Quantum Mechanics: SERP-CHEM Final written test: 11/1/2013

Problem 1

A 1-dimensional harmonic oscillator with frequency ω is, at the time t = 0, in the state

$$\psi_0 = \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle + \frac{1}{\sqrt{5}} |3\rangle$$
 (1)

where $|n\rangle$, with n = 0, 1, 2, ... are the normalized energy eigenstates with eigenvalues $E_n = \hbar \omega \left(n + \frac{1}{2}\right)$.

- a) What is the average value of the energy at the time t > 0? (Points 5%)
- b) What is the average value of \hat{x}^2 on this state at t > 0? (Points 25%)
- c) What is the probability that at the time t > 0 the system is still in the state ψ_0 ? (Points 20%)

Problem 2

A free uni-dimensional particle of mass m is confined to the interval [0, L] of the real line. The particle is, at the time t = 0, in the state described by the (non-normalized) wave function

$$\psi_0(x) = \sin \frac{\pi x}{L} + \sin \frac{2\pi x}{L} \tag{2}$$

- a) What is the average energy of this state at time t > 0? (Points 5%)
- b) What is the probability that at the time t > 0 the particle is in the region $0 < x < \frac{L}{4}$? (Points 15%)
- c) Compute the probability that at the time t > 0 the particle be in the state described by the wave function

$$\psi_1(x) = \sin \frac{\pi x}{L} - \sin \frac{2\pi x}{L} \tag{3}$$

(Points 10%)

Problem 3

Find the *normalized* radial wave function of the bound state of the hydrogenoid with angular momentum l = 1 and lowest possible energy, by looking for solutions of the radial Schrödinger equation with l = 1 of the form

$$\chi(\rho) = C \,\rho^2 \,\mathrm{e}^{-A \,\rho} \tag{4}$$

where

$$\rho = \frac{Z r}{a_B} \tag{5}$$

Z is the atomic number, a_B the Bohr radius, r the radial coordinate, $\chi(r) = r R(r)$ with R(r) the radial wave function, A and C are two numerical constants to be determined.

- a) Determine A, C and the eigenvalue E of the energy corresponding to this eigenfunction. (Points 15%)
- b) Explain why the eigenvalue E is the lowest possible with the given angular momentum. (Points 5%)

Final written test: 11/1/2013: Solution

Problem 1

The norm squared of ψ_0 is

$$\langle \psi_0, \psi_0 \rangle = \frac{1}{2} + \frac{1}{2} + \frac{1}{5} = \frac{6}{5}$$
 (6)

Let us consider the normalized state

$$\tilde{\psi}_{0} = \frac{\sqrt{5}}{\sqrt{6}} \left(\frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle + \frac{1}{\sqrt{5}} |3\rangle \right) = = \frac{\sqrt{5}}{\sqrt{12}} |0\rangle + \frac{\sqrt{5}}{\sqrt{12}} |1\rangle + \frac{1}{\sqrt{6}} |3\rangle$$
(7)

The normalized state at time t is

$$\tilde{\psi}_{0}(t) = e^{-i\frac{\omega}{2}t} \left(\frac{\sqrt{5}}{\sqrt{12}} |0\rangle + e^{-i\omega t} \frac{\sqrt{5}}{\sqrt{12}} |1\rangle + e^{-3i\omega t} \frac{1}{\sqrt{6}} |3\rangle \right)$$
(8)

The average of the energy is independent of time and is

$$\langle \tilde{\psi}_0, \hat{H}\,\tilde{\psi}_0 \rangle = \hbar\,\omega\,\left(\frac{1}{2} \times \frac{5}{12} + \frac{3}{2} \times \frac{5}{12} + \frac{7}{2} \times \frac{1}{6}\right) = \frac{17}{12}\,\hbar\,\omega$$
 (9)

 \hat{x} and \hat{p} can be expressed in terms of a and a^{\dagger}

$$\hat{x} = i\lambda(a-a^{\dagger})$$
 $\hat{p} = \frac{\hbar}{2\lambda}(a+a^{\dagger})$ $\lambda \equiv \sqrt{\frac{\hbar}{2\,m\,\omega}}$ (10)

Therefore

$$\hat{x}\,\tilde{\psi}_{0}(t) = i\lambda\,\mathrm{e}^{-i\frac{\omega}{2}t}\left(-\frac{\sqrt{5}}{\sqrt{12}}\,|1\rangle + \frac{\sqrt{5}}{\sqrt{12}}\,\mathrm{e}^{-i\,\omega\,t}\,|0\rangle - \frac{\sqrt{5}\,\sqrt{2}}{\sqrt{12}}\,\mathrm{e}^{-i\,\omega\,t}\,|2\rangle + \frac{\sqrt{3}}{\sqrt{6}}\,\mathrm{e}^{-3\,i\,\omega\,t}\,|2\rangle - \frac{\sqrt{4}}{\sqrt{6}}\,\mathrm{e}^{-3\,i\,\omega\,t}\,|4\rangle\right) \tag{11}$$

Therefore the average of \hat{x}^2 is

$$\langle \tilde{\psi}_{0}(t), \hat{x}^{2} \, \tilde{\psi}_{0}(t) \rangle = \langle \hat{x} \, \tilde{\psi}_{0}, \hat{x} \, \tilde{\psi}_{0} \rangle = \lambda^{2} \left(\frac{5}{12} + \frac{5}{12} + \frac{5}{6} \left| 1 - \frac{\sqrt{3}}{\sqrt{5}} e^{-2i\omega t} \right|^{2} + \frac{4}{6} \right) = \lambda^{2} \left[\frac{3}{2} + \frac{5}{6} \left(1 + \frac{3}{5} - \frac{2\sqrt{3}}{\sqrt{5}} \cos(2\omega t) \right) \right] = \lambda^{2} \left[\frac{17}{6} - \frac{\sqrt{5}}{\sqrt{3}} \cos(2\omega t) \right]$$
(12)

The probability amplitude to be in the state ψ_0 at the time t

$$\langle \tilde{\psi}_0, \tilde{\psi}_0(t) \rangle = e^{\frac{-i\omega t}{2}} \left[\frac{5}{12} + \frac{5}{12} e^{-i\omega t} + \frac{1}{6} e^{-3i\omega t} \right]$$
 (13)

The corresponding probability is

$$P(t) = \left| \langle \tilde{\psi}_{0}, \tilde{\psi}_{0}(t) \rangle \right|^{2} = \left| \frac{5}{12} + \frac{5}{12} e^{-i\omega t} + \frac{1}{6} e^{-3i\omega t} \right|^{2} = \frac{1}{12^{2}} \left| 5 + 5 e^{-i\omega t} + 2 e^{-3i\omega t} \right|^{2} = \frac{1}{12^{2}} \left[54 + 50 \cos(\omega t) + 20 \cos(2\omega t) + 20 \cos(3\omega t) \right] = \frac{1}{72} \left[27 + 25 \cos(\omega t) + 10 \cos(2\omega t) + 10 \cos(3\omega t) \right]$$
(14)

Of course

$$P(0) = 1 \tag{15}$$

This function is periodic with period $T = \frac{2\pi}{\omega}$. Therefore at the times

$$t_n = \frac{2\pi n}{\omega} \tag{16}$$

the probability P(t) equals one

$$P(t_n) = 1 \tag{17}$$

Problem 2

The eigenvalues of the energy of the system are

$$E_n = \frac{n^2 \hbar^2 \pi^2}{2 m L^2} \qquad n = 1, 2, \dots$$
(18)

Since $\psi_0(x)$ is the sum of the two eigenfunctions with energies E_1 and E_2 the average energy is

$$\langle H \rangle = \frac{1}{2} \left(E_1 + E_2 \right) = \frac{5 \hbar^2 \pi^2}{4 m L^2}$$
 (19)

This is also the average energy at the time t, since the average of the energy does not depend on time.

Let us compute the norm of the wave function

$$\langle \psi_0, \psi_0 \rangle = \int_0^L dx \left(\sin \frac{\pi x}{L} + \sin \frac{2 \pi x}{L} \right)^2 = \int_0^L dx \left(\sin^2 \frac{\pi x}{L} + \sin^2 \frac{2 \pi x}{L} \right) = \frac{L}{2} + \frac{L}{2} = L$$
(20)

The *normalized* wave function at the time t is

$$\tilde{\psi}_0(x,t) = \frac{1}{\sqrt{L}} \left[e^{-i\frac{\hbar\pi^2 t}{2mL^2}} \sin\frac{\pi x}{L} + e^{-i\frac{2\hbar\pi^2 t}{mL^2}} \sin\frac{2\pi x}{L} \right]$$
(21)

The probability P(t) that particle is in the interval $[0, \frac{L}{3}]$ at the time t > 0 is

$$P(t) = \frac{1}{L} \int_{0}^{\frac{L}{4}} dx \left[\sin^{2} \frac{\pi x}{L} + \sin^{2} \frac{2\pi x}{L} + 2 \cos \frac{3\hbar\pi^{2} t}{2mL^{2}} \sin \frac{\pi x}{L} \sin \frac{2\pi x}{L} \right] = \\ = \int_{0}^{\frac{1}{4}} dy \left[\sin^{2}(\pi y) + \sin^{2}(2\pi y) + 2 \cos \frac{3\hbar\pi^{2} t}{2mL^{2}} \sin(\pi y) \sin(2\pi y) \right] = \\ = \left[\frac{1}{8} - \frac{1}{4\pi} + \frac{1}{8} + 2 \cos \frac{3\hbar\pi^{2} t}{2mL^{2}} \left(\frac{1}{3\sqrt{2}\pi} \right) \right] = \\ = \frac{1}{4} - \frac{1}{4\pi} + \frac{\sqrt{2}}{3\pi} \cos \frac{3\hbar\pi^{2} t}{2mL^{2}}$$
(22)

The normalized wave function corresponding to $\psi_1(x)$ is

$$\tilde{\psi}_1(x) = \frac{1}{\sqrt{L}} \left(\sin \frac{\pi x}{L} - \sin \frac{2\pi x}{L} \right)$$
(23)

The amplitude probability that at the time t the particle is in the state ψ_1 is

$$A_1(t) = \langle \tilde{\psi}_1, \tilde{\psi}_0(t) \rangle = \frac{1}{2} e^{-i\frac{\hbar\pi^2 t}{2mL^2}} - \frac{1}{2} e^{-i\frac{2\hbar\pi^2 t}{mL^2}}$$
(24)

The probability is

$$P_1(t) = |A_1(t)|^2 = \frac{1}{2} \left(1 - \cos \frac{3\hbar\pi^2 t}{2mL^2} \right) = \sin^2 \frac{3\hbar\pi^2 t}{2mL^2}$$
(25)

Problem 3

The Schrödinger equation for $\chi_{E,l}(\rho)$ of angular momentum l , written in adimensional variables is

$$-\chi_{E,l}''(\rho) + \left(-\frac{2}{\rho} + \frac{l(l+1)}{\rho^2}\right)\chi_{E,l}(\rho) = 2\,\epsilon\,\chi_{E,l}(\rho)$$
(26)

where ϵ is an adimensional quantity proportional to the energy

$$\epsilon = \frac{E}{\frac{Z^2 \hbar^2}{m \, a_B^2}} \tag{27}$$

For l = 1 this is

$$\frac{\chi_{E,1}''}{\chi_{E,1}} = -\frac{2}{\rho} + \frac{2}{\rho^2} - 2\epsilon$$
(28)

Using

$$\chi_{E,1}(\rho) = C \,\rho^2 \,\mathrm{e}^{-A\,\rho} \tag{29}$$

we obtain

$$A^{2} + \frac{2}{\rho^{2}} - \frac{4A}{\rho} = -\frac{2}{\rho} + \frac{2}{\rho^{2}} - 2\epsilon$$
(30)

Therefore we arrive to the equations

$$4A = 2 \qquad A^2 = -2\epsilon \tag{31}$$

whose solution is

$$A = \frac{1}{2} \qquad \epsilon = -\frac{1}{8} \tag{32}$$

This value of ϵ corresponds to the energy

$$E = -\frac{Z^2 \hbar^2}{2 m a_B^2} \frac{1}{2^2}$$
(33)

which is indeed the Bohr level with n = 2.

The normalization factor C is determined from

$$\int_0^\infty d\rho \,\chi_{E,l}^2(\rho) = \frac{Z}{a_B} \tag{34}$$

From the expression (29) we obtain

$$C^2 \int_0^\infty d\rho \,\rho^4 \,\mathrm{e}^{-\rho} = C^2 \,24 = \frac{Z}{a_B} \tag{35}$$

Hence

$$C = \sqrt{\frac{Z}{a_B}} \frac{1}{2\sqrt{6}} \tag{36}$$

Since the wave function has no nodes for $\rho > 0$ it must correspond to the lowest energy eigenvalue with the given l.

Quantum Mechanics: SERP-CHEM Final written test: 28/1/2013

Problem 1

A quantum system is described by a 3-dimensional space of states. Let $\{|1\rangle, |2\rangle, |3\rangle\}$ be an orthonormal base and let the Hamiltonian of the system in this base be

$$H = \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix}$$
(37)

where ϵ and η are real parameters.

• a) Show that the operator represented in the given basis by the matrix

$$T = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$$
(38)

commutes with the Hamiltonian H. Does T correspond to an observable? (Points 10%)

- b) Compute the eigenvectors and eigenvalues of T. (Points 20%)
- c) Using (b), determine the eigenvectors and eigenvalues of H. (Points 25%)
- d) Suppose that at the time t = 0 the system is in the state $\psi_0 = |1\rangle$. What is the probability that the system is again in the same state $|1\rangle$ after a time t? (Points 20%)

Problem 2

An harmonic oscillator of frequency ω is, at time t = 0, in the state

$$\psi_0 = |2\rangle + 2|3\rangle \tag{39}$$

- a) Compute the average of \hat{x} and \hat{x}^2 at the time t. (Points 15%)
- b) Write the normalized wave function of the state at the time t in the Schrödinger representation. (Points 10%)

Final written test: 29/1/2013: Solution

Problem 1

a)

$$HT = \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} = \begin{pmatrix} \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \\ \epsilon & \eta & \eta \end{pmatrix}$$
$$TH = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix} = \begin{pmatrix} \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \\ \epsilon & \eta & \eta \end{pmatrix}$$
(40)

T is not hermitian, and thus it does not correspond to an observable. b)

The eigenvalues of T are given by the secular equation

$$P(t) = \det \begin{pmatrix} -t & 1 & 0\\ 0 & -t & 1\\ 1 & 0 & -t \end{pmatrix} = -t(t^2) - 1(-1) = 0$$
(41)

Hence

$$t_1 = e^{\frac{2\pi}{3}i}$$
 $t_2 = t_1^2 = t_1^{-1}$ $t_3 = t_1^3 = 1$ (42)

The eigenvectors relative to t_k ,

$$v_k = \begin{pmatrix} x \\ y \\ z \end{pmatrix} \qquad \text{with } k = 1, 2, 3, \tag{43}$$

satisfy the equations

$$y = t_k x \qquad z = t_k y = t_k^2 z \tag{44}$$

Hence the normalized eigenvectors are

$$v_k = \frac{1}{\sqrt{3}} \begin{pmatrix} 1\\t_k\\t_k^2\\t_k^2 \end{pmatrix} \qquad \text{with } k = 1, 2, 3 \tag{45}$$

or, explicitly

$$v_1 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1\\t_1\\t_1^{-1} \end{pmatrix} \qquad v_2 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1\\t_1^{-1}\\t_1 \end{pmatrix} \qquad v_3 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1\\1\\1 \end{pmatrix}$$
(46)

c) Consider:

$$T v_k = t_k v_k \tag{47}$$

Then

$$HT v_k = t_k H v_k \tag{48}$$

Since H and T commute, we obtain

$$T(Hv_k) = t_k (Hv_k) \tag{49}$$

Therefore $H v_k$ is an eigenvector of T with eigenvalue t_k . Since t_k is nondegenerate and its only eigenvector is v_k , $H v_k$ must be proportional to v_k :

$$H v_k = E_k v_k \tag{50}$$

Thus, v_k are simultaneous eigenvectors of H and T. We can compute the eigenvalues E_k by acting on v_k with H

$$H v_1 = \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix} \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ t_1 \\ t_1^{-1} \end{pmatrix} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ t_1 \\ t_1^{-1} \end{pmatrix} (\epsilon + \eta(t_1 + t_1^{-1}))$$
(51)

Thus

$$E_1 = \epsilon + \eta \left(t_1 + t_1^{-1} \right) = \epsilon + 2 \eta \cos \frac{2\pi}{3} = \epsilon - \eta$$
 (52)

Analogously

$$Hv_2 = \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix} \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ t_1^{-1} \\ t_1 \end{pmatrix} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ t_1^{-1} \\ t_1 \end{pmatrix} (\epsilon + \eta(t_1^{-1} + t_1))$$
(53)

from which

$$E_2 = E_1 = \epsilon - \eta \tag{54}$$

Finally

$$H v_3 = \begin{pmatrix} \epsilon & \eta & \eta \\ \eta & \epsilon & \eta \\ \eta & \eta & \epsilon \end{pmatrix} \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} (\epsilon + 2\eta)$$
(55)

and

$$E_3 = \epsilon + 2\eta \tag{56}$$

d) Let us decompose the state $|1\rangle$ in the base $\{v_1, v_2, v_3\}$:

$$|1\rangle = v_1 \langle v_1 | 1 \rangle + v_2 \langle v_2 | 1 \rangle + v_3 \langle v_3 | 1 \rangle = = \frac{1}{\sqrt{3}} (v_1 + v_2 + v_3)$$
(57)

Hence the state at the time t is

$$\psi(t) = \frac{1}{\sqrt{3}} e^{-\frac{iE_1t}{\hbar}} \left(v_1 + v_2 + e^{-\frac{i(E_3 - E_1)t}{\hbar}} v_3 \right) = \frac{1}{\sqrt{3}} e^{-\frac{iE_1t}{\hbar}} \left(v_1 + v_2 + e^{-\frac{i3\eta t}{\hbar}} v_3 \right)$$
(58)

The amplitude for this state to be into the state $|1\rangle$ is

$$A(t) = \langle 1|\psi(t)\rangle = \frac{1}{3} e^{-\frac{iE_1 t}{\hbar}} \left(2 + e^{-\frac{i3\eta t}{\hbar}}\right)$$
(59)

The corresponding probability

$$P(t) = |A(t)|^2 = \frac{1}{9} \left(5 + 4 \cos \frac{3 \eta t}{\hbar} \right)$$
(60)

Problem 2

a)

Let us write \hat{x} in terms of a and a^{\dagger}

$$\hat{x} = i \lambda (a - a^{\dagger}) \qquad \lambda = \sqrt{\frac{\hbar}{2 m \omega}}$$
 (61)

The system is at the time t in the state described by the normalized vector

$$\psi(t) = \frac{1}{\sqrt{5}} e^{-i\frac{5}{2}\omega t} \left(|2\rangle + 2 e^{-i\omega t} |3\rangle \right)$$
(62)

Hence

$$\hat{x}\,\psi(t) = \frac{i\,\lambda}{\sqrt{5}}\,\mathrm{e}^{-i\,\frac{5}{2}\,\omega\,t}\left(\sqrt{2}\,|1\rangle - \sqrt{3}\,|3\rangle + 2\,\sqrt{3}\,\mathrm{e}^{-i\,\omega\,t}\,|2\rangle - 4\,\mathrm{e}^{-i\,\omega\,t}\,|4\rangle\right) \quad (63)$$

Therefore

$$\langle \psi(t) | \hat{x} | \psi(t) \rangle = \frac{i\lambda}{5} \left(2\sqrt{3} e^{-i\omega t} - 2\sqrt{3} e^{i\omega t} \right) = \frac{4\sqrt{3}\lambda}{5} \sin \omega t \tag{64}$$

and

$$\langle \psi(t) | \hat{x}^2 | \psi(t) \rangle = \langle \hat{x} \, \psi(t), \hat{x} \, \psi(t) \rangle = \frac{\lambda^2}{5} \left(2 + 3 + 12 + 16 \right) = \frac{33 \, \lambda^2}{5} \tag{65}$$

b) The state at the time t is

$$\psi(t) = \frac{1}{\sqrt{5}} e^{-i\frac{5}{2}\omega t} \left(|2\rangle + 2 e^{-i\omega t} |3\rangle \right)$$
(66)

The normalized wave function of the ground state in the Schrödinger representation is

$$\psi_0(x) = \frac{1}{(2\pi)^{\frac{1}{4}} \lambda^{\frac{1}{2}}} e^{-\frac{x^2}{4\lambda^2}} = \frac{1}{(2\pi)^{\frac{1}{4}} \lambda^{\frac{1}{2}}} e^{-\frac{y^2}{4}} \qquad y \equiv \frac{x}{\lambda}$$
(67)

Therefore the normalized wave function of the first excited state is

$$\psi_1(x) = -i\left(\lambda \frac{d}{dx} - \frac{1}{2} \frac{x}{\lambda^2}\right)\psi_0(x) = -i\left(\frac{d}{dy} - \frac{1}{2}y\right)\psi_0 = -i\left(-\frac{1}{2}y - \frac{1}{2}y\right)\psi_0 = iy\psi_0$$
(68)

and

$$\psi_{2}(x) = \frac{1}{\sqrt{2}} \left(\frac{d}{dy} - \frac{1}{2} y \right) y \psi_{0} = \frac{1}{\sqrt{2}} \left(-y^{2} \psi_{0} + \psi_{0} \right) = \frac{1}{\sqrt{2}} \left(1 - y^{2} \right) \psi_{0}$$

$$\psi_{3}(x) = \frac{-i}{\sqrt{6}} \left(\frac{d}{dy} - \frac{1}{2} y \right) \left(\left(1 - y^{2} \right) \psi_{0} \right) =$$

$$= \frac{-i}{\sqrt{6}} \left[-\left(1 - y^{2} \right) y \psi_{0} - 2 y \psi_{0} \right] = \frac{-i}{\sqrt{6}} \left[y^{3} - 3 y \right] \psi_{0}$$
(69)

Therefore the normalized wave-function of the state at the time t is

$$\psi(x,t) = \frac{1}{\sqrt{5}} e^{-i\frac{5}{2}\omega t} \left(\frac{1}{\sqrt{2}} \left(1-y^2\right) - \frac{2ie^{-i\omega t}}{\sqrt{6}} \left(y^3-3y\right)\right)\psi_0 = \\ = \frac{e^{-i\frac{5}{2}\omega t}}{\sqrt{5} \left(2\pi\right)^{\frac{1}{4}} \lambda^{\frac{1}{2}}} \left(\frac{1}{\sqrt{2}} \left(1-y^2\right) - \frac{2ie^{-i\omega t}}{\sqrt{6}} \left(y^3-3y\right)\right) e^{-\frac{y^2}{4}}$$
(70)

Quantum Mechanics: SERP-CHEM Final written test: 12/7/2013

Problem 1

A quantum system is described by a 2-dimensional space of states. Let $\{|1\rangle, |2\rangle\}$ be an orthonormal base and let the Hamiltonian of the system in this base be

$$H = \begin{pmatrix} a & b \\ b^* & a \end{pmatrix} \tag{71}$$

where a and b are *complex* numbers.

- a) What are the restriction on the numbers a and b for H to be hermitian? (Points 10%)
- b) Compute the eigenvalues and the normalized eigenvectors of *H*, when *H* is hermitian. (Points 20%)
- c) Suppose that the system be at the time t = 0 in the state $|1\rangle$:

$$\psi(0) = |1\rangle \tag{72}$$

Compute the probability at the time t that the system be in the state $|2\rangle$. (Points 25%)

Problem 2

An otherwise free one-dimensional particle of mass m is confined to a box of size L. Let us denote by $|E_n\rangle$ the eigenstates of the energy, with $n = 1, 2, 3, \ldots$

- a) What is the normalized wave function of the fundamental state in the *momentum representation*? (Points 20%)
- b) What is the normalized distribution probability of the momentum when the particle is in the fundamental state? (Points 10%)
- c) What is the average of \hat{p}^2 in the fundamental state? (Points 10%)
- d) Suppose the particle is in the state $\psi(0) = |E_1\rangle + 3i |E_2\rangle$ at the time t = 0. What is the smallest probability that the particle be in the state $\psi(0)$ at any time t > 0? (Points 15%)

Final written test: 12/7/2013: Solution

Problem 1

• a) The hermitian conjugate of H is

$$H^{\dagger} = \begin{pmatrix} a^* & b\\ b^* & a^* \end{pmatrix} \tag{73}$$

Therefore the condition for H to be hermitian is

$$a = a^* \tag{74}$$

i.e. a real.

• b) The eigenvalues are determined by the equation

$$0 = \det(H - E \mathbb{I}) = \det\begin{pmatrix} a - E & b \\ b^* & a - E \end{pmatrix} = = E^2 - 2 a E + a^2 - |b|^2$$
(75)

that is

$$E = E_{\pm} = a \pm |b| \tag{76}$$

The corresponding eigenvectors

$$|E_{\pm}\rangle = x_{\pm}|1\rangle + y_{\pm}|2\rangle \tag{77}$$

are determined by the equations

$$\begin{pmatrix} a - E_{\pm} & b \\ b^* & a - E_{\pm} \end{pmatrix} \begin{pmatrix} x_{\pm} \\ y_{\pm} \end{pmatrix} = 0$$
(78)

that is

$$(a - E_{\pm}) x_{\pm} = -b y_{\pm} \tag{79}$$

that is

$$\mp |b| x_{\pm} = -b y_{\pm} \Rightarrow x_{\pm} = \pm \frac{b}{|b|} y_{\pm} = \pm e^{i\theta_b} y_{\pm}$$
(80)

where θ_b is the argument of b

$$b = |b| e^{i\theta_b} \tag{81}$$

Thus the normalized eigenvectors are

$$|E_{\pm}\rangle = \frac{1}{\sqrt{2}} \left(\pm e^{i\theta_b} |1\rangle + |2\rangle \right) \tag{82}$$

• c) Let us decompose $|1\rangle$ in the basis $\{|E_+\rangle, |E_-\rangle\}$:

$$\psi(0) = |1\rangle = \frac{1}{\sqrt{2}} e^{-i\theta_b} \left(|E_+\rangle - |E_-\rangle \right)$$
(83)

The state at the time t is therefore

$$\psi(t) = \frac{1}{\sqrt{2}} e^{-i\theta_b} \left(e^{-i\frac{E_+}{\hbar}t} |E_+\rangle - e^{-i\frac{E_-}{\hbar}t} |E_-\rangle \right) =$$
$$= \frac{1}{\sqrt{2}} e^{-i\theta_b} e^{-i\frac{a}{\hbar}t} \left(e^{-i\frac{|b|}{\hbar}t} |E_+\rangle - e^{i\frac{|b|}{\hbar}t} |E_-\rangle \right)$$
(84)

The amplitude to be in the state $|2\rangle$ at the time t is therefore

$$A_{1\to2}(t) = \langle 2|\psi(t)\rangle = \frac{1}{\sqrt{2}} e^{-i\theta_b} e^{-i\frac{a}{\hbar}t} \left(e^{-i\frac{|b|}{\hbar}t} \langle 2|E_+\rangle + -e^{i\frac{|b|}{\hbar}t} \langle 2|E_-\rangle \right) =$$
$$= \frac{1}{2} e^{-i\theta_b} e^{-i\frac{a}{\hbar}t} \left(e^{-i\frac{|b|}{\hbar}t} - e^{i\frac{|b|}{\hbar}t} \right) =$$
$$= i e^{-i\theta_b} e^{-i\frac{a}{\hbar}t} \sin\frac{|b|}{\hbar} t$$
(85)

and the probability

$$P_{1\to 2}(t) = |A_{1\to 2}(t)|^2 = \sin^2 \frac{|b|}{\hbar} t$$
(86)

Remark: By changing basis, and using $|1'\rangle = e^{i\theta_b} |1\rangle$ and $|2'\rangle = |2\rangle$, we can always change the matrix H to a *real* matrix

$$H' = \begin{pmatrix} a & |b|\\|b| & a \end{pmatrix}$$
(87)

Therefore physical quantities cannot depend on the phase θ_b .

Problem 2

• a) The normalized wave function in the coordinate representation is

$$\psi_1(x) = \sqrt{\frac{2}{L}} \sin \frac{\pi x}{L} \tag{88}$$

if **x** is the coordinate in the interval [0, L]. The wave function in the momentum representation is

$$\tilde{\psi}_{1}(p) = \int_{0}^{L} \psi_{1}(x) e^{-i\frac{p}{\hbar}x} \frac{dx}{\sqrt{2\pi\hbar}} = \\ = \int_{0}^{L} \left[e^{i(\frac{\pi}{L} - \frac{p}{\hbar})x} - e^{-i(\frac{\pi}{L} + \frac{p}{\hbar})x} \right] \frac{dx}{2i\sqrt{\pi\hbar L}} = \\ = \left[\frac{e^{i(\frac{\pi}{L} - \frac{p}{\hbar})L} - 1}{\frac{\pi}{L} - \frac{p}{\hbar}} + \frac{e^{-i(\frac{\pi}{L} + \frac{p}{\hbar})L} - 1}{\frac{\pi}{L} + \frac{p}{\hbar}} \right] \frac{-1}{2\sqrt{\pi\hbar L}} = \\ = \frac{\frac{\pi}{L}}{\frac{\pi^{2}}{L^{2}} - \frac{p^{2}}{\hbar^{2}}} \frac{e^{-i\frac{p}{\hbar}L} + 1}{\sqrt{\pi\hbar L}} = \frac{e^{-i\frac{p}{\hbar}L} + 1}{\frac{\hbar^{2}\pi^{2}}{L^{2}} - p^{2}} \sqrt{\frac{\pi\hbar^{3}}{L^{3}}}$$
(89)

• b) The probability distribution of the momentum is therefore

$$|\tilde{\psi}_1(p)|^2 = \frac{\left|e^{-i\frac{p}{\hbar}L} + 1\right|^2}{\left(\frac{\hbar^2\pi^2}{L^2} - p^2\right)^2} \frac{\pi\hbar^3}{L^3} = \frac{\cos^2\left(\frac{p}{2\hbar}L\right)}{\left(\frac{\hbar^2\pi^2}{L^2} - p^2\right)^2} \frac{4\pi\hbar^3}{L^3}$$
(90)

• c) $\psi_1(x)$ is an eigenfunction of \hat{p}^2 , therefore the average of \hat{p}^2 on this state is its eigenvalue

$$\langle \psi_1, \hat{p}^2 \, \psi_1 \rangle = \frac{\hbar^2 \, \pi^2}{L^2} \tag{91}$$

• d) The normalized state at the time t is

$$\psi(t) = \frac{1}{\sqrt{10}} e^{-it\frac{E_1}{\hbar}} |E_1\rangle + \frac{3}{\sqrt{10}} i e^{-it\frac{E_2}{\hbar}} |E_2\rangle$$
(92)

The amplitude to be in the same state $\psi(0)$ at the time t is

$$\left\langle \frac{1}{\sqrt{10}} \psi(0), \psi(t) \right\rangle = \frac{1}{10} e^{-it\frac{E_1}{\hbar}} + \frac{9}{10} e^{-it\frac{E_2}{\hbar}}$$
(93)

The corresponding probability is

$$P(t) = \left| \left\langle \frac{1}{\sqrt{10}} \psi(0), \psi(t) \right\rangle \right|^2 = \frac{82}{100} + \frac{18}{100} \cos \frac{E_2 - E_1}{\hbar} t = = \frac{82}{100} + \frac{18}{100} \cos \frac{3\pi^2 \hbar t}{2 m L^2}$$
(94)

The smallest probability to be in the fundamental state is therefore

$$P_{min} = \frac{82}{100} - \frac{18}{100} = \frac{64}{100} \tag{95}$$

Quantum Mechanics: SERP-CHEM Final written test: 7/1/2014

Problem 1

An harmonic oscillator of frequency ω and mass m is in the state

$$\psi = |0\rangle + i |1\rangle$$

- a) Compute the *normalized* wave function of the state in the *momentum* representation. (Points 20%)
- b) Compute the average value on this state of the operator $\hat{O} = \hat{x} \hat{p} + \hat{p} \hat{x}$. (Points 20%)
- c) Compute the average of \hat{O} at the time t. (Points 20%)

Problem 2

A quantum system is described by a 2-dimensional space of states. In the base $\{|1\rangle, |2\rangle\}$ the Hamiltonian of the system corresponds to the following matrix

$$H = \begin{pmatrix} a & b \\ b & 3 a \end{pmatrix}$$

where a and b are *real* numbers.

- a) What are the normalized energy eigenstates and energy eigenvalues of the system? (Points 20%)
- b) Suppose that the system be at the time t = 0 in the state $\psi(0) = |1\rangle$. What is the probability P(t) that the system is in the ground state at the time t > 0? (Points 20%)
- c) Suppose that the system above is used to describe a molecule made up of two different atoms (or sub-molecules) with two bonding electrons¹.

What is the ground state energy of the molecule in the independent electron approximation and its binding energy? (Points 30%)

¹*Reminder*: In this schematization, when the two atoms are separated one of the two valence electrons is attached to the first atom (state $|1\rangle$) while the other electron localized around the second atom (state $|2\rangle$). When the two atoms get closer, *b* becomes non-vanishing and the valence electrons can be shared by the two atoms (or sub-molecules).

Final written test: 7/1/2014: Solution

Problem 1

a) The normalized state is

$$\psi' = \frac{1}{\sqrt{2}} \left| 0 \right\rangle + i \frac{1}{\sqrt{2}} \left| 1 \right\rangle$$

The ground state satisfies the equation

$$a \left| 0 \right\rangle = \left(\hat{p} - i \, m \, \omega \, \hat{x} \right) \left| 0 \right\rangle = 0$$

In the momentum representation this becomes

$$m\,\omega\,\hbar\,\frac{d\tilde{\psi}_0(p)}{dp} = -p\,\tilde{\psi}_0(p)$$

where $\tilde{\psi}_0(p)$ is the wave function of the ground state $|0\rangle$ in the momentum representation. The solution of this differential equation is

$$\tilde{\psi}_0(p) = C \,\mathrm{e}^{-\frac{p^2}{2\,m\,\omega\,\hbar}}$$

The normalization factor C is fixed by

$$1 = \int_{-\infty}^{\infty} dp \, |\tilde{\psi}_0(p)|^2 = |C|^2 \, \int_{-\infty}^{\infty} dp \, \mathrm{e}^{-\frac{p^2}{m\,\omega\,\hbar}} = |C|^2 \, \sqrt{\pi} \, \sqrt{m\,\omega\,\hbar}$$

We can therefore take

$$\tilde{\psi}_0(p) = \frac{1}{(\pi \, m \, \omega \, \hbar)^{\frac{1}{4}}} \, \mathrm{e}^{-\frac{p^2}{2 \, m \, \omega \, \hbar}} = \frac{2^{\frac{1}{4}}}{(\pi)^{\frac{1}{4}}} \, \frac{\sqrt{\lambda}}{\sqrt{\hbar}} \mathrm{e}^{-\frac{\lambda^2 \, p^2}{\hbar^2}}$$

where $\lambda = \sqrt{\frac{\hbar}{2\omega m}}$. The first excited state is obtain from

$$|1\rangle = a^{\dagger} |0\rangle = \frac{\lambda}{\hbar} \left(\hat{p} + i \, m \, \omega \, \hat{x} \right) |0\rangle$$

Hence

$$\tilde{\psi}_1(p) = C \frac{\lambda}{\hbar} \left(p - m \,\omega \,\hbar \,\frac{d}{dp} \right) e^{-\frac{p^2}{2m\omega\hbar}} =$$
$$= C \frac{\lambda}{\hbar} \left(p + p \right) e^{-\frac{p^2}{2m\omega\hbar}} = C \frac{2 p \,\lambda}{\hbar} e^{-\frac{p^2}{2m\omega\hbar}}$$

In conclusion the normalized wave function of ψ in the momentum representation is

$$\tilde{\psi}(p) = \frac{C}{\sqrt{2}} \left(1 + i \frac{2p\lambda}{\hbar} \right) e^{-\frac{p^2}{2m\omega\hbar}} = \frac{1}{(2\pi)^{\frac{1}{4}}} \frac{\sqrt{\lambda}}{\sqrt{\hbar}} \left(1 + i \frac{2p\lambda}{\hbar} \right) e^{-\frac{\lambda^2 p^2}{\hbar^2}}$$

b) Let us write \hat{O} in terms of a and a^{\dagger} :

$$\hat{O} = \hat{x}\,\hat{p} + \hat{p}\,\hat{x} = \frac{\hbar^2}{4\,\lambda^2}\,\frac{i}{m\,\omega}\left[(a+a^{\dagger})\,(a-a^{\dagger}) + (a-a^{\dagger})\,(a+a^{\dagger})\right] = \\ = \frac{\hbar^2}{4\,\lambda^2}\,\frac{i}{m\,\omega}\left[(a+a^{\dagger})\,(a-a^{\dagger}) + (a-a^{\dagger})\,(a+a^{\dagger})\right] = i\,\hbar\left[a^2 - (a^{\dagger})^2\right]$$

Then

$$\begin{aligned} \langle \psi, \hat{O} \, \psi \rangle &= i \, \hbar \, \langle \psi, [a^2 - (a^{\dagger})^2] (|0\rangle + i \, |1\rangle) \rangle = \\ &= -i \, \hbar \, \langle \psi, (\sqrt{2} \, |2\rangle + i \, \sqrt{6} \, |3\rangle) \rangle = 0 \end{aligned}$$

Hence the average of \hat{O} on this state is zero.

c) At the time t, the state is

$$\psi(t) = \mathrm{e}^{-i\frac{\omega t}{2}} \left(|0\rangle + i \, \mathrm{e}^{-i\,\omega t} \, |1\rangle \right)$$

Hence

$$\begin{aligned} \langle \psi(t), \hat{O}\,\psi(t) \rangle &= \mathrm{e}^{-i\frac{\omega t}{2}} \,i\,\hbar\,\langle\psi(t), [a^2 - (a^{\dagger})^2](|0\rangle + i\,\mathrm{e}^{-i\,\omega\,t}\,|1\rangle) \rangle = \\ &= -i\,\hbar\,\mathrm{e}^{-i\frac{\omega t}{2}}\,\langle\psi, (\sqrt{2}\,|2\rangle + i\,\sqrt{6}\,\mathrm{e}^{-i\,\omega\,t}\,|3\rangle) \rangle = 0 \end{aligned}$$

This means that the average of the operator vanishes for any time t. Problem 2

a)

Energy eigenvalues E are given by the solutions of the equation

$$0 = \det \begin{pmatrix} a - E & b \\ b & 3a - E \end{pmatrix} = E^2 - 4aE + 3a^2 - b^2$$

Thus

$$E_{1,2} = 2 \, a \pm \sqrt{a^2 + b^2}$$

The corresponding eigenvectors $\psi_{1,2} = \begin{pmatrix} x_{1,2} \\ y_{1,2} \end{pmatrix}$ are (for $b \neq 0$)

$$\psi_{1,2} = C_{1,2} \begin{pmatrix} b \\ E_{1,2} - a \end{pmatrix} = C_{1,2} \begin{pmatrix} b \\ a \pm \sqrt{a^2 + b^2} \end{pmatrix}$$

The normalized eigenvectors are determined by the condition

$$1 = |C_{1,2}|^2 2 \left(a^2 + b^2 \pm a \sqrt{a^2 + b^2}\right)$$

which means

$$|C_{1,2}|^2 = \frac{1}{2\sqrt{a^2 + b^2}} \frac{1}{\sqrt{a^2 + b^2} \pm a}$$

The ground state is ψ_2 since $E_2 < E_1$ for any value of a and b.

b) The state $|1\rangle$ writes in the basis $\{\psi_1, \psi_2\}$

$$|1\rangle = C_1 \, b \, \psi_1 + C_2 \, b \, \psi_2$$

The state at the time t is

$$\psi(t) = e^{-i\frac{E_1 t}{\hbar}} C_1 b \psi_1 + e^{-i\frac{E_2 t}{\hbar}} C_2 b \psi_2$$

The probability that at the time t the system is in the ground state ψ_2 is independent of t and is

$$P(t) = |C_2|^2 b^2 = \frac{b^2}{2\sqrt{a^2 + b^2}} \frac{1}{\sqrt{a^2 + b^2} - a}$$

c) The ground state of the molecule is the one in which both electrons occupy the level E_2 , therefore the energy of the ground state of the molecule is

$$E = 2E_2 = 4a - 2\sqrt{a^2 + b^2}$$

The energy of the molecule when the electrons are each around its own atom is

$$E_0 = 3a + a = 4a$$

The binding energy is therefore

$$E_{binding} = E - E_0 = 2\sqrt{a^2 + b^2}$$

Quantum Mechanics: SERP-CHEM Final written test: 27/1/2014

Problem 1

Consider the state ψ of an harmonic oscillator of frequency ω and mass m which satisfies the equation

$$a\,\psi = \alpha\,\psi \tag{96}$$

where α is a *complex number*.

a) Find the *normalized* wave function of ψ in the Schrödinger representation. (Points 20%)

b) Compute Δx and Δp on the state ψ and show it is a state of minimal indeterminacy. (Points 20%)

c) Consider the state

$$\tilde{\psi} = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \left| n \right\rangle \tag{97}$$

where $|n\rangle$ with n = 0, 1, 2, ... are the normalized eigenstates of the harmonic oscillator Hamiltonian. Show that $\tilde{\psi}$ satisfies (96). (Points 20%)

d) Compute the norm of $\tilde{\psi}$. (Points 20%)

Problem 2

Consider the homogenous polynomial of degree 3

$$P(x, y, z) = x y z$$

a) Show that P(x, y, z) harmonic. (Points 10%)

b) Decompose P(x, y, z) as linear combination of harmonic polynomials $H_{3,m}(x, y, z)$ of definite L_3 eigenvalues. (Points 20%)

Final written test: 27/1/2014: Solution

Problem 1

Solution:

a) The wave-function $\psi(x)$ satisfies the differential equation

$$-i\hbar\psi'(x) - im\omega x\psi(x) = \frac{\hbar\alpha}{\lambda}\psi(x)$$
(98)

that is

$$\frac{\psi'(x)}{\psi(x)} = \frac{i\,\alpha}{\lambda} - \frac{x}{2\,\lambda^2} \tag{99}$$

The solution is

$$\psi(x) = C e^{i \alpha \frac{x}{\lambda} - \frac{x^2}{4\lambda^2}}$$
(100)

The normalized wave-function is obtained by imposing

$$1 = \langle \psi, \psi \rangle = |C|^2 \int_{-\infty}^{\infty} dx \, \mathrm{e}^{i \, (\alpha - \bar{\alpha}) \frac{x}{\lambda} - \frac{x^2}{2 \, \lambda^2}} = |C|^2 \, \lambda \int_{-\infty}^{\infty} dy \, \mathrm{e}^{i \, (\alpha - \bar{\alpha}) \, y - \frac{y^2}{2}} = \\ = |C|^2 \, \lambda \, \sqrt{2 \, \pi} \, \mathrm{e}^{-\frac{(\alpha - \bar{\alpha})^2}{2}} \tag{101}$$

Hence

$$\psi(x) = \frac{1}{(2\pi)^{\frac{1}{4}}} \frac{1}{\lambda^{\frac{1}{2}}} e^{\frac{(\alpha - \bar{\alpha})^2}{4}} e^{i\,\alpha\frac{x}{\lambda} - \frac{x^2}{4\,\lambda^2}}$$
(102)

b) Let us compute Δx by making use of the Schrödinger representation:

$$\langle \psi, \hat{x} \psi \rangle = |C|^2 \int_{-\infty}^{\infty} dx \, x \, \mathrm{e}^{i \, (\alpha - \bar{\alpha}) \frac{x}{\lambda} - \frac{x^2}{2 \, \lambda^2}} = |C|^2 \, \lambda^2 \int_{-\infty}^{\infty} dy \, y \, \mathrm{e}^{i \, (\alpha - \bar{\alpha}) \, y - \frac{y^2}{2}} =$$
$$= -i \, |C|^2 \, \lambda^2 \, \sqrt{2 \pi} \, \frac{\partial}{\partial (\alpha - \bar{\alpha})} \mathrm{e}^{-\frac{(\alpha - \bar{\alpha})^2}{2}} =$$
$$= i \, |C|^2 \, \lambda^2 \, \sqrt{2 \pi} \, (\alpha - \bar{\alpha}) \, \mathrm{e}^{-\frac{(\alpha - \bar{\alpha})^2}{2}} = i \, (\alpha - \bar{\alpha}) \, \lambda$$
(103)

and

$$\langle \psi, \hat{x}^2 \,\psi \rangle = |C|^2 \,\int_{-\infty}^{\infty} dx \, x^2 \mathrm{e}^{i \,(\alpha - \bar{\alpha}) \frac{x}{\lambda} - \frac{x^2}{2\,\lambda^2}} = |C|^2 \,\lambda^3 \int_{-\infty}^{\infty} dy \, y^2 \,\mathrm{e}^{i \,(\alpha - \bar{\alpha}) \,y - \frac{y^2}{2}} = \\ = -|C|^2 \,\lambda^3 \,\sqrt{2\,\pi} \,\frac{\partial^2}{\partial(\alpha - \bar{\alpha})^2} \,\mathrm{e}^{-\frac{(\alpha - \bar{\alpha})^2}{2}} = \\ = |C|^2 \,\lambda^3 \,\sqrt{2\,\pi} \left(1 - (\alpha - \bar{\alpha})^2\right) \,\mathrm{e}^{-\frac{(\alpha - \bar{\alpha})^2}{2}} = \lambda^2 \left(1 - (\alpha - \bar{\alpha})^2\right) \tag{104}$$

Thus

$$\Delta x^2 = \lambda^2 \left[1 - (\alpha - \bar{\alpha})^2 + (\alpha - \bar{\alpha})^2 \right] = \lambda^2 \tag{105}$$

For Δp , let us show how to perform the computation by using raising and lowering operators together with Eq. (96). As a matter of fact, this method is simpler:

$$\langle \psi, \hat{p} \, \psi \rangle = \frac{\hbar}{2\,\lambda} \, \langle \psi, (a + a^{\dagger}) \, \psi \rangle = \frac{\hbar}{2\,\lambda} \, \langle \psi, (\alpha + a^{\dagger}) \, \psi \rangle =$$
$$= \frac{\hbar}{2\,\lambda} \, (\alpha + \langle a \, \psi, \psi \rangle) = \frac{\hbar}{2\,\lambda} \, (\alpha + \bar{\alpha})$$
(106)

And

$$\langle \psi, \hat{p}^2 \psi \rangle = \frac{\hbar^2}{4\lambda^2} \langle (a + a^{\dagger})\psi, (a + a^{\dagger})\psi \rangle =$$

$$= \frac{\hbar^2}{4\lambda^2} \left[|\alpha|^2 + \langle a^{\dagger}\psi, a^{\dagger}\psi \rangle + \alpha^2 + \bar{\alpha}^2 \right] =$$

$$= \frac{\hbar^2}{4\lambda^2} \left[|\alpha|^2 + |\alpha|^2 + 1 + \alpha^2 + \bar{\alpha}^2 \right] =$$

$$= \frac{\hbar^2}{4\lambda^2} \left[2 |\alpha|^2 + 1 + \alpha^2 + \bar{\alpha}^2 \right]$$

$$(107)$$

Hence

$$\Delta p^2 = \frac{\hbar^2}{4\,\lambda^2} \left[2\,|\alpha|^2 + 1 + \alpha^2 + \bar{\alpha}^2 - \alpha^2 - \bar{\alpha}^2 - 2\,|\alpha|^2 \right] = \frac{\hbar^2}{4\,\lambda^2} \tag{108}$$

We verified that the state is indeed a state of minimal indeterminacy

$$\Delta x \, \Delta p = \frac{\hbar}{2} \tag{109}$$

c)

$$a\,\tilde{\psi} = \sum_{n=1}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \sqrt{n} \,|n-1\rangle = \sum_{n=0}^{\infty} \frac{\alpha^{n+1}}{\sqrt{(n+1)!}} \sqrt{n+1} \,|n\rangle =$$
$$= \alpha \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \,|n\rangle = \alpha\,\tilde{\psi}$$
(110)

d)

$$\langle \tilde{\psi}, \tilde{\psi} \rangle = \sum_{n=0}^{\infty} \frac{\bar{\alpha}^n}{\sqrt{n!}} \frac{\alpha^n}{\sqrt{n!}} = \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} = e^{|\alpha|^2}$$
(111)

Problem 2

 ${\cal P}(x,y,z)$ is homogenous of degree 3 and harmonic

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right) x y z = 0$$
(112)

We have

$$H_{3,\pm3} = x_{\pm}^{3} \qquad H_{3,\pm2} = x_{\pm}^{2} z$$

$$H_{3,1} = x_{\pm} z^{2} - \frac{1}{4} x_{\pm}^{2} x_{\pm} \qquad H_{3,-1} = x_{\pm} z^{2} - \frac{1}{4} x_{\pm}^{2} x_{\pm}$$

$$H_{3,0} = z^{3} - \frac{3}{2} x_{\pm} x_{\pm} z$$
(113)

Hence

$$P(x, y, z) = \frac{(x_{+}^{2} - x_{-}^{2})}{4i} z = \frac{H_{3,2} - H_{3,-2}}{4i}$$
(114)

Quantum Mechanics: SERP-CHEM Final written test: 8/1/2016

Problem 1

A particle of spin 1/2 is subject to a uniform magnetic field which has components both along the y direction and the z direction: the interaction with the magnetic field is described by

$$\hat{H} = -\mu \ B_y \, \hat{s}_y - \mu \, B_z \, \hat{s}_z$$

with $\mu > 0$. Suppose that at the time t = 0 the particle is in the state with definite value $+\frac{\hbar}{2}$ of \hat{s}_z .

- a) What is the probability that at the time t = 0 the measurement of \hat{s}_x gives the value $+\frac{\hbar}{2}$?(Points 10%)
- b) What are the energy eigenvalues and the normalized energy eigenvectors in terms of the eigenvectors $|\pm\rangle$ of \hat{s}_z ? (Points 40 %)
- c) What is the probability that at the time t > 0 the particle be in the lowest energy state? Suppose $B_y = \sqrt{3} B_z$, give the numerical value for this probability. (Points 20%)
- d) What is the probability that at the time t > 0 the measurement of \hat{s}_z gives the value $+\frac{\hbar}{2}$? (Points 40 %)

Final written test: 8/1/2016: Solution

Problem 1

• a) The state with $s_x = \frac{\hbar}{2}$ is

$$|s_x = +\rangle = \frac{1}{\sqrt{2}} \left(|+\rangle + |-\rangle \right)$$

The probability that at the time t=0 the measurement of \hat{s}_x gives the value $+\frac{\hbar}{2}$ is therefore

$$P(s_x = \frac{\hbar}{2}) = |\langle s_x = +|+\rangle|^2 = \frac{1}{2}$$

• b) The energy eigenvalues are determined by the equation

$$\det \begin{pmatrix} -\frac{\hbar}{2} \mu B_z - E & i\frac{\hbar}{2} \mu B_y \\ -i\frac{\hbar}{2} \mu B_y & \frac{\hbar}{2} \mu B_z - E \end{pmatrix} = E^2 - \frac{\hbar^2}{4} \mu^2 \left(B_z^2 + B_y^2 \right) = 0$$

that is

$$E_{\pm} = \pm \frac{\hbar}{2} \,\mu \sqrt{B_z^2 + B_y^2} \equiv \pm \frac{\hbar}{2} \,\mu \left| \vec{B} \right|$$

The corresponding eigenstates are

$$\psi_{\pm} = x_{\pm} \left| + \right\rangle + y_{\pm} \left| - \right\rangle$$

where

$$i B_y y_{\pm} = \left(\frac{E_{\pm}}{\frac{\hbar}{2}\mu} + B_z\right) x_{\pm} = \left(\pm |\vec{B}| + B_z\right) x_{\pm}$$

For normalized eigenvectors we must take

$$|x_{\pm}|^2 + |y_{\pm}|^2 = 1$$

$$y_{\pm} = N_{\pm} (\pm |\vec{B}| + B_z) \qquad x_{\pm} = N_{\pm} i B_y$$
$$\frac{1}{N_{\pm}^2} = (B_z \pm |\vec{B}|)^2 + B_y^2 = 2 |\vec{B}| (|\vec{B}| \pm B_z)$$
$$N_{\pm}^2 = \frac{|\vec{B}| \mp B_z}{2 |\vec{B}| B_y^2}$$
$$x_{\pm} = i \sqrt{\frac{|\vec{B}| \mp B_z}{2 |\vec{B}|}} \qquad y_{\pm} = \pm \sqrt{\frac{|\vec{B}| \pm B_z}{2 |\vec{B}|}}$$

with N_{\pm} real. At the time t = 0 the system is in the state $|+\rangle$, whose decomposition in terms of energy eigenstates is

$$\psi(0) = |+\rangle = \psi_+ \langle \psi_+ |+\rangle + \psi_- \langle \psi_- |+\rangle =$$
$$= \psi_+ x_+^* + \psi_- x_-^*$$

At the time t the system is in the state

$$\begin{split} \psi(t) &= \psi_{+} x_{+}^{*} e^{-i \frac{\hbar}{2} \frac{\mu |\vec{B}|}{\hbar} t} + \psi_{-} x_{-}^{*} e^{i \frac{\hbar}{2} \frac{\mu |\vec{B}|}{\hbar} t} = \\ &= (x_{+} |+\rangle + y_{+} |-\rangle) x_{+}^{*} e^{-i \frac{\mu |\vec{B}|}{2} t} + \\ &+ (x_{-} |+\rangle + y_{-} |-\rangle) x_{-}^{*} e^{i \frac{\mu |\vec{B}|}{2} t} = \\ &= \left(|x_{+}|^{2} e^{-i \frac{\mu |\vec{B}|}{2} t} + |x_{-}|^{2} e^{i \frac{\mu |\vec{B}|}{2} t} \right) |+\rangle + \left(y_{+} x_{+}^{*} e^{-i \frac{\mu |\vec{B}|}{2} t} + \\ &+ y_{-} x_{-}^{*} e^{i \frac{\mu |\vec{B}|}{2} t} \right) |-\rangle \end{split}$$

• c) The probability that at the time t > 0 the particle be in the lowest energy state is

$$P_{\psi_{-}}(t) = \left| \langle \psi_{-} | \psi(t) \rangle \right|^{2} = |x_{-}^{*} e^{-i \frac{\mu |\vec{B}|}{2} t}|^{2} = |x_{-}^{*}|^{2} = N_{-}^{2} B_{y}^{2} = \frac{1}{2} \left(1 + \frac{B_{z}}{|\vec{B}|} \right)$$

For $B_y = \sqrt{3} B_z$ we have

$$P_{\psi_{-}}(t) = \frac{1}{2}\left(1 + \frac{1}{2}\right) = \frac{3}{4}$$

- d) The probability *amplitude* that at the time t the measurement of \hat{s}_z gives the value $+\frac{\hbar}{2}$ is

$$\langle +|\psi(t)\rangle = \langle +|\psi_{+}\rangle \, x_{+}^{*} \, \mathrm{e}^{-i\frac{\mu|\vec{B}|}{2}t} + \langle +|\psi_{-}\rangle \, x_{-}^{*} \, \mathrm{e}^{i\frac{\mu|\vec{B}|}{2}t} = \\ = |x_{+}|^{2} \, \mathrm{e}^{-i\frac{\mu|\vec{B}|}{2}t} + |x_{-}|^{2} \, \mathrm{e}^{i\frac{\mu|\vec{B}|}{2}t}$$

The corresponding probability is

$$P_{+}(t) = \left(|x_{+}|^{2} e^{-i\frac{\mu|\vec{B}|}{2}t} + |x_{-}|^{2} e^{i\frac{\mu|\vec{B}|}{2}t}\right) \left(|x_{+}|^{2} e^{i\frac{\mu|\vec{B}|}{2}t} + |x_{-}|^{2} e^{-i\frac{\mu|\vec{B}|}{2}t}\right) =$$

$$= |x_{+}|^{4} + |x_{-}|^{4} + 2|x_{+}|^{2}|x_{-}|^{2} \cos \mu |\vec{B}| t =$$

$$= B_{y}^{4} \left(N_{+}^{4} + N_{-}^{4} + 2N_{+}^{2}N_{-}^{2}\frac{\mu|\vec{B}|t}{\hbar}\right) =$$

$$= \frac{1}{4|\vec{B}|^{2}} \left((|\vec{B}| - B_{z})^{2} + (|\vec{B}| + B_{z})^{2} + 2B_{y}^{2} \cos \mu |\vec{B}| t\right) =$$

$$= \frac{1}{2|\vec{B}|^{2}} \left(|\vec{B}|^{2} + B_{z}^{2} + B_{y}^{2} \cos \mu |\vec{B}| t\right) =$$

$$= 1 - \frac{B_{y}^{2}}{2|\vec{B}|^{2}} \left(1 - \cos \mu |\vec{B}| t\right) = 1 - \frac{B_{y}^{2}}{|\vec{B}|^{2}} \sin^{2} \frac{\mu |\vec{B}|t}{2}$$

Quantum Mechanics: SERP-CHEM Final written test: 1/2/2016

Problem 1

An electron moves freely in a 1-dimensional box of length $L = 10^{-10}m$. At the time t = 0 the particle is in the state described by the wave-function

$$\psi_0(x) = \sin \frac{\pi x}{L} + \frac{\sqrt{2}}{3} i \sin \frac{4\pi x}{L}$$

- a) What is the average energy in electronvolts of the particle at the time t = 0? What is the average energy at the time t > 0? (Points 20%)
- b) What is the average of \hat{x} at the time t = 0? What is the average of \hat{x} at the time t > 0? (Points 30 %)
- c) Compute the probability that at the time t > 0 the system be in the state $\psi_f(x) = \sin \frac{\pi x}{L} \sin \frac{4\pi x}{L}$. (Points 20%)

Problem 2

The space of states for electron in a diatomic molecule is a 2-dimensional vector space with the orthonormal basis $\{|1\rangle, |2\rangle\}$ where $|1\rangle$ ($|2\rangle$) are the states in which the electron is attached to the first (second) molecule. The energy levels of the electron are described by the hamiltonian

$$H = \begin{pmatrix} E_0 & c \\ c & E_1 \end{pmatrix}$$

where c is a real parameter and $E_1 - E_0 = \Delta E > 0$. Suppose $\Delta E = 1eV$ and c = 0.67eV. (Points 40%)

• a) What is the probability that the electron in the ground state is attached to the first molecule?

Final written test: 1/2/2016: Solution

Problem 1

• a) What is the average energy of the particle at the time t = 0? What is the average energy at the time t > 0?

The normalized wave function at the time t = 0 is

$$\hat{\psi}_0 = \frac{3}{\sqrt{11}} \left(\psi_1 + \frac{\sqrt{2}}{3} i \, \psi_4 \right) \tag{115}$$

where $\psi_n = \sqrt{\frac{2}{L}} \sin \frac{n\pi}{L}$ are the normalized energy eigenfunctions. Hence

$$\langle H \rangle = \frac{\hbar^2 \pi^2}{2 m L^2} \left(\frac{9}{11} + 16 \frac{2}{11} \right) = \frac{\hbar^2 \pi^2}{2 m L^2} \frac{41}{11} = = \frac{41}{11} 13.6 \, eV \, \frac{a_{Bohr}^2}{L^2} = 14.24 \, eV$$
 (116)

which does not depend on time.

• b) At time t = 0:

$$\langle \hat{x} \rangle = \langle \hat{\psi}_{0}, \hat{x} \, \hat{\psi}_{0} \rangle = \frac{9}{11} \langle \psi_{1}, \hat{x} \, \psi_{1} \rangle + \frac{2}{11} \langle \psi_{4}, \hat{x} \, \psi_{4} \rangle + + \frac{3\sqrt{2}i}{11} \left(\langle \psi_{1}, \hat{x} \, \psi_{4} \rangle - \langle \psi_{4}, \hat{x} \, \psi_{1} \rangle \right) = = \frac{9}{11} \frac{L}{2} + \frac{2}{11} \frac{L}{2} + + \frac{3\sqrt{2}i}{11} \left(\langle \psi_{1}, \hat{x} \, \psi_{4} \rangle - \langle \psi_{4}, \hat{x} \, \psi_{1} \rangle \right) = \frac{L}{2}$$
(117)

Putting $\Delta E = E_4 - E_1 = \frac{15 \hbar^2 \pi^2}{2 m L^2}$, we have at time t > 0

$$\langle \hat{x}(t) \rangle = \langle \hat{\psi}_{0}(t), \hat{x} \, \hat{\psi}_{0}(t) \rangle = \frac{9}{11} \langle \psi_{1}, \hat{x} \, \psi_{1} \rangle + \frac{2}{11} \langle \psi_{4}, \hat{x} \, \psi_{4} \rangle + + \frac{3\sqrt{2}i}{11} \left(e^{-i\frac{\Delta Et}{\hbar}} \langle \psi_{1}, \hat{x} \, \psi_{4} \rangle - e^{i\frac{\Delta Et}{\hbar}} \langle \psi_{4}, \hat{x} \, \psi_{1} \rangle \right) = = \frac{9}{11} \frac{L}{2} + \frac{2}{11} \frac{L}{2} + + \frac{3\sqrt{2}i}{11} \frac{(-32)L}{225\pi^{2}} \left(e^{-i\frac{\Delta Et}{\hbar}} - e^{i\frac{\Delta Et}{\hbar}} \right) = = \frac{L}{2} - \frac{\sqrt{2} \, 64 \, L}{825 \, \pi^{2}} \sin \frac{\Delta E t}{\hbar}$$
(118)

• c)

The norm of ψ_f is

$$\langle \psi_f, \psi_f \rangle = \frac{L}{2} + \frac{L}{2} = L \tag{119}$$

and the normalized final state is

$$\hat{\psi}_f(x) = \frac{1}{\sqrt{2}} \left(\psi_1(x) - \psi_4(x) \right)$$
(120)

The transition amplitude is

$$\langle \hat{\psi}_f, \hat{\psi}_0(t) \rangle = \frac{1}{\sqrt{2}} \frac{3}{\sqrt{11}} e^{-i\frac{E_1 t}{\hbar}} - \frac{1}{\sqrt{2}} \frac{\sqrt{2}}{\sqrt{11}} i e^{-i\frac{E_4 t}{\hbar}}$$
(121)

and the probability amplitude is

$$P_{i \to f}(t) = \frac{9}{2 \times 11} + \frac{1}{11} - \frac{6}{11\sqrt{2}} \sin \frac{\Delta E t}{\hbar} = \frac{1}{2} - \frac{6}{11\sqrt{2}} \sin \frac{\Delta E t}{\hbar}$$
(122)

Problem 2

• a)

The energy levels of the electron are

$$E_{\pm} = E_0 + \frac{\Delta E \pm \sqrt{\Delta E^2 + 4c^2}}{2}$$
(123)

The corresponding eigenstates are

$$|E_{\pm}\rangle = x_{\pm}|1\rangle + y_{\pm}|2\rangle \tag{124}$$

where

$$(E_{\pm} - E_0) x_{\pm} = c y_{\pm} \tag{125}$$

The (not normalized) ground state is

$$|E_{-}\rangle = |1\rangle + \frac{\Delta E - \sqrt{\Delta E^2 + 4c^2}}{2c} |2\rangle \tag{126}$$

The probability for an electron in the ground state to be attached to the first atom is

$$P_{1} = \frac{|\langle 1|E_{-}\rangle|^{2}}{\langle E_{-}|E_{-}\rangle} = \frac{1}{1 + \frac{(\Delta E - \sqrt{\Delta E^{2} + 4c^{2}})^{2}}{4c^{2}}} = \frac{4x^{2}}{4x^{2} + (1 - \sqrt{1 + 4x^{2}})^{2}}$$
(127)

where

$$x = \frac{c}{\Delta E} \tag{128}$$

For

$$x = 0.67$$
 (129)

$$P_1 = 0.80 \tag{130}$$