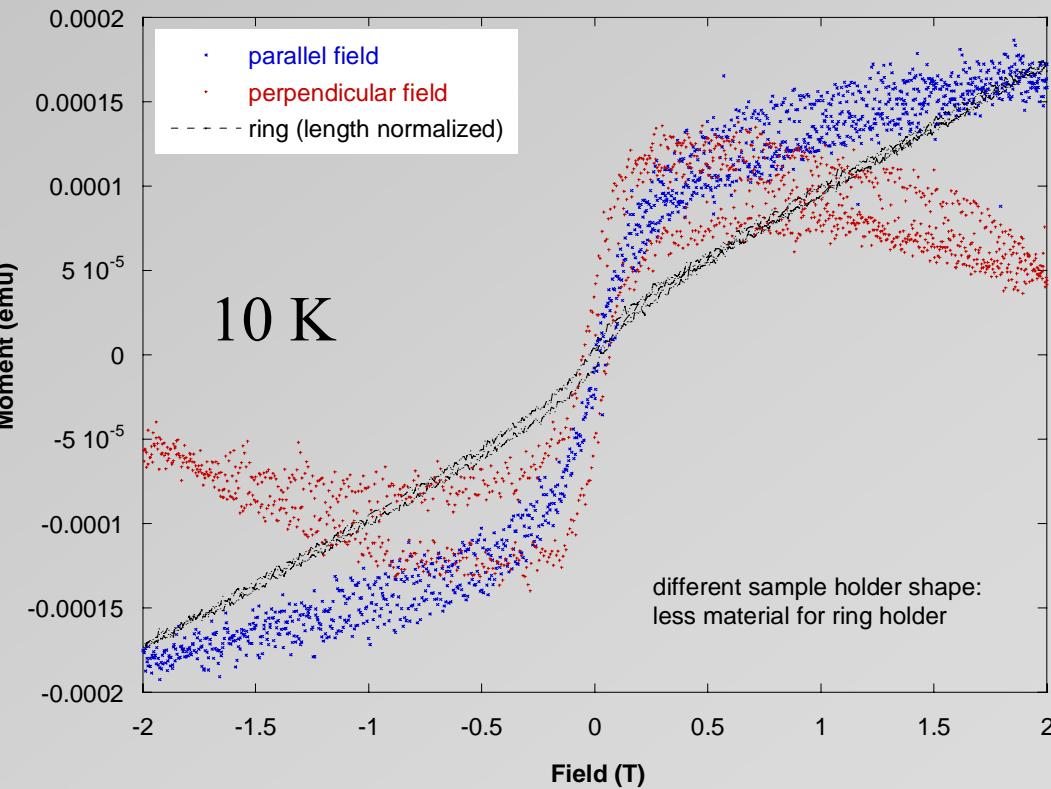
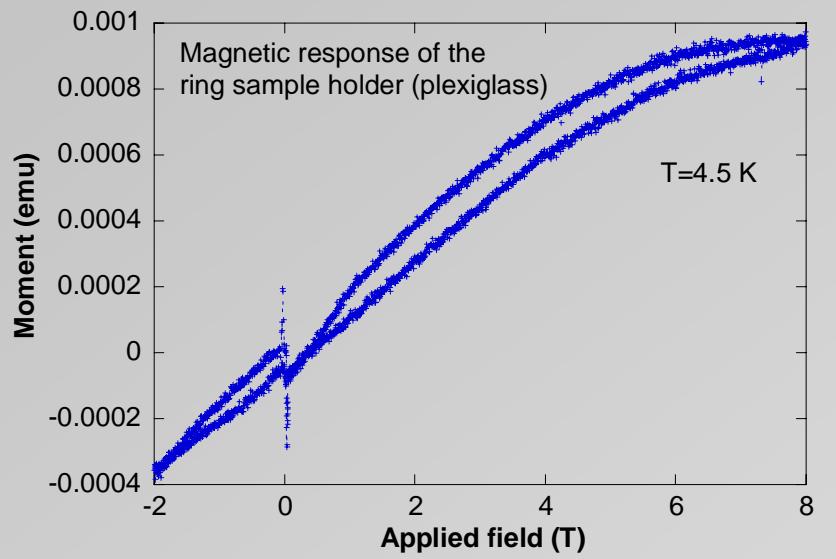


In practical measurement for a NbTi wire we expect a symmetric curve with respect the x axis. Of course our measure includes the sample holder.

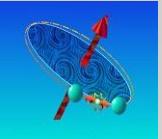




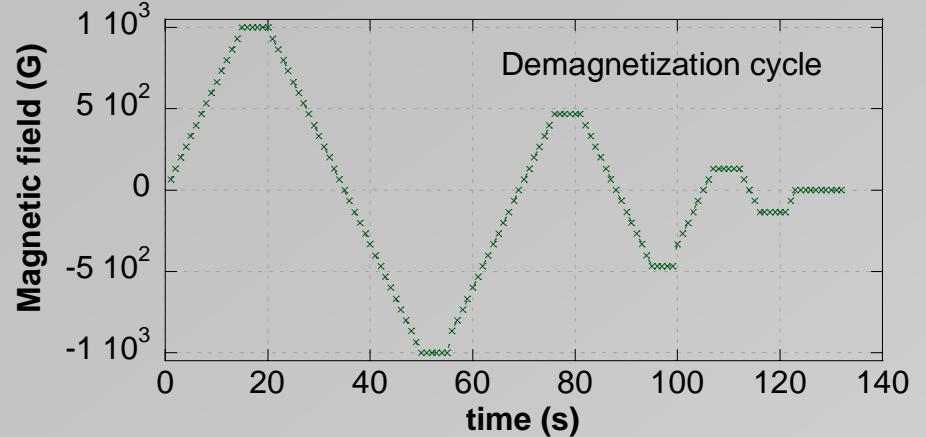
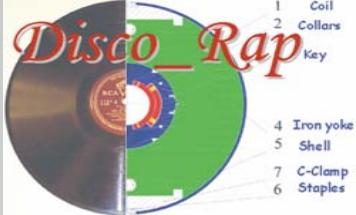
Wire + sample holder →



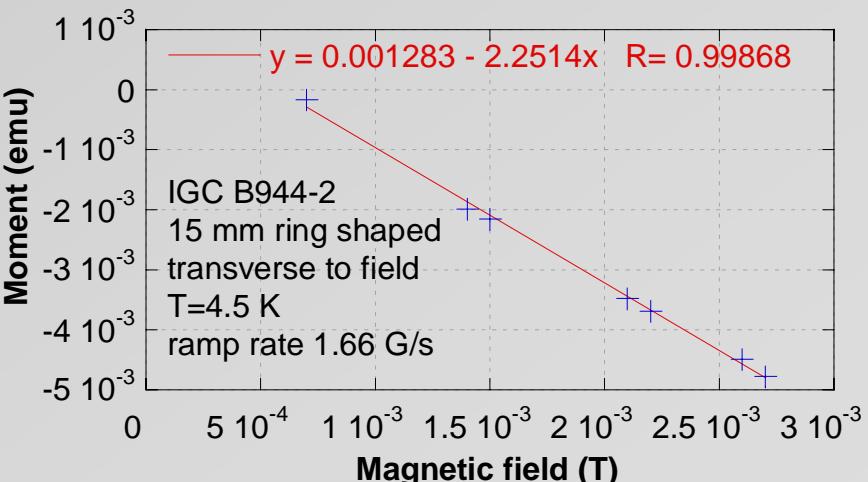
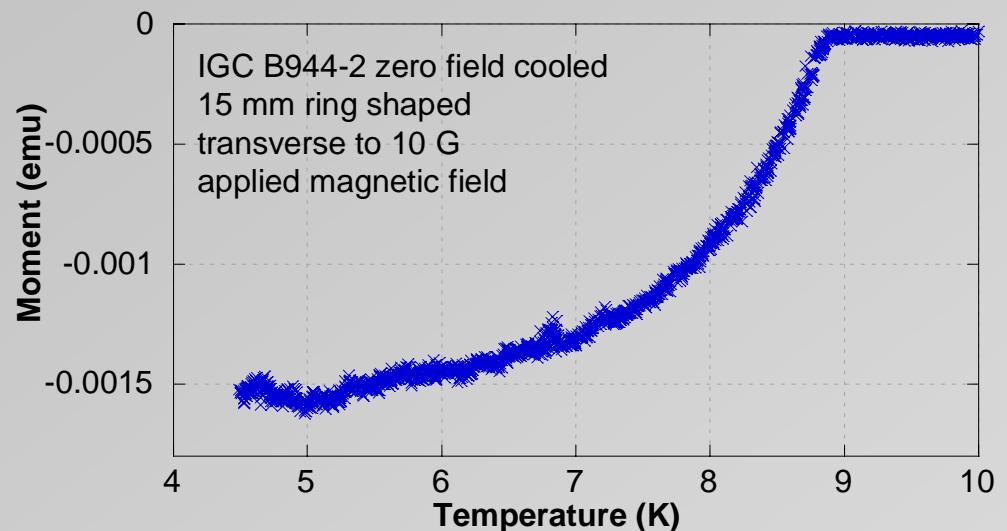
← sample holder



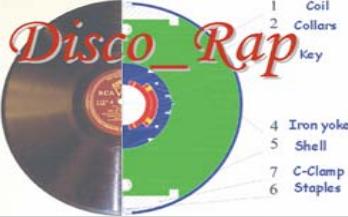
The shielded volume low field measure



To achieve low field the magnet itself has to be demagnetized



Limits arising from very thin filaments



When using standard $V_{\text{shield}} = m/(\chi_0 B)$ the ratio λ/d_{fil} should be negligible to avoid large errors: e.g. for sample SL 86969 ($d_{\text{fil}}=2.5 \mu\text{m}$) in 10 G we had $V_{\text{shield}}=0.753 \text{ mm}^3$ while the wire section is 2.87 mm^3 and $\alpha=1.55$ thus $V_{\text{nominal}}= 1.131 \text{ mm}^3$.

For Luvata 1 wire we found $SC=0.45 \text{ mm}^3$ to be compared to $SC_{\text{nominal}}= 0.64 \text{ mm}^3$

χ_0 must be corrected according to R/λ ratio but $\lambda(T)$ has to be determined

$$\chi_0 = - \left[1 - \frac{2\lambda}{R} \frac{\frac{1}{\pi} \int_0^{\pi} e^{R/\lambda \cos t} \cos t dt}{\frac{1}{\pi} \int_0^{\pi} e^{R/\lambda \cos t} dt} \right]$$



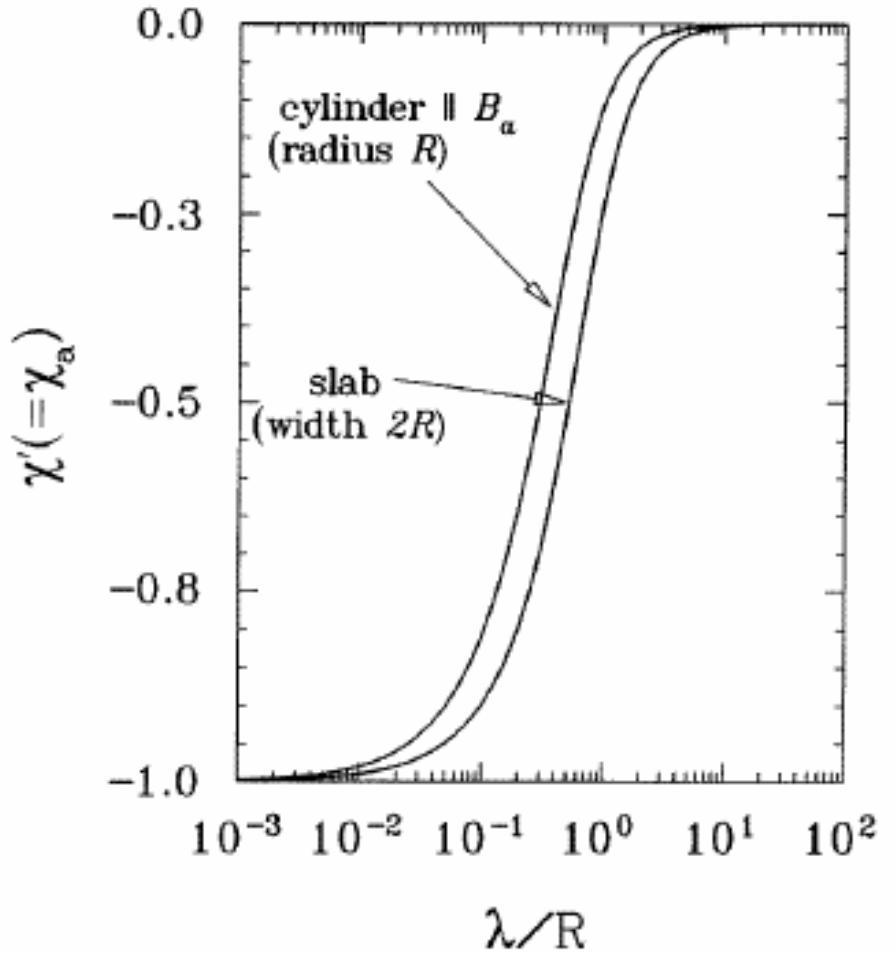
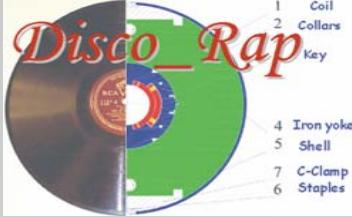
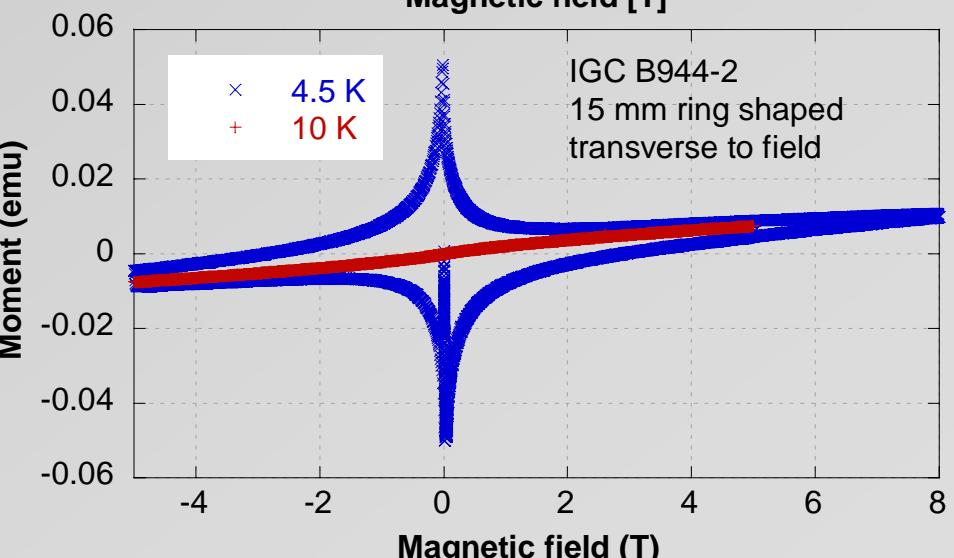
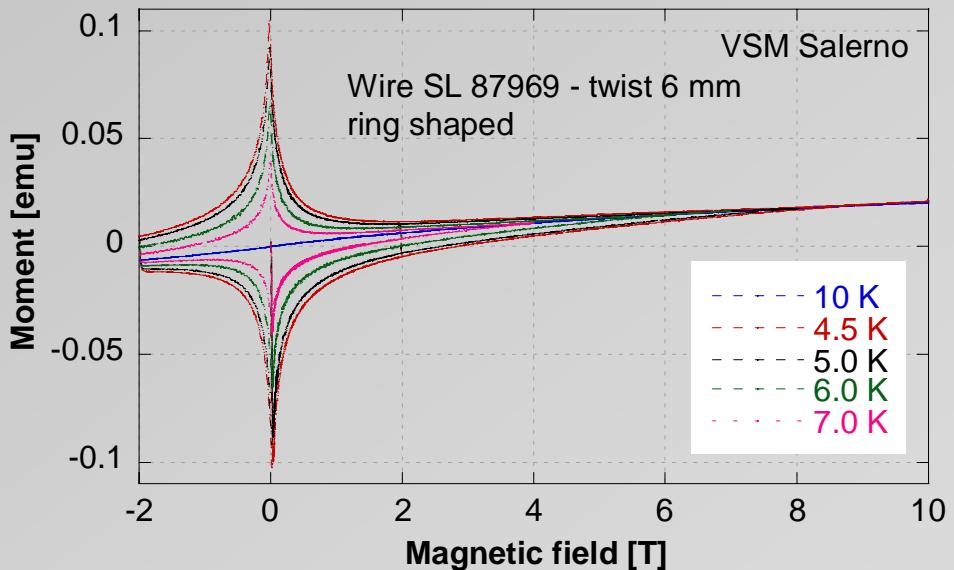
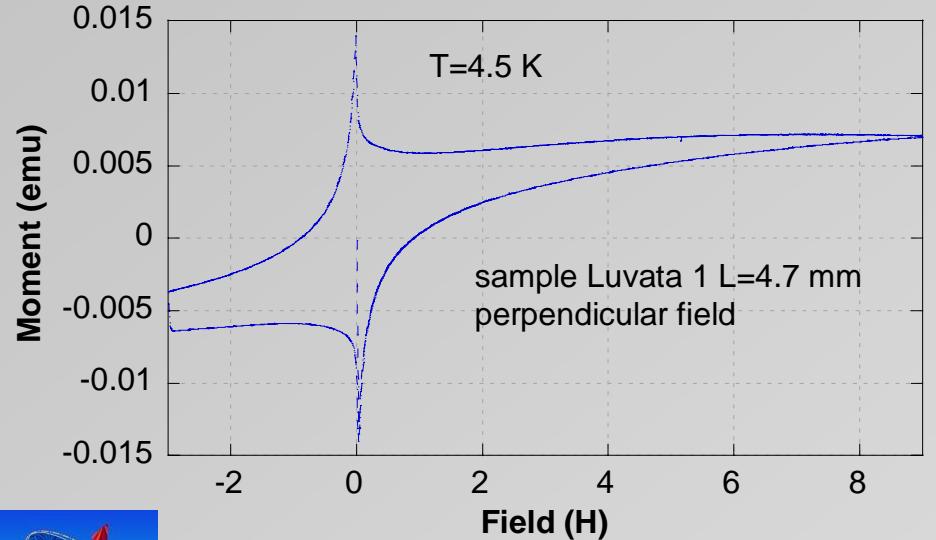
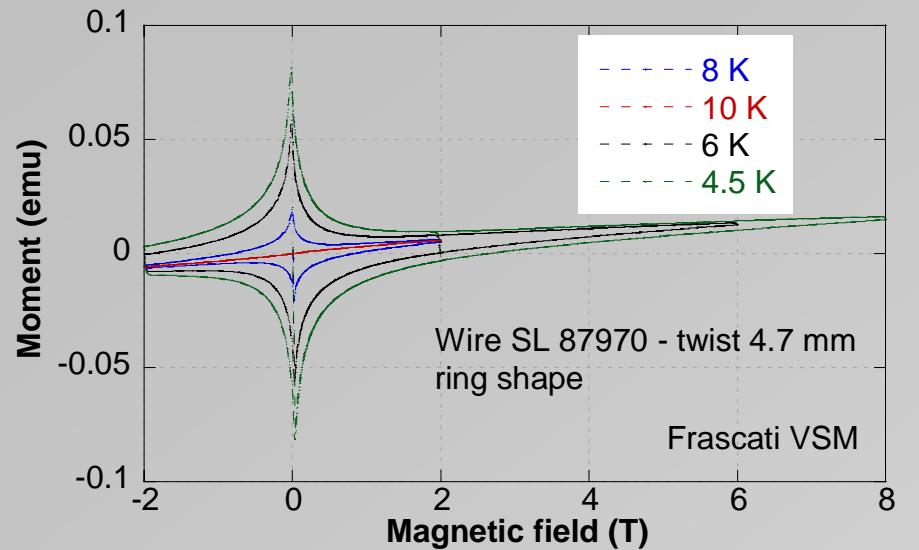
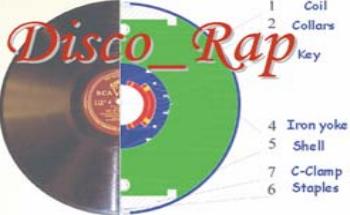


Figure 5. AC susceptibility in the regime of reversible screening.

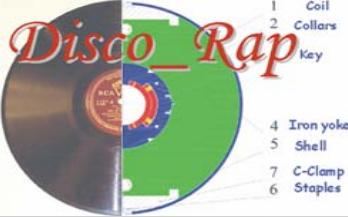
Fedor Gomory
SuST **10** (1997) 523–542



The presence of CuMn



Deconvolution of contributions



It is not a simple task, as the signal comes from:

1. the reversible magnetization of the specimen
2. the paramagnetism of Mn
3. the sample holder contribution

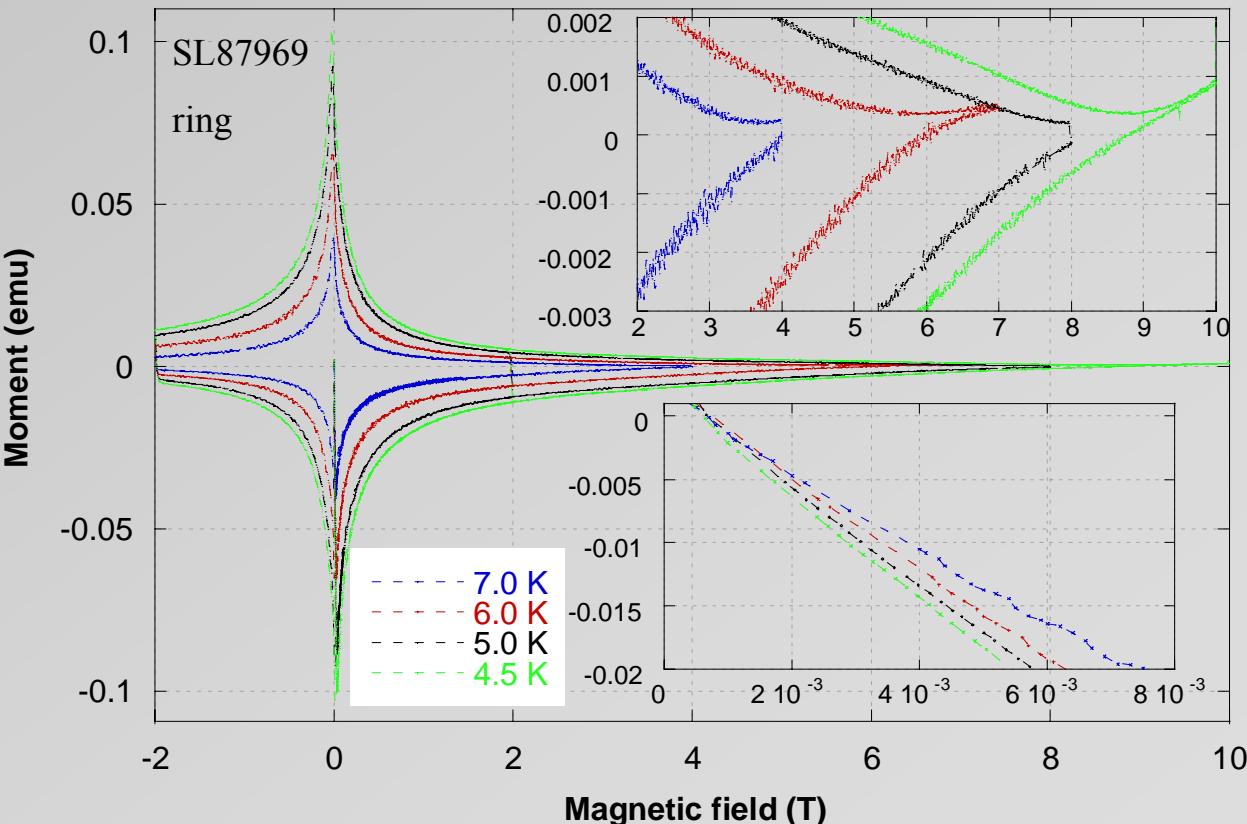
At $H_{c1} \ll H < H_{c2}$

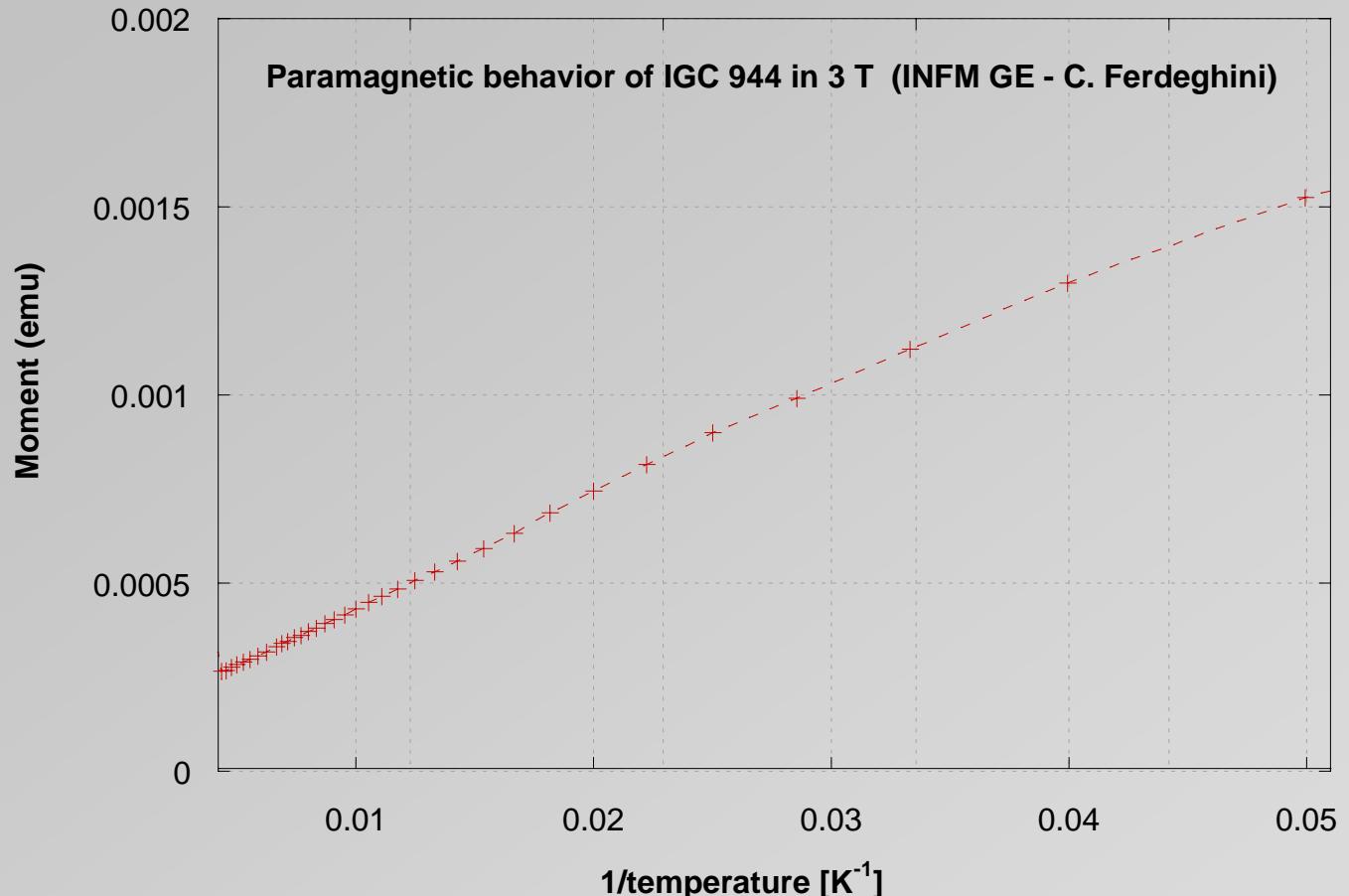
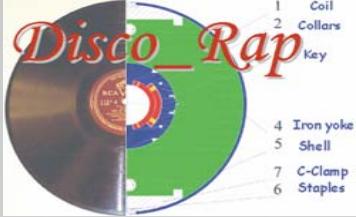
$$\frac{dM}{dH} = \frac{1}{4\pi\beta_a(2\kappa^2(T)-1)}$$

$$\kappa = \frac{1}{\sqrt{2}} \frac{H_{c2}}{H_c}$$

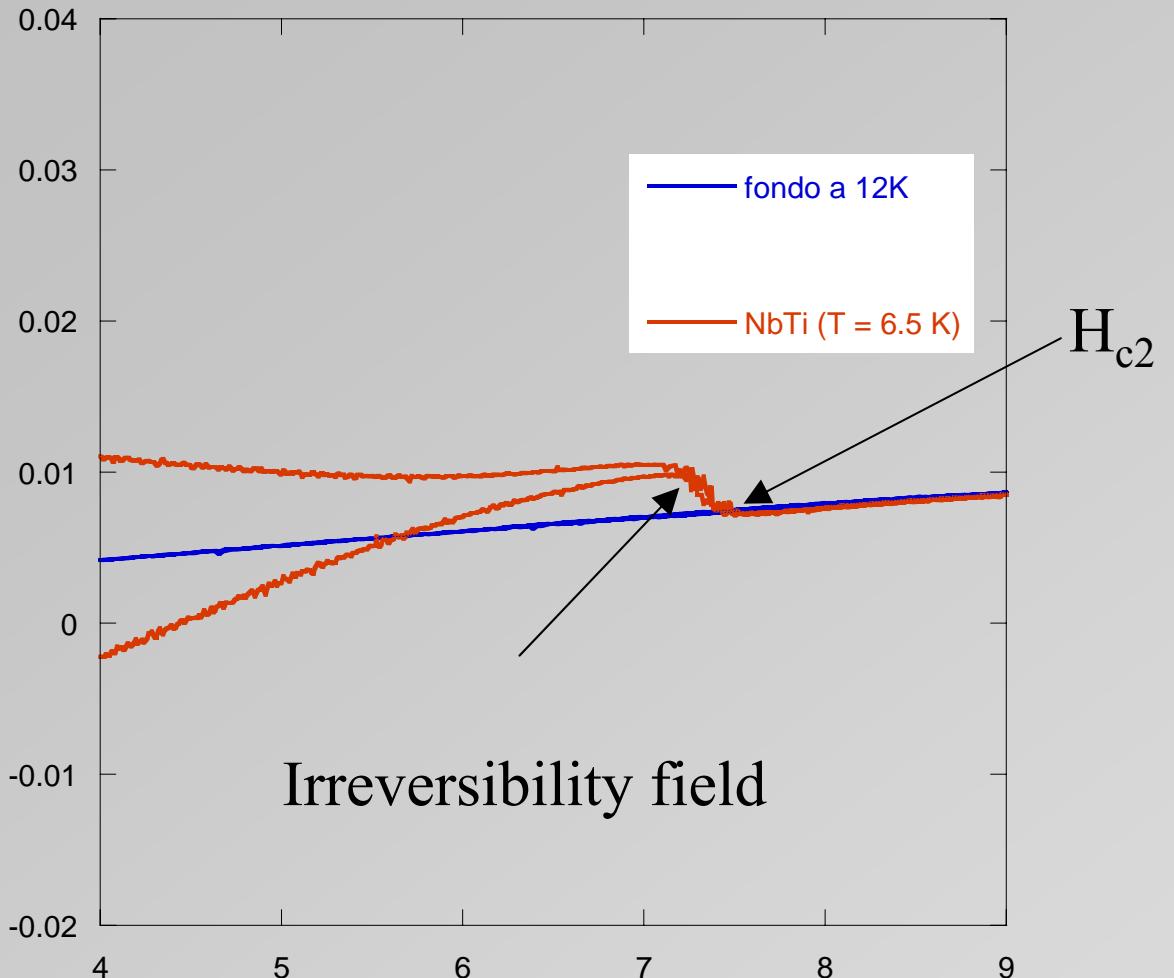
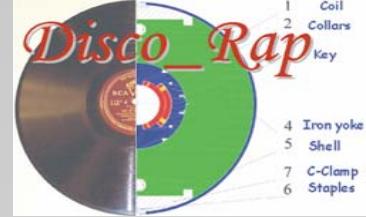
At $H \ll H_{c1}$

dM/dH vs T from ZFC may be used to estimate $\lambda(T)$



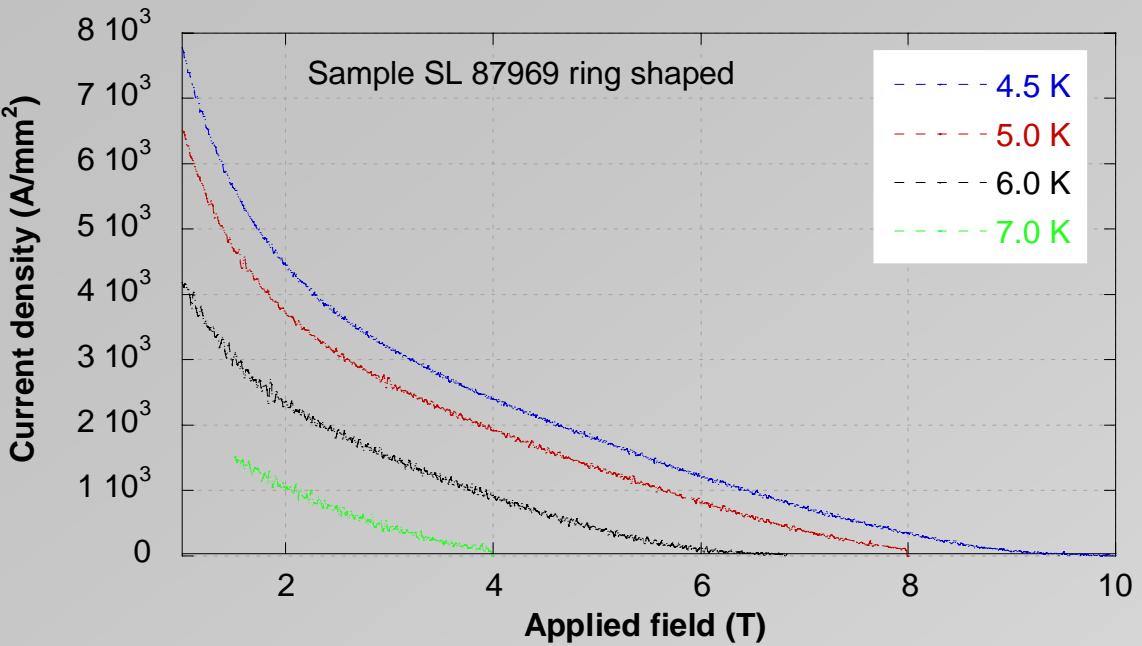
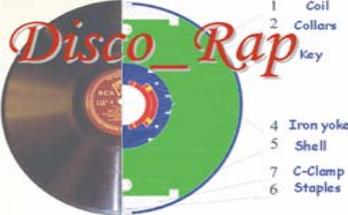


Other fundamental properties



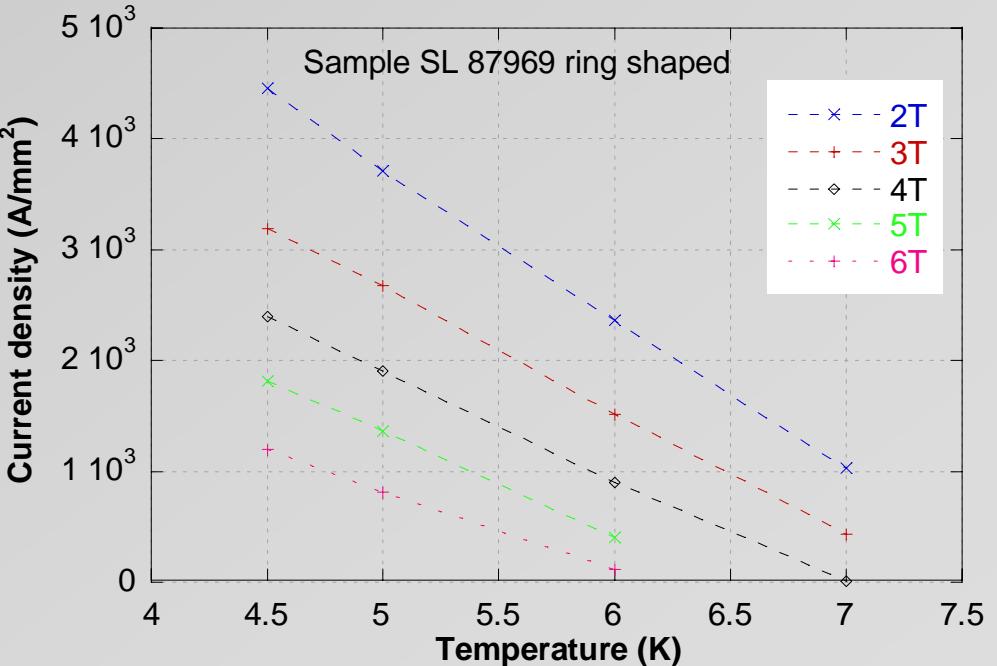
Luvata NbTi wire,
 $6 \mu\text{m}$ filaments, 18 mm ring shape measured at 6.5K



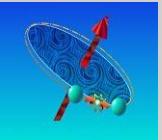
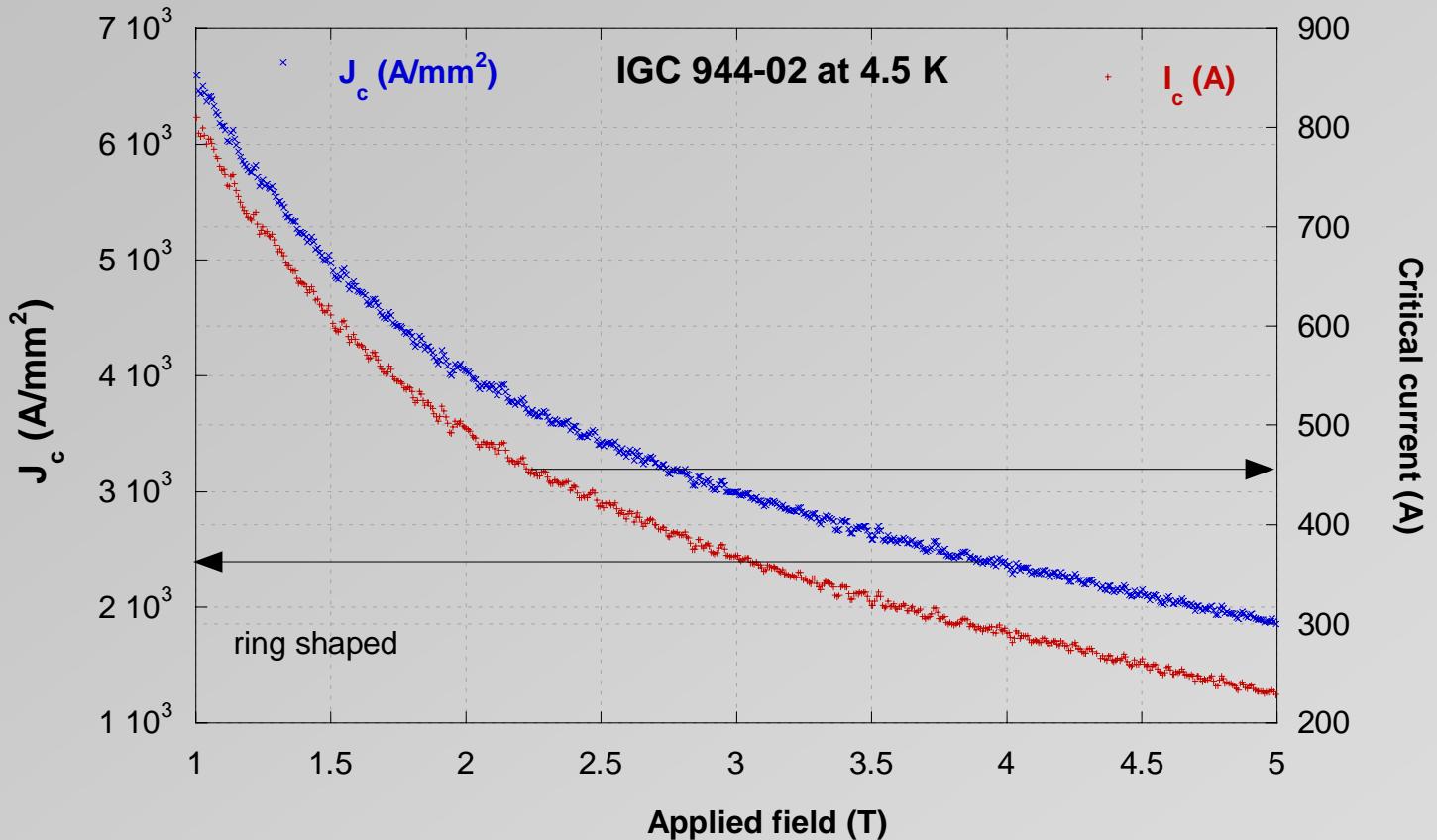
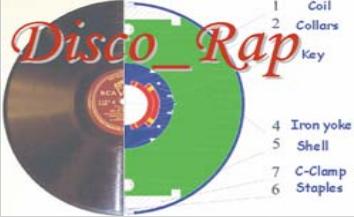


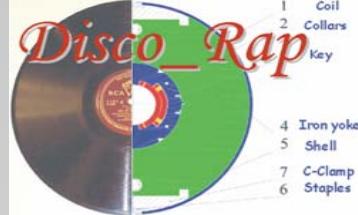
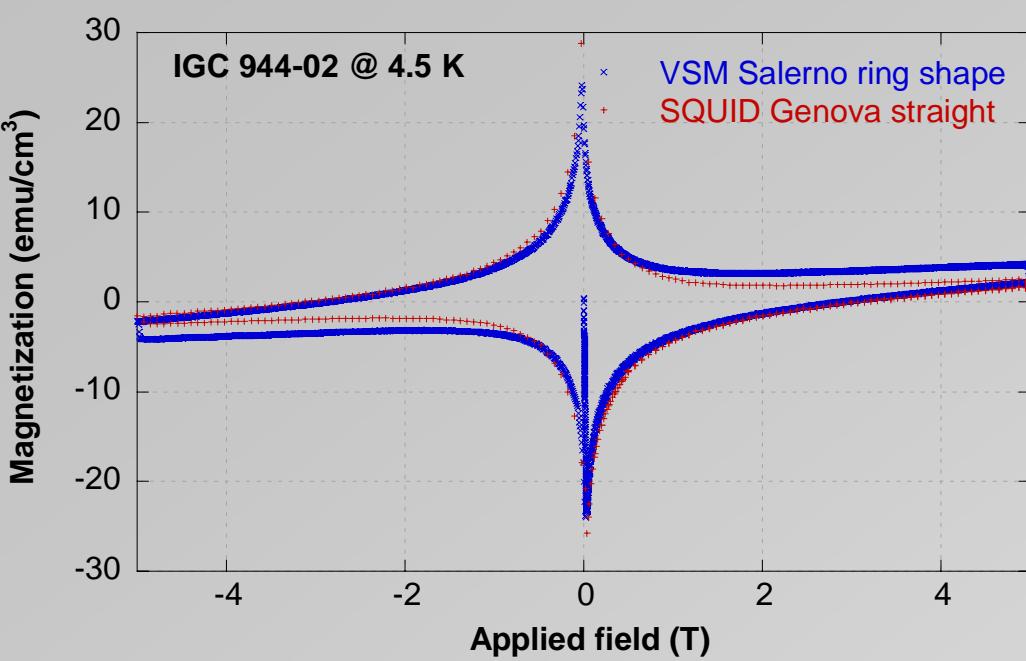
SL 87969 wire

Critical current density as derived from magnetization measurements in a 17mm long wire ring shaped

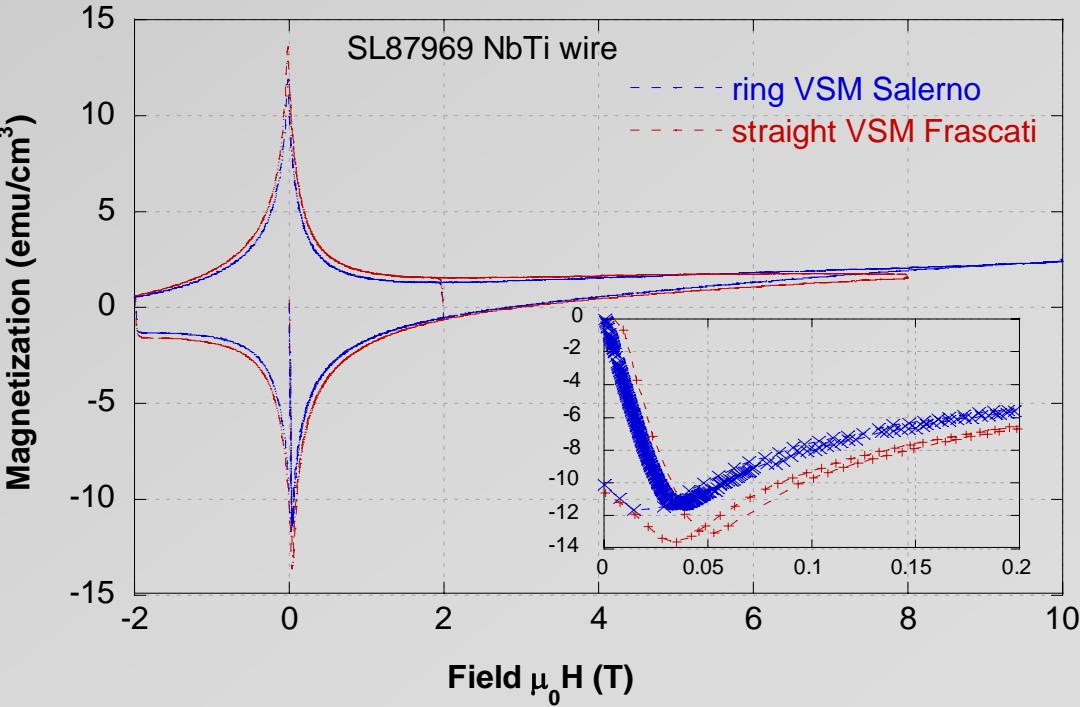


Critical currents in IGC B944-03

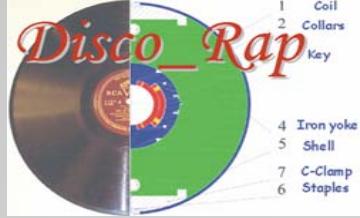




Comparison among different magnetometers



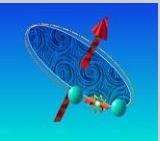
Losses & filaments coupling



$$Q_{mag} = \oint m(B) dB$$

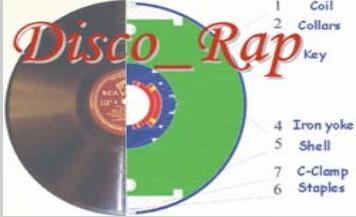
In the IGC 944 wire we have $Q_{mag} \approx 1.23 \cdot 10^{-1}$ mJ in a ± 5 T cycle

or $Q_{mag} \approx 58.5$ mJ/cm³ in a ± 5 T cycle



Cycles 4±0.5T at different ramp rates

(Frascati VSM only)



Luvata 6 μm NbTi wire

$$Q_c = \frac{B_m^2}{2\mu_0} \frac{8\tau}{T_m} \quad (\text{cfr. Wilson 8.53})$$

$n\tau = 11.2$ ms @ 4.8K

$n\tau = 10.7$ ms @ 6.0K

($n\tau = 6.8$ ms @ 5.0K,
ref. Gislon *et al.*, IEEE 2003)

